

SUSY QCD corrections to electroweak gauge boson production with an associated jet

Ryan Gavin
University of Wisconsin Madison

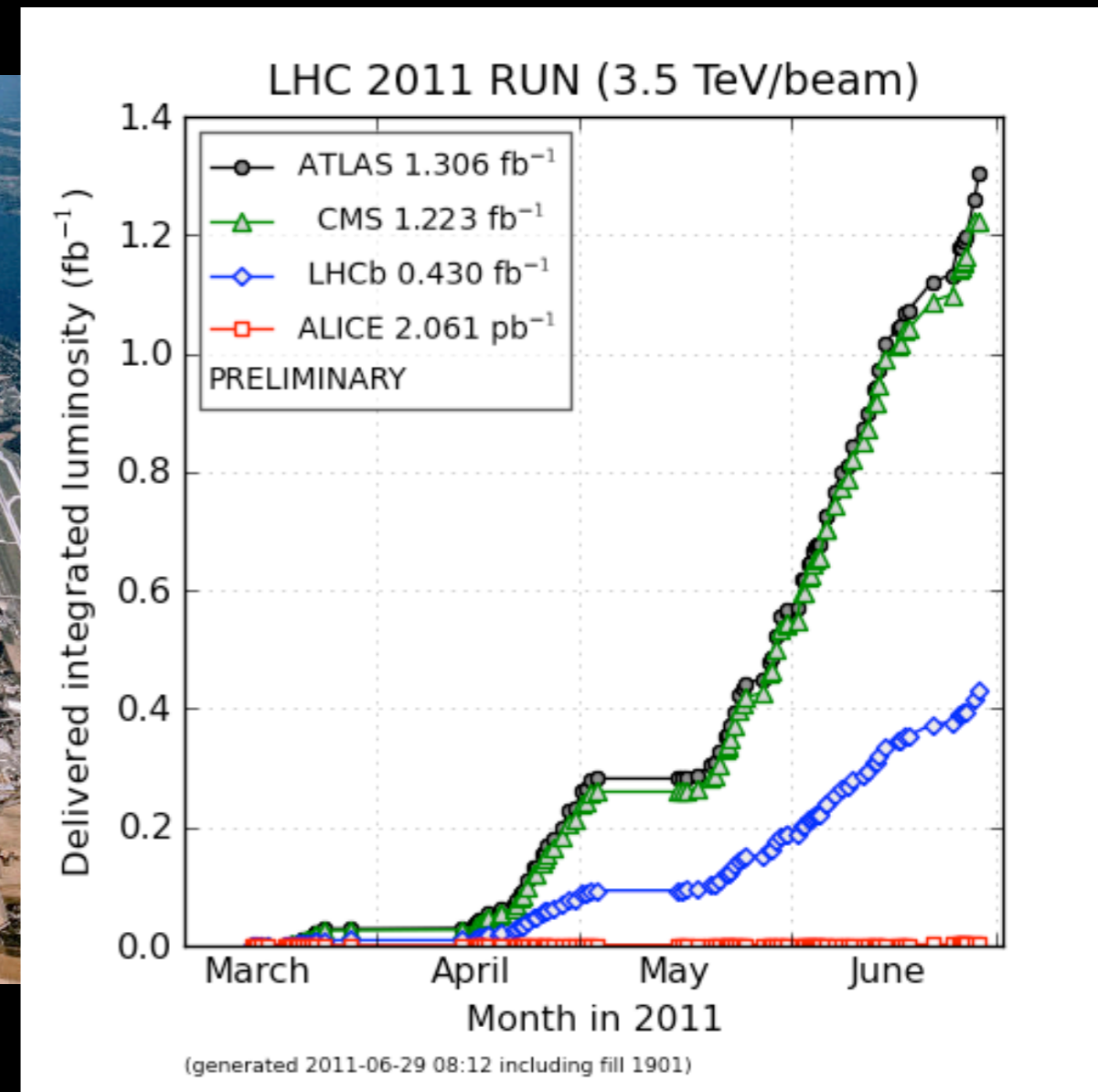
with F Petriello, M Trenkel
in progress

PSI Particle Theory Seminar
June 30, 2011

Outline

- EW Gauge Boson Production
- at Higher Orders
- sQCD Corrections to Z+Jet
- Numerical Results
- sQCD Corrections to W+Jet
- Concluding Remarks

The LHC



Running at 7 TeV since
March 13, 2011

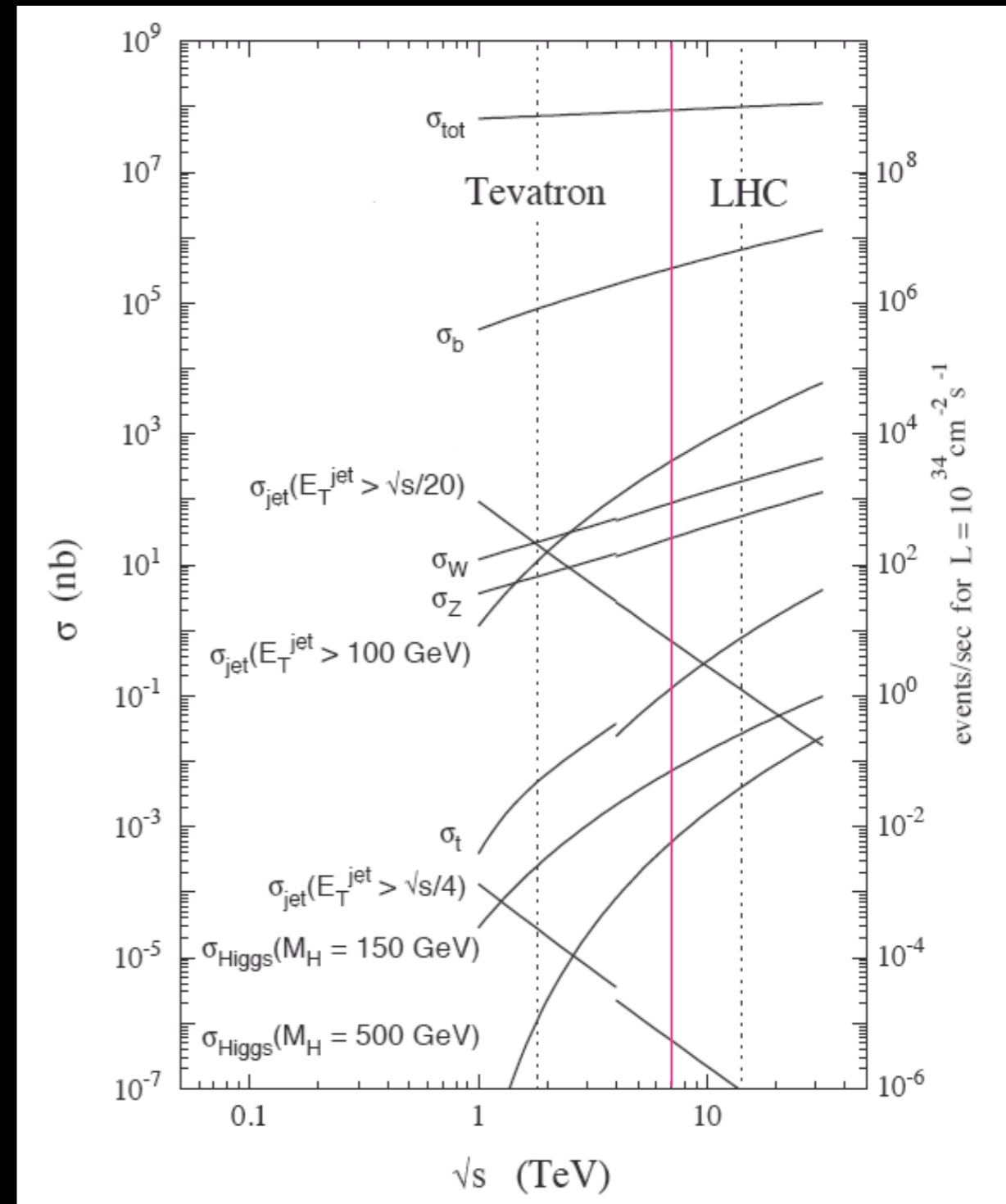
$$\mathcal{L}_{int} \sim 1.3 \text{ fb}^{-1}$$

