Time evolutions of quantum systems and quantum Zeno effect

> Saverio Pascazio Dipartimento di Fisica and INFN Bari, Italy

> > Paul Scherrer Institut, 19 April 2024

### Zeno of Elea

Zeno was an Eleatic philosopher, a native of Elea in Italy, son of Teleutagoras, and the favorite disciple of Parmenides. He was born about 488 BC, and at the age of forty accompanied Parmenides to Athens





#### The flying arrow is at rest.

At any given moment it is in a space equal to its own length, and therefore is at rest at that moment. So, it is at rest at all moments. The sum of an infinite number of these positions of rest is not a motion.



### Quantum Zeno effect quantum system $\psi$ Hamiltonian HSchrödinger equation $\psi_t = e^{-iHt}\psi_0$ $\mathcal{A}(t) = \langle \psi_0 | \psi(t) \rangle = \langle \psi_0 | e^{-iHt} | \psi_0 \rangle$ $p(t) = |\mathcal{A}(t)|^2 = |\langle \psi_0 | e^{-iHt} | \psi_0 \rangle|^2$ survival amplitude and probability

(always valid; both in QM and QFT)

 $p(t) = |\langle \psi_t | \psi_0 \rangle|^2 = 1 - t^2 / \tau_z^2$  $\tau_{\rm Z}^{-2} \equiv \langle \psi_0 | H^2 | \psi_0 \rangle - \langle \psi_0 | H | \psi_0 \rangle^2$ 



# $[p(t/N)]^N \simeq \left(1 - \frac{t^2}{\tau_Z^2 N^2}\right)^N \stackrel{N \to \infty}{\longrightarrow} 1$

Misra and Sudarshan 1977

### always valid, also for "unstable" systems





$$p(t) = \left| \left\langle \Psi(0) \left| \Psi(t) \right\rangle \right|^2 = \cos\left(\frac{t}{\tau_z}\right)$$
$$p^{(N)}(t) = p\left(\frac{t}{N}\right)^N = \cos\left(\frac{t}{N\tau_z}\right)^{2N}$$
$$p^{(N)}(t) \xrightarrow{N \to \infty} 1$$

P., Namiki, Badurek, Rauch, Phys. Lett. A 169, 155 (1993)

$$[p(t/N)]^{N} \simeq \left(1 - \frac{t^{2}}{\tau_{Z}^{2}N^{2}}\right)^{N} \stackrel{N \to \infty}{\longrightarrow} 1$$
Quantum Zeno effect



. .

many experiments on many physical systems (and applications: quantum control)

### Quantum Zeno effect

- A consequence of general principles of quantum physics
- A quantum measurement perturbs the system under observation
- And entails a projection ("collapse") of the wave function
- Projection onto the state that is the outcome of the measurement

$$\mathcal{A}(t) = \langle \psi_0 | \psi(t) \rangle = \langle \psi_0 | e^{-iHt} | \psi_0 \rangle$$
$$p(t) = |\mathcal{A}(t)|^2 = |\langle \psi_0 | e^{-iHt} | \psi_0 \rangle|^2$$

Fourier/Laplace transform

$$\mathcal{A}(t) = \int_{B} \frac{dE}{2\pi} e^{-iEt} \mathcal{A}(E) = \frac{i}{2\pi} \int_{B} dE \frac{e^{-iEt}}{E - \omega_0 - \Sigma(E)}$$
  
Bromwich path   
B

$\mathcal{A}(t) = \frac{i}{2\pi} \int_{B} dE \; \frac{e^{-iEt}}{E - \omega_0 - \Sigma(E)}$	$\mathcal{A} = \mathcal{A}_{\text{pole}}(t) + \mathcal{A}_{\text{cut}}(t)$ "small"
$E_{\rm pole} - \omega_0 - \Sigma_{\rm II}(E_{\rm pole}) = 0$	
Weisskopf-Wigner	ω
$P(t) \simeq  \mathcal{A}_{\text{pole}}(t) ^2 = Ze^{-\gamma t}$ $Z =  1 - \Sigma'_{\text{II}}(E_{\text{pole}}) ^{-2}$ wave-function renormalization	× E <sub>pole</sub> II
Araki etal (1957), Schwinger 1960	

### notice

# wave-function renormalization





Schwinger (simple poles); Arakí et al (proof of Fermí "Golden rule")



### 

in general

### therefore

Z (wf renormalisation)
pole + cut
what else?