Z & W's at Hadron Colliders

- Z & W production still very interesting
 - playing an important role in LHC physics

Z & W's at Hadron Colliders

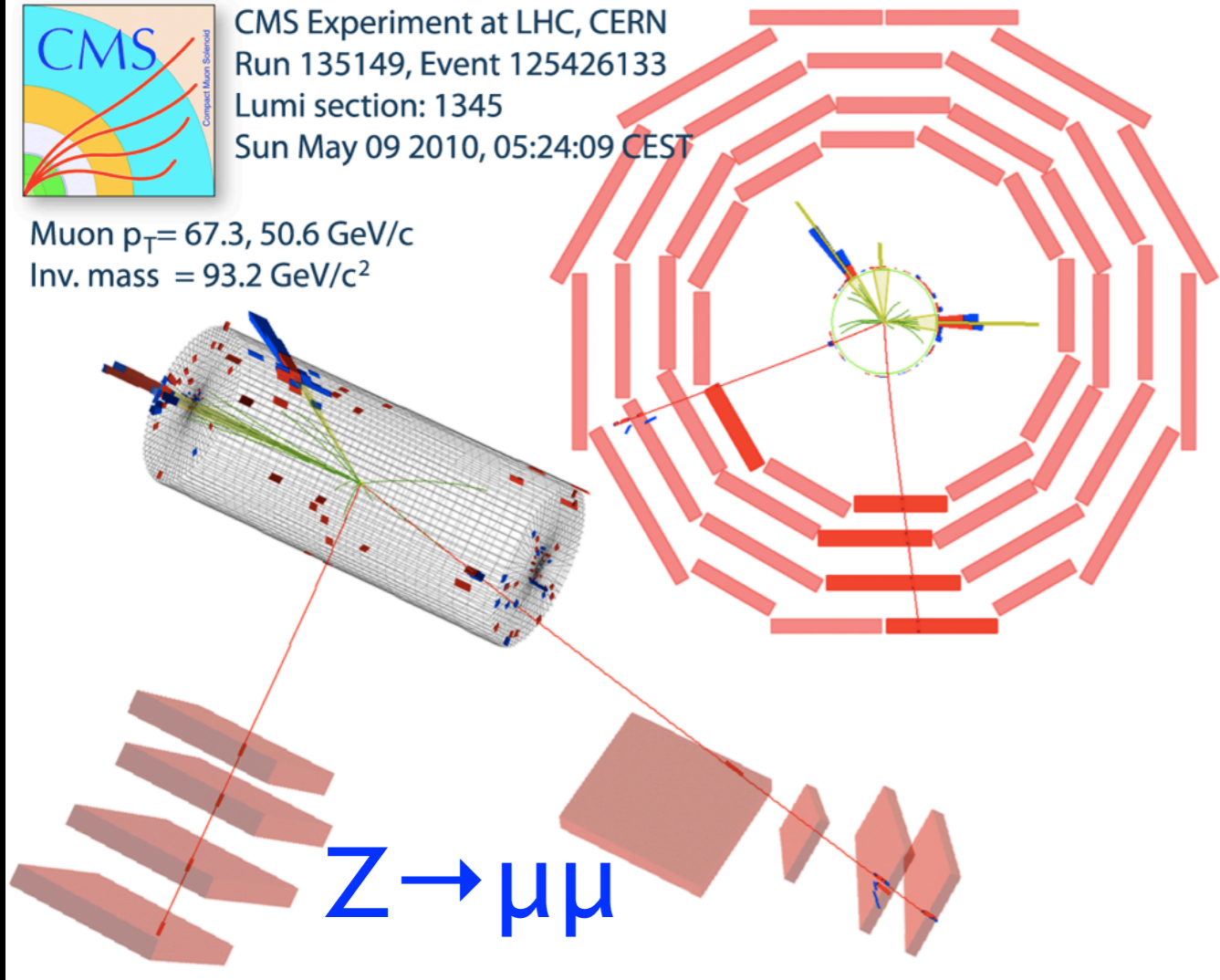
- Z & W production still very interesting
 - playing an important role in LHC physics
- σ_Z, σ_W
 - large production cross sections
 - 1 to 10's of nbs





CMS Experiment at LHC, CERN
 Run 135149, Event 125426133
 Lumi section: 1345
 Sun May 09 2010, 05:24:09 CEST

Muon $p_T = 67.3, 50.6 \text{ GeV}/c$
 Inv. mass = $93.2 \text{ GeV}/c^2$

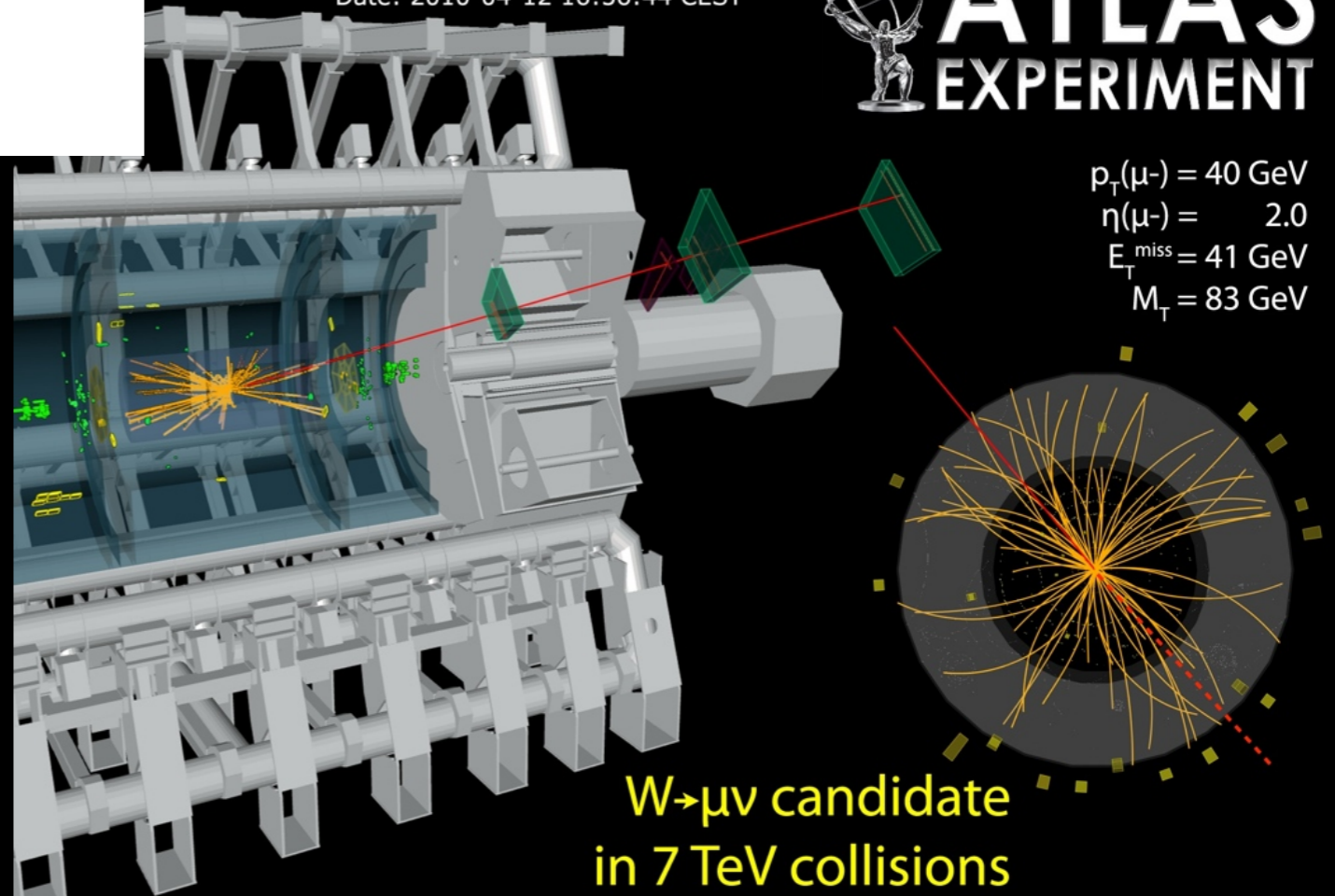


Z & W candidates

Run: 152845, Event: 3338173
 Date: 2010-04-12 16:56:44 CEST



$p_T(\mu^-) = 40 \text{ GeV}$
 $\eta(\mu^-) = 2.0$
 $E_T^{\text{miss}} = 41 \text{ GeV}$
 $M_T = 83 \text{ GeV}$

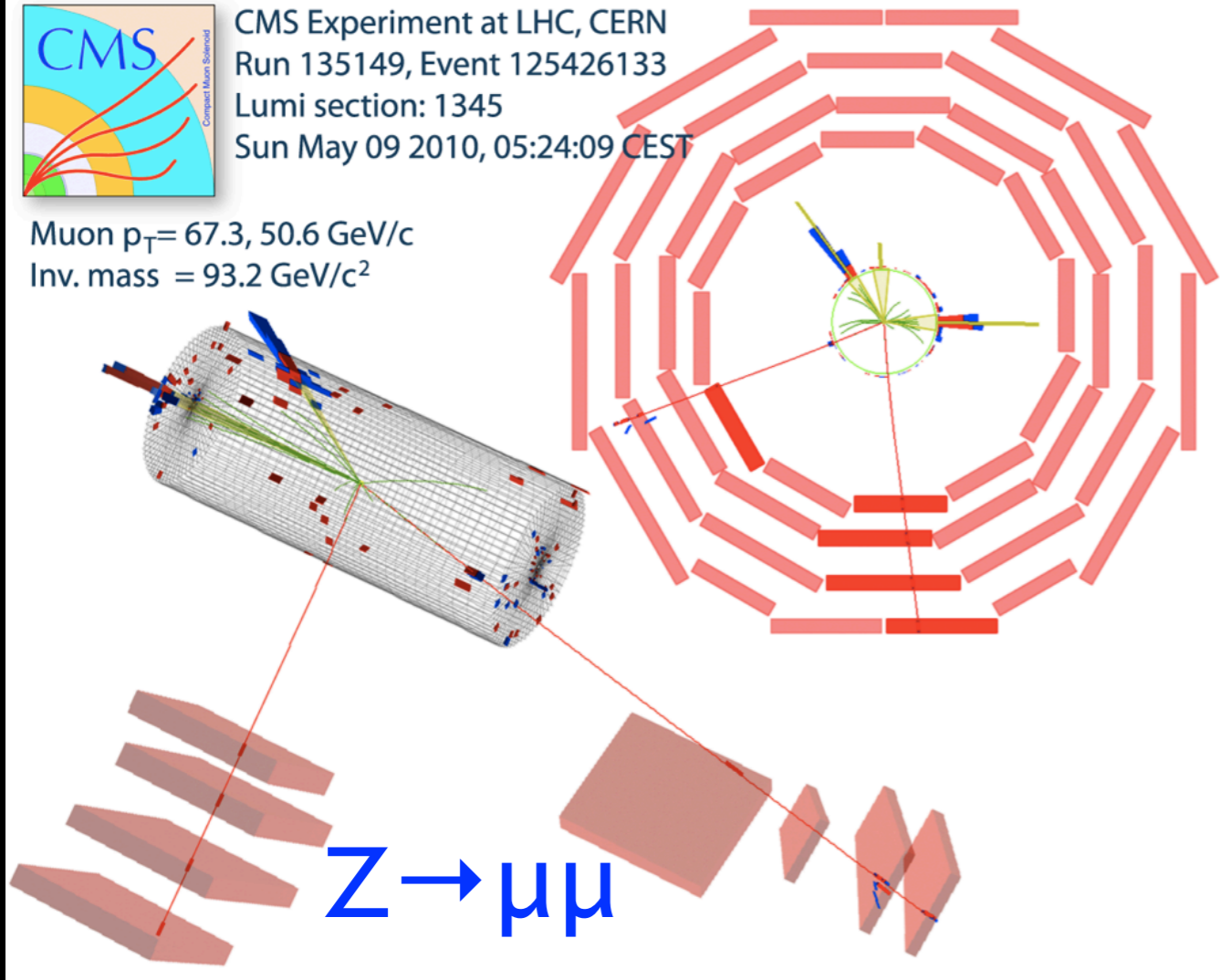


100k events/yr



CMS Experiment at LHC, CERN
 Run 135149, Event 125426133
 Lumi section: 1345
 Sun May 09 2010, 05:24:09 CEST