Zeno

# Paley and Wiener

 $\int_{-\infty}^{+\infty} \frac{|\log p(t)|}{1+t^2} \mathrm{d}t < +\infty$ 

if spectrum of H bounded from below

incompatible with  $p(t) \sim e^{-\alpha t^{\beta}}$ 



### in a real system











#### PHYSICAL REVIEW LETTERS



#### From the Quantum Zeno to the Inverse Quantum Zeno Effect

P. Facchi,<sup>1,3</sup> H. Nakazato,<sup>2</sup> and S. Pascazio<sup>3</sup>

<sup>1</sup>Atominstitut der Österreichischen Universitäten, Stadionallee 2, A-1020, Wien, Austria <sup>2</sup>Department of Physics, Waseda University, Tokyo 169-8555, Japan <sup>3</sup>Dipartimento di Fisica, Università di Bari and Istituto Nazionale di Fisica Nucleare, Sezione di Bari, I-70126 Bari, Italy (Received 21 June 2000)

The temporal evolution of an unstable quantum mechanical system undergoing repeated measurements is investigated. In general, by changing the time interval between successive measurements, the decay can be accelerated (inverse quantum Zeno effect) or slowed down (quantum Zeno effect), depending on the features of the interaction Hamiltonian. A geometric criterion is propried for a transition to occur between these two regimes.

DOI: 10.1103/PhysRevLett.86.2699



#### $Z \simeq 1.23$

n-

ın

is

 $x(t) = \sqrt{Z} e^{-\gamma_0 t/2 - i\alpha(t)} + x_{\text{cut}}(t),$ (9)  $Z = |1 - \Sigma'(E_{\text{pole}})|^{-2}$ ,

where the exponential term (first term) is due to the contribution of a simple pole  $E_{pole}$  on the second Riemannian sheet in the complex energy plane, while the second term is the result of a contour integration [1]. The lifetime  $\gamma_0^{-1}$  is given by the Fermi golden rule, computed according to the Weisskopf-Wigner approximation. The quantity Z is the square of the residue of pole of the propagator (yielding wave function renormalization in quantum field the formula to the formula to the formula term.

#### alth 1gh for a stable state Z < 1 (due to probability c servition in the Källén-Lehmann representation), for uns

of order (coupling constant)<sup>2</sup> and modifies the exponential law both at short and long times, yielding the characteristic quadratic and power-law behaviors. The survival probabilpower law





 $p(t) \sim \frac{1}{t^3}$ 

#### notice interference

### back to short times

#### Seminal ideas

Recent experiments



$$p(t) = \left| \left\langle \Psi(0) \left| \Psi(t) \right\rangle \right|^2 = \cos\left(\frac{t}{\tau_z}\right)$$
$$p^{(N)}(t) = p\left(\frac{t}{N}\right)^N = \cos\left(\frac{t}{N\tau_z}\right)^{2N}$$
$$p^{(N)}(t) \xrightarrow{N \to \infty} 1$$