Muon $p_T = 67.3, 50.6 \text{ GeV}/c$
 Inv. mass = $93.2 \text{ GeV}/c^2$

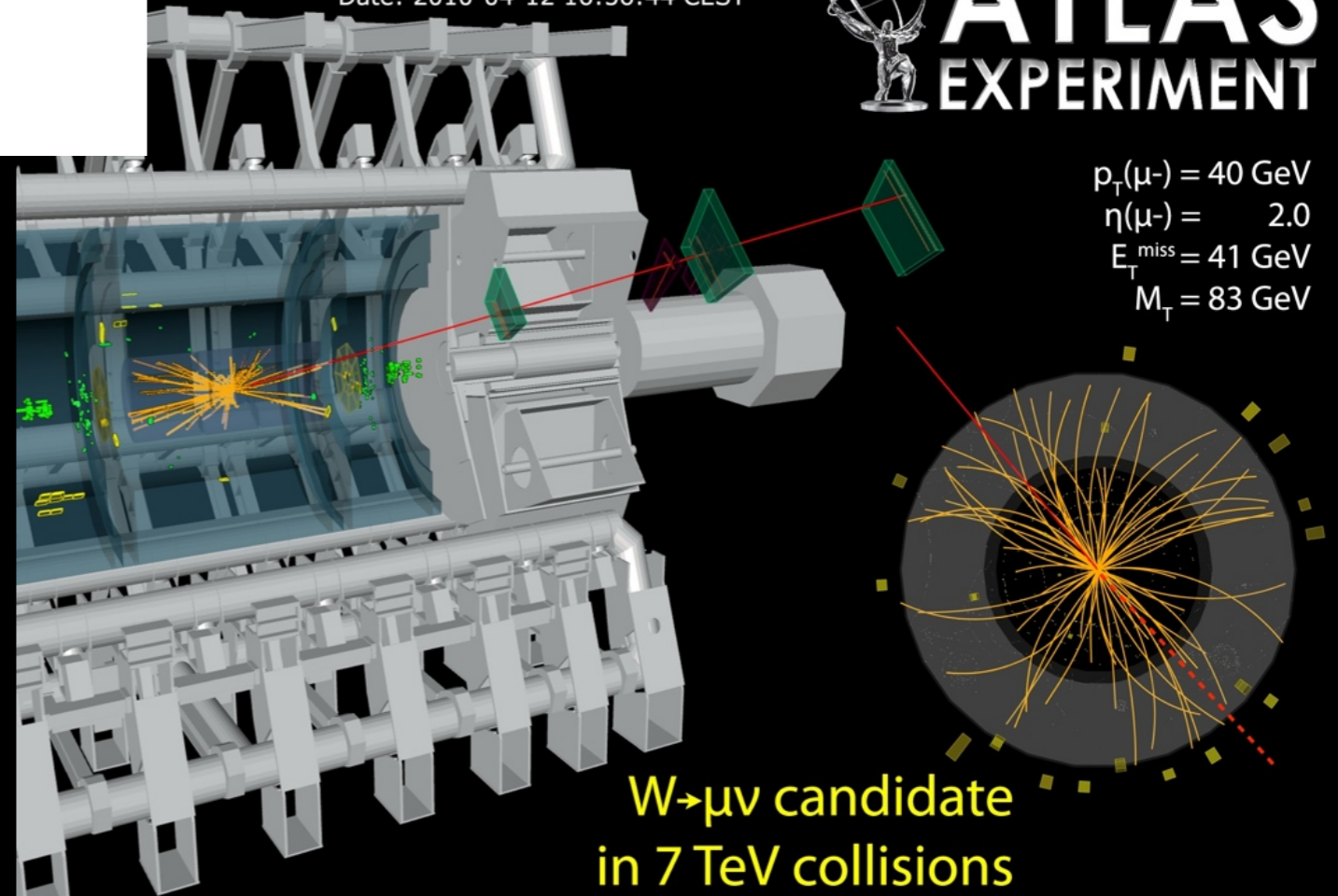


Z & W candidates

Run: 152845, Event: 3338173
 Date: 2010-04-12 16:56:44 CEST



$p_T(\mu^-) = 40 \text{ GeV}$
 $\eta(\mu^-) = 2.0$
 $E_T^{\text{miss}} = 41 \text{ GeV}$
 $M_T = 83 \text{ GeV}$



Over
 One Year Old!

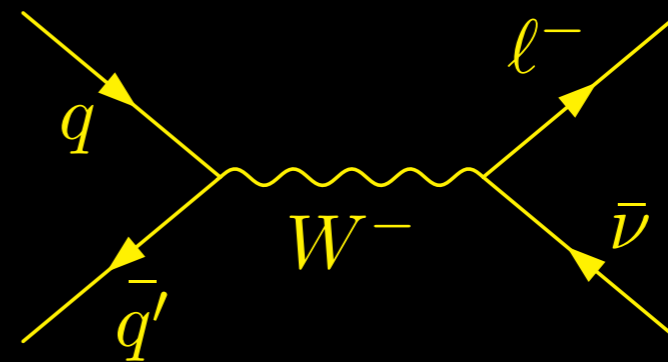
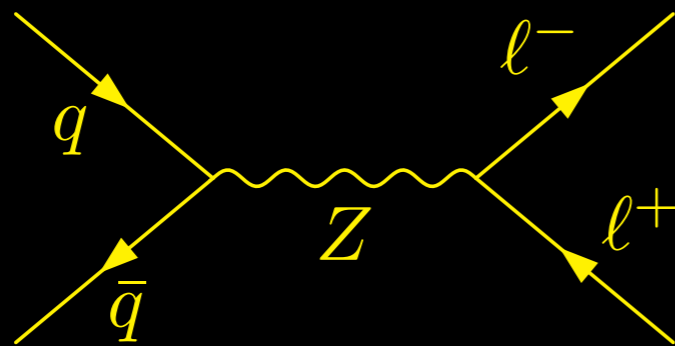
100k events/yr

Z & W's at Hadron Colliders

- Electroweak (EW) gauge boson production

- decay to leptonic final state

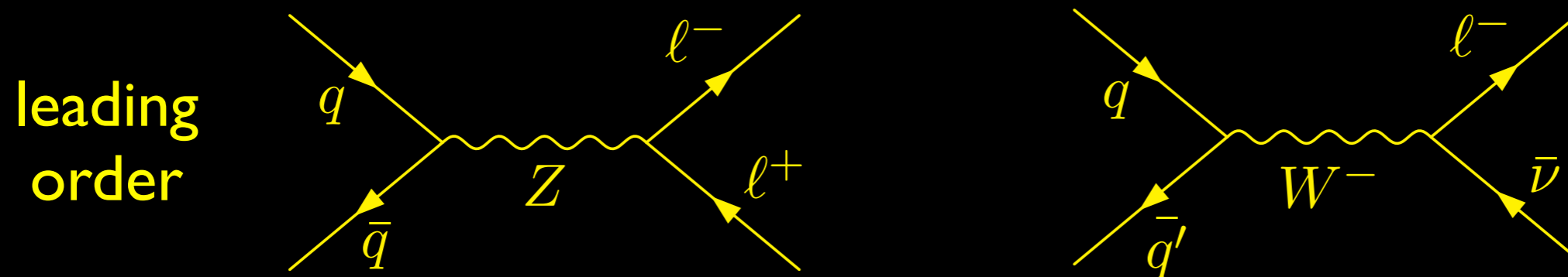
leading
order



Z & W's at Hadron Colliders

- Electroweak (EW) gauge boson production

- decay to leptonic final state



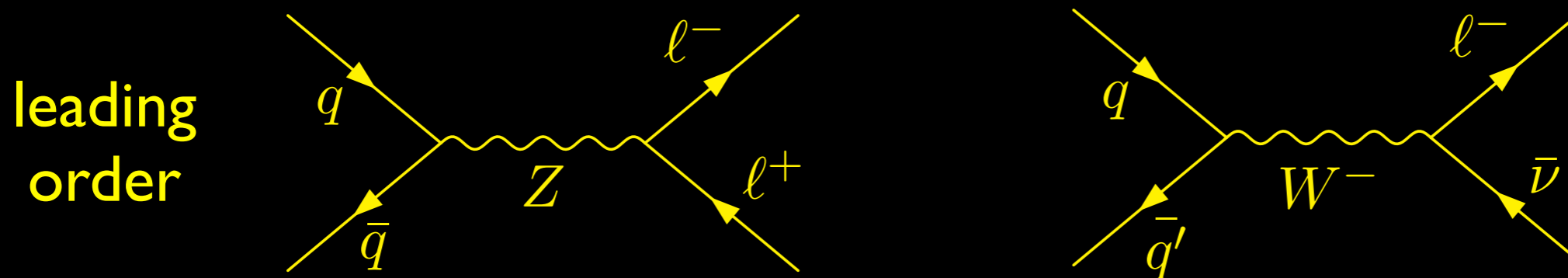
- Hadronic cross section

$$d\sigma_{AB} = \sum_{a,b} \int dx_a f_{a/A}(x_a, \mu_F) \int dx_b f_{b/B}(x_b, \mu_F) d\hat{\sigma}_{ab}$$

Z & W's at Hadron Colliders

- Electroweak (EW) gauge boson production

- decay to leptonic final state



- Hadronic cross section

$$d\sigma_{AB} = \sum_{a,b} \int dx_a \underline{f_{a/A}(x_a, \mu_F)} \int dx_b \underline{f_{b/B}(x_b, \mu_F)} d\hat{\sigma}_{ab}$$

parton distribution functions (PDF)

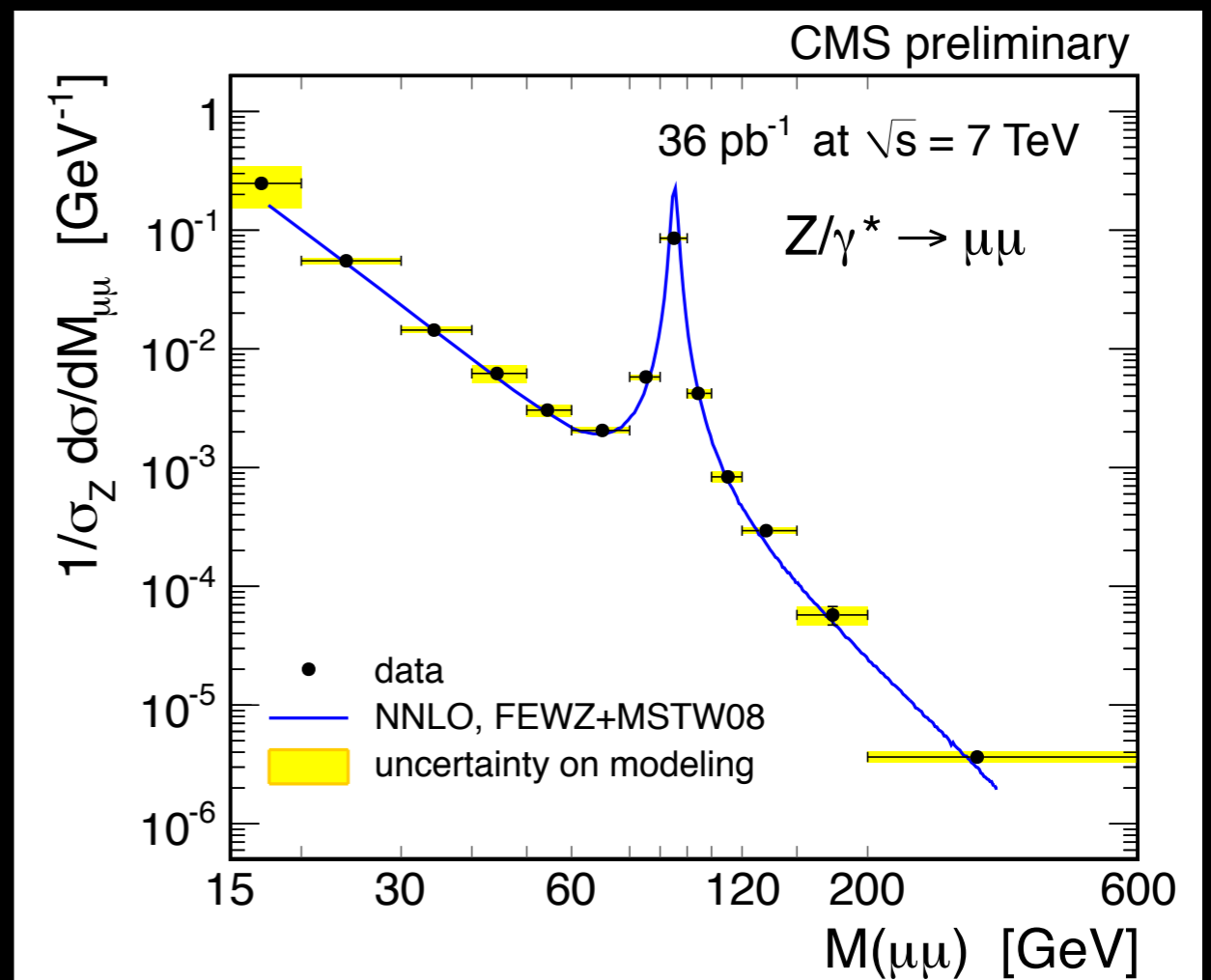
- probability to find a parton with momentum fraction x
- process independent
- μ_F factorization scale

Z & W's at Hadron Colliders

- Clear collider signature
 - leptonic final states

$$pp \rightarrow Z/\gamma^* \rightarrow l^- l^+$$

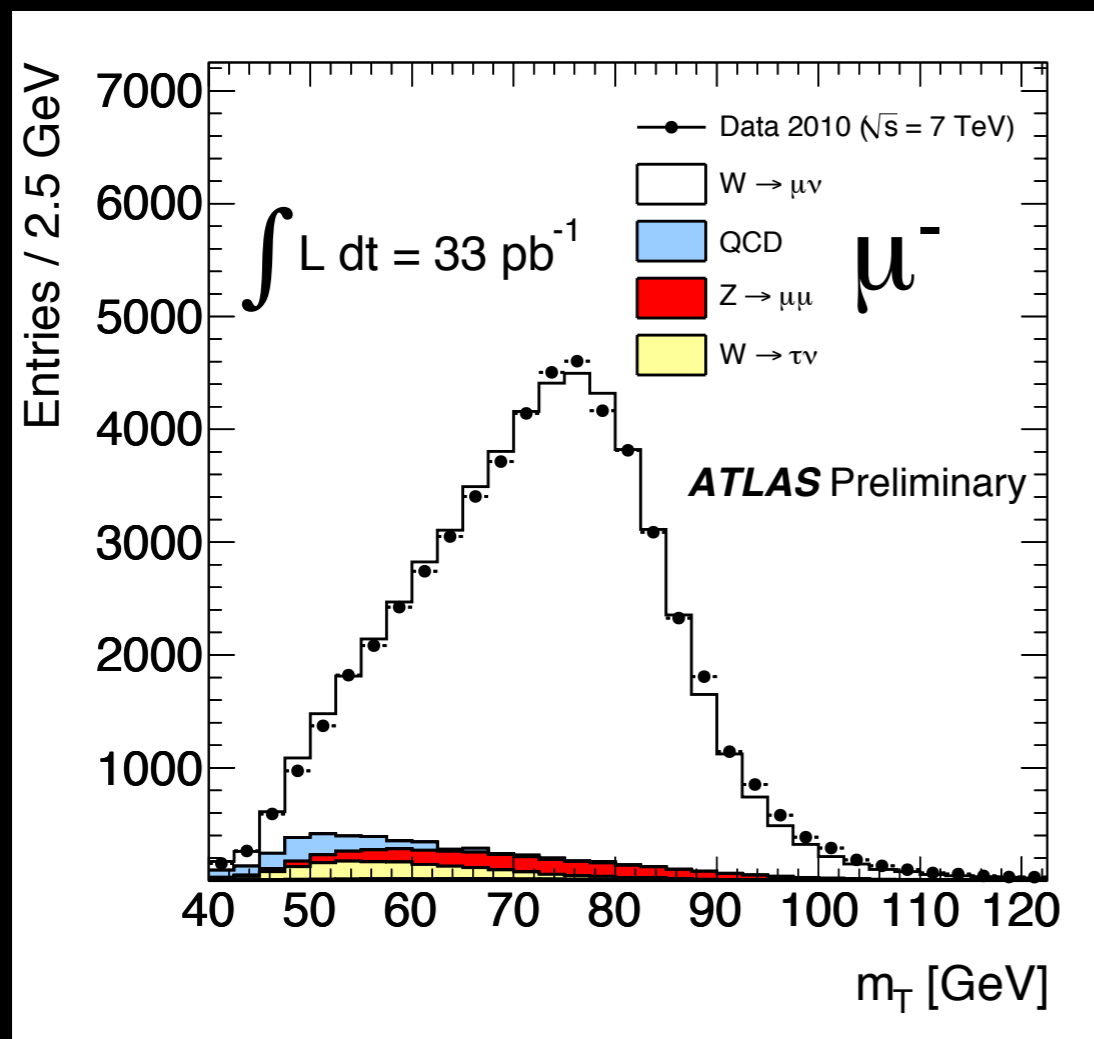
- two opposite sign candidates
- high p_T
- isolation & identification
- peak in $M_{l^+ l^-}$