P., Namiki, Badurek, Rauch, Phys. Lett. A 169, 155 (1993)

von Neumann, 1932 rization, nuclear spin isomers HISTORY Beskow and Nilsson, 1967, neutron spin Landau-Zener tunneling Bose-Einstein condensates Khalfin 1968 Friedman 1972 Applications: decoherence in quantum computing Misra and Spearshan, 1977 efficient preservation of spin polarized gases Zeno tomography (MAIN) EXPERIMENTS (Cook 1988) Itano, Heinzen, Bollinger, and Wineland 1990 Nagels, Hermans, and Chapovsky 1997 Kwiat, White, Mitchell, Nairz, Weihs, Weinfurter and Zeilinger 1999 Wunderlich, Balzer, and Toschek, 2001 Fischer, Gutierrez-Medina, Raizen, 2001 Streed, Mun, Boyd, Campbell, Medley, Ketterle, Pritchard, 2006 Bernu, Sayrin, Kuhr, Dotsenko, Brune, Raimond, Haroche 2008 ONE DIMENSIONAL ZENO EFFECT!



system remains in subspace defined by (negative result, nonselective) measurement what if there is a Hamiltonian? Quantum Zeno DYNAMICS in Quantum Zeno SUBSPACE

 $|m\rangle$ 



#### Zeno subspace



Continuous coupling

(continuous measurement)

#### e.g.: let us look at continuous coupling



### What causes the Zeno effect

#### Zeno subspace





Coupling or N so

### Quantum Zeno subspaces

Dynamical superselection sectors





### Experiments

Experimental realization of quantum zeno dynamics
F. Schäfer, I. Herrera, S. Cherukattil, C. Lovecchio, F.S. Cataliotti, F. Caruso & A. Smerzi
Nature Communications 5, 3194 (2014)

Confined quantum Zeno dynamics of a watched atomic arrow Adrien Signoles, Adrien Facon, Dorian Grosso, Igor Dotsenko, Serge Haroche, Jean-Michel Raimond, Michel Brune & Sébastien Gleyzes Nature Physics 10, 715–719 (2014)

Quantum dynamics of an electromagnetic mode that cannot have N photons L. Bretheau, P. Campagne-Ibarcq, E. Flurin, F. Mallet, B. Huard Science, 348, 776-779 (2015)



### Hamiltonians



 $P = 1 - |s\rangle\langle s|$ 

# laser $H = \alpha a^{\dagger} + \alpha^* a$ cavity $V = \frac{\hbar\Omega}{2} (|e\rangle \langle g|a + |g\rangle \langle e|a^{\dagger})$

implement measurement of s

### Zeno Hamíltonían $H_Z = P_{<s}HP_{<s} + P_{>s}HP_{>s} = H_{<s} + H_{>s}$



5=4

## s=6: phase space view



### Quantum Zeno dynamics



exclusion circle in phase space

think of the way the measurement is performed!

#### FIRST EXPERIMENT Schaefer, Herrera, Cherukattil, Lovecchio, Cataliotti, Caruso, Smerzi Florence 2014





Free particle in n dimensions  $H = \frac{p^2}{2M} = -\frac{\hbar^2 \Delta}{2M}, \quad U(t) = \exp(-iHt/\hbar) \quad \text{in } L^2(\mathbb{R}^n)$  $\Omega \subset \mathbb{R}^n$  compact domain,  $P = \chi_{\Omega}(x)$  spatial projection Zeno dynamics  $V_Z(t) = \lim_{N \to \infty} \left[ V\left(\frac{t}{N}\right) \right]^N$ , V(s) = PU(s)PHow does the particle move inside  $\Omega$ ? Does it leak out? The weak limit  $V_Z(t)$  exists and yields  $V_Z(t) = \exp(-iH_\Omega t/\hbar)P,$ a unitary group in  $L^2(\Omega)$ with Zeno Hamiltonian  $H_{\Omega} = -\frac{\hbar^2 \Delta}{2M}, \quad D(H_{\Omega}) = H^2(\Omega) \cap H_0^1(\Omega)$ Free particle in a box with perfectly Dirichlet boundary conditions reflecting hard walls

Facchi, Marmo, Pascazio, Scardicchio, Sudarshan 2003 ....although there is NO wall! Exner and Ichinose, 2005

#### School of Athens (Raphael's Room of the Signature, Vatican Museums)