CMS PAS EVK-10-007

Z & W's at Hadron Colliders

- Clear collider signature
 - leptonic final states



$$pp \rightarrow W^- \rightarrow l^- \bar{\nu}$$

- single charged candidate
- high lepton p_T
- lepton isolation & identification
- missing E_T
- peak in M_T

$$M_T = [2p_T^l p_T^\nu (1 - \cos \phi_{l\nu})]^{1/2}$$

ATLAS-CONF-2011-041

Z & W's at Hadron Colliders

- Make for good analyses on early data
 - integrated luminosity of nb^{-1} 's to pb^{-1} 's
- Rediscovering the EW gauge bosons at LHC

Z & W's at Hadron Colliders

- Make for good analyses on early data
 - integrated luminosity of nb^{-1} 's to pb^{-1} 's
- Rediscovering the EW gauge bosons at LHC
- Z/W properties are well known - reliable predictions
 - $M_Z, M_W, \Gamma_Z, \Gamma_W$, cross sections

Z & W's at Hadron Colliders

- Make for good analyses on early data
 - integrated luminosity of nb^{-1} 's to pb^{-1} 's
- Rediscovering the EW gauge bosons at LHC
- Z/W properties are well known - reliable predictions
 - $M_Z, M_W, \Gamma_Z, \Gamma_W$, cross sections
- LHC 'Standard Candles'



Drell-Yan as Standard Candles

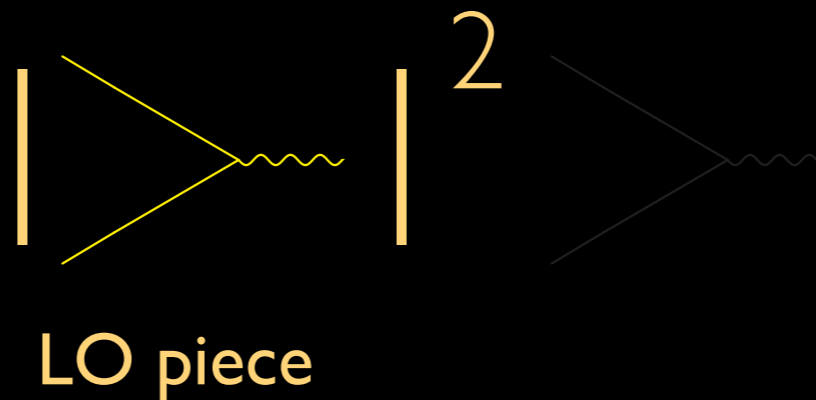
- Detector calibration & performance
 - M_Z, Γ_M from M_{ll} distributions at Z resonance
- Luminosity monitoring (Dittmar et. al.)
 - σ_W, σ_Z
- Important for measuring EW parameter
 - precise measurements of $M_W, \Gamma_W, \sin^2\theta_{\text{eff}}$
- Look at new analysis tools
 - $a_\tau, \phi_\eta^* \rightarrow$ low Z p_τ study
 - reduce uncertainty due to energy resolution

Drell-Yan as Standard Candles

- Study of perturbative QCD
 - p_T distributions
 - Z $p_T > 0$ starting at NLO

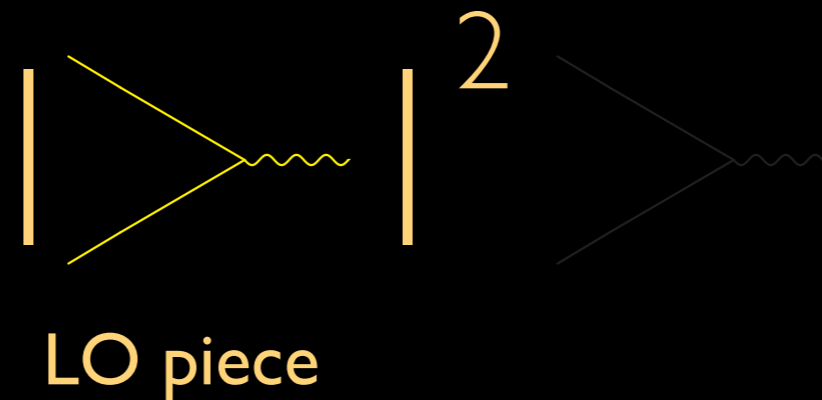
Drell-Yan as Standard Candles

- Study of perturbative QCD
 - p_T distributions
 - Z $p_T > 0$ starting at NLO



Drell-Yan as Standard Candles

- Study of perturbative QCD
 - p_T distributions
 - Z $p_T > 0$ starting at NLO



$$2 \operatorname{Re} \left(\text{tree} \times \text{NLO} \right) + \left| \text{NLO} \right|^2$$

NLO contributions

Drell-Yan as Standard Candles

- PDF measurements
 - Z rapidity distributions

Drell-Yan as Standard Candles

- PDF measurements
 - Z rapidity distributions

$$d\sigma_{AB} = \sum_{a,b} \int dx_a f_{a/A}(x_a) \int dx_b f_{b/B}(x_b) d\hat{\sigma}_{ab}$$

Drell-Yan as Standard Candles

- PDF measurements
 - Z rapidity distributions

$$d\sigma_{AB} = \sum_{a,b} \int dx_a f_{a/A}(x_a) \int dx_b f_{b/B}(x_b) d\hat{\sigma}_{ab}$$

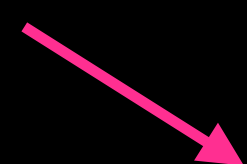
$$x_a = \frac{Q}{\sqrt{s}} e^Y \quad \& \quad x_b = \frac{Q}{\sqrt{s}} e^{-Y}$$

Drell-Yan as Standard Candles

- PDF measurements
 - Z rapidity distributions

$$d\sigma_{AB} = \sum_{a,b} \int dx_a f_{a/A}(x_a) \int dx_b f_{b/B}(x_b) d\hat{\sigma}_{ab}$$

$$x_a = \frac{Q}{\sqrt{s}} e^Y \quad \& \quad x_b = \frac{Q}{\sqrt{s}} e^{-Y}$$


$$Y = \frac{1}{2} \ln\left(\frac{x_a}{x_b}\right)$$

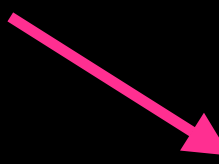
Drell-Yan as Standard Candles

- PDF measurements
 - Z rapidity distributions

$$d\sigma_{AB} = \sum_{a,b} \int dx_a f_{a/A}(x_a) \int dx_b f_{b/B}(x_b) d\hat{\sigma}_{ab}$$

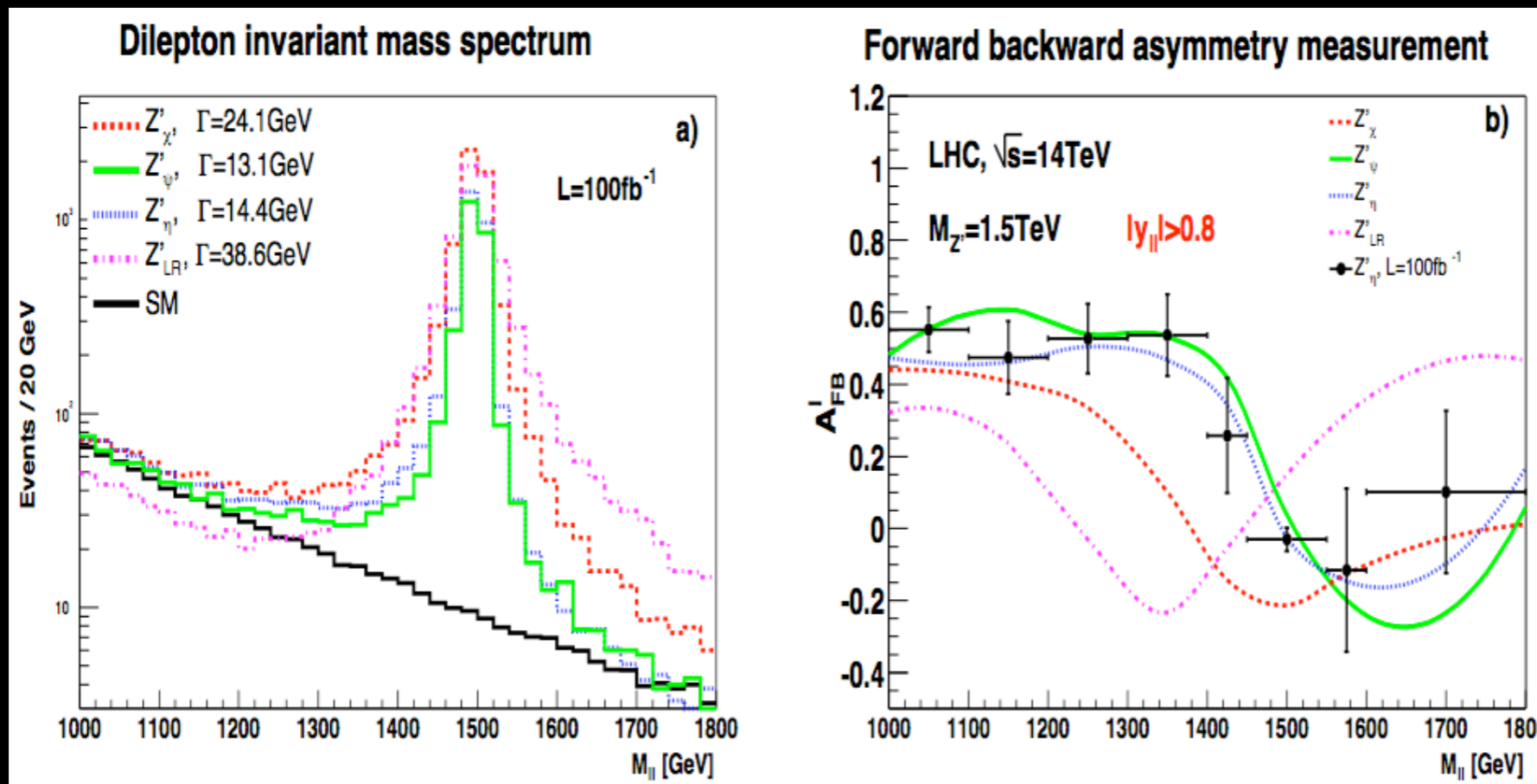
$$x_a = \frac{Q}{\sqrt{s}} e^Y \quad \& \quad x_b = \frac{Q}{\sqrt{s}} e^{-Y}$$

- Distributions in rapidity
 - measure/constrain PDFs


$$Y = \frac{1}{2} \ln\left(\frac{x_a}{x_b}\right)$$

Drell-Yan as Standard Candles

- Constraint on some BSM physics
 - deviations in M_{ll} distributions (e.g. Z')
 - deviations in forward/backward asymmetry, A_{FB}



Dittmar, Nicollrat, Djouadi
 arXiv:hep-ph/0307020

Need for Higher Order?

Need for Higher Order?

- LHC will produce large amount of data
 - small statistical error
 - measurements limited by systematics & theoretical error

Need for Higher Order?

- LHC will produce large amount of data
 - small statistical error
 - measurements limited by systematics & theoretical error
- As LHC continues to run
 - expect systematics to improve

Need for Higher Order?

- LHC will produce large amount of data
 - small statistical error
 - measurements limited by systematics & theoretical error
- As LHC continues to run
 - expect systematics to improve
- Can expect percent level physics
 - need to reduce theoretical uncertainties

Need for Higher Order?

- Measurements require theory input
- Acceptances derived from calculation and simulation
- Precision physics → need higher order calculations

Need for Higher Order?

- Measurements require theory input

$$\sigma = \frac{N_{sig}}{\epsilon \cdot A \cdot \mathcal{L}_{int}}$$

- Acceptances derived from calculation and simulation
- Precision physics → need higher order calculations

Need for Higher Order?

- Measurements require theory input

$$\sigma = \frac{N_{sig}}{\epsilon \cdot A \cdot \mathcal{L}_{int}}$$

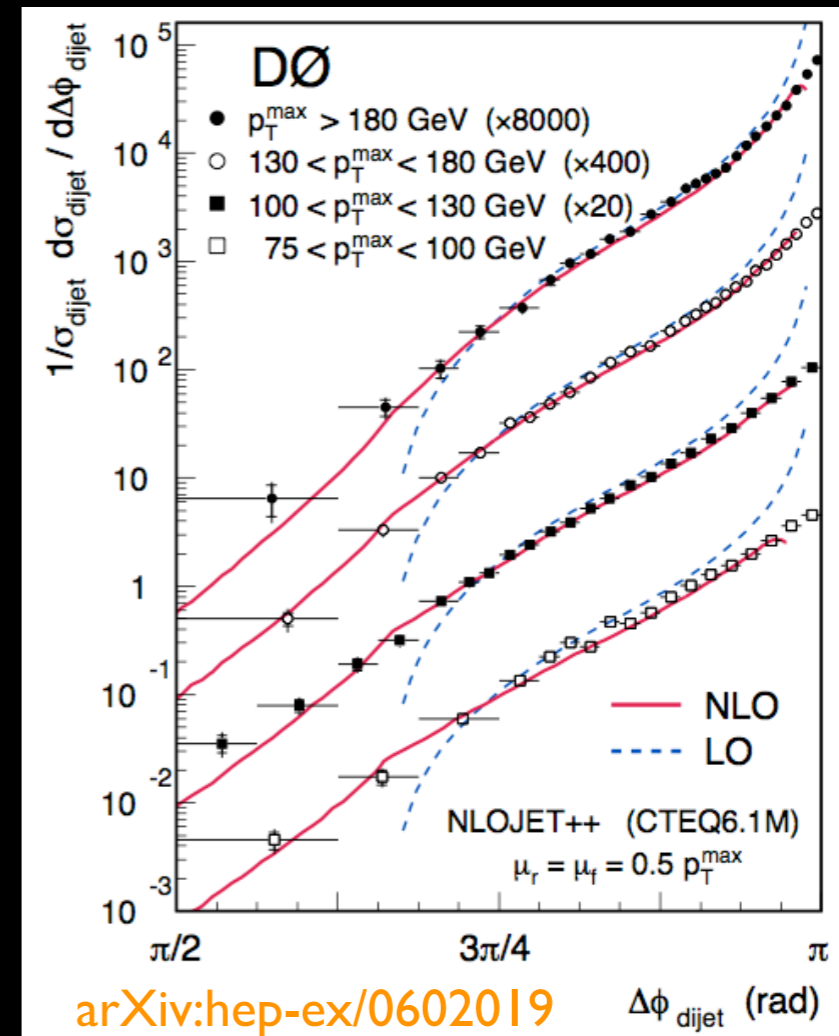
- Acceptances derived from calculation and simulation
- Precision physics → need higher order calculations

Higher Orders

- Quantitative predictive power begins at NLO
- Improves normalization
- Reduces uncertainties
- More accurately describes distribution shapes

Higher Orders

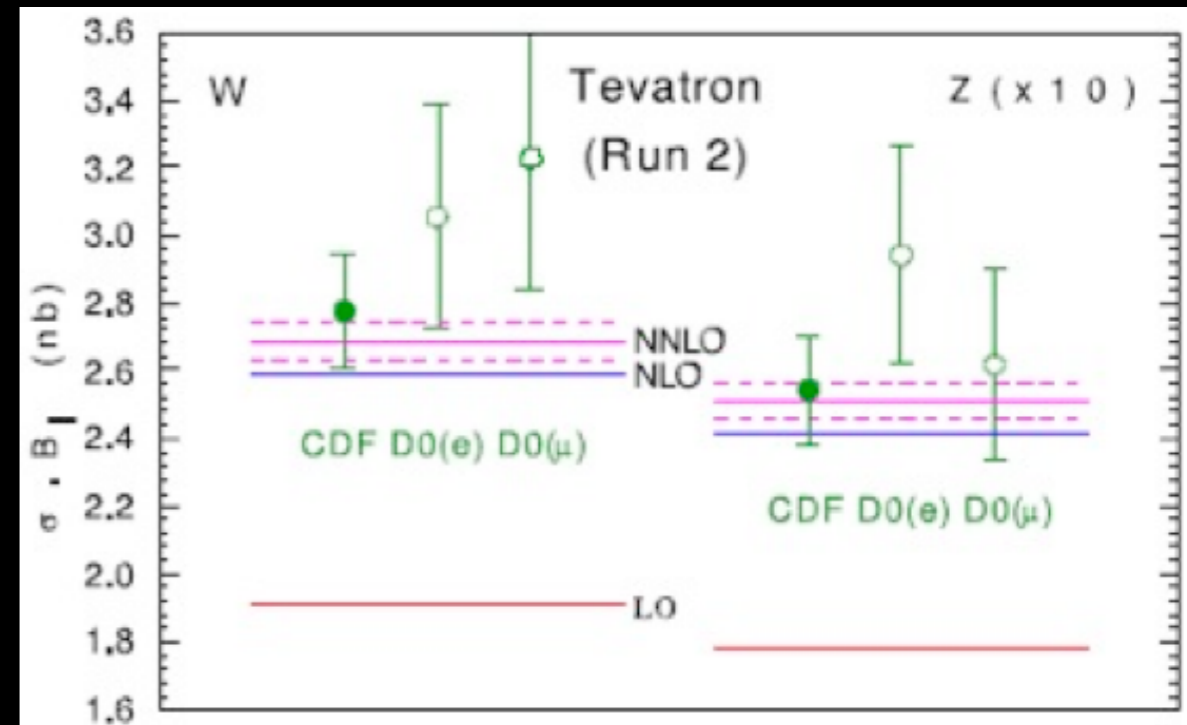
- Quantitative predictive power begins at NLO
- Improves normalization
- Reduces uncertainties
- More accurately describes distribution shapes



↑
dijet production

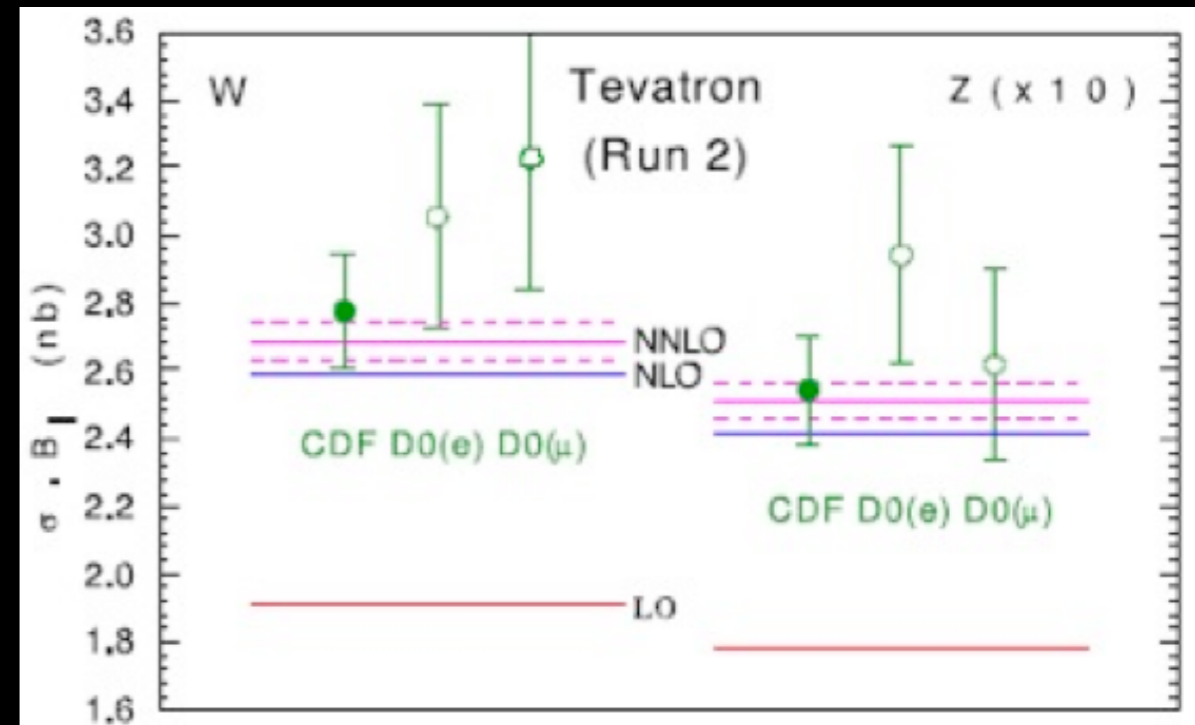
Higher Orders

- Quantitative predictive power begins at NLO
- Improves normalization
- Reduces uncertainties
- More accurately describes distribution shapes



Higher Orders

- Quantitative predictive power begins at NLO
- Improves normalization
- Reduces uncertainties
- More accurately describes distribution shapes



number of jets	CDF	LO	NLO
1	53.5 ± 5.6	$41.40(0.02)^{+7.59}_{-5.94}$	$57.83(0.12)^{+4.36}_{-4.00}$
2	6.8 ± 1.1	$6.159(0.004)^{+2.41}_{-1.58}$	$7.62(0.04)^{+0.62}_{-0.86}$
3	0.84 ± 0.24	$0.796(0.001)^{+0.488}_{-0.276}$	$0.882(0.005)^{+0.057}_{-0.138}$

← inclusive cross section for W + n jets (pb)

Berger, et al.
arXiv:0907.1984

DY + Jets Production

- Z/W boson production purely EW $\sigma_{DY} \propto \alpha_{em}^2$

DY + Jets Production

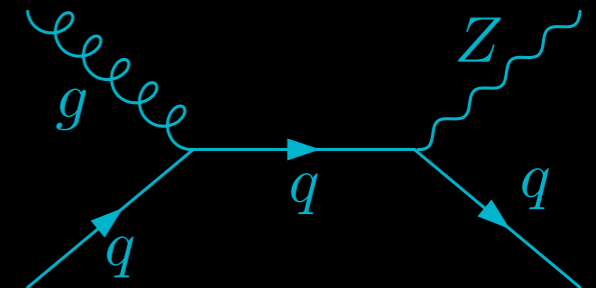
- Z/W boson production purely EW $\sigma_{DY} \propto \alpha_{em}^2$
- Introduce hard jets associated with DY
 - introduce strong coupling at tree level $\sigma_{DY+n jets} \propto \alpha_{em}^2 \alpha_s^n$

DY + Jets Production

- Z/W boson production purely EW $\sigma_{DY} \propto \alpha_{em}^2$
- Introduce hard jets associated with DY
 - introduce strong coupling at tree level $\sigma_{DY+n\text{ jets}} \propto \alpha_{em}^2 \alpha_s^n$
- Emitted parton can be collinear or soft
 - well defined jet has a p_T cutoff - regulates IR divergences

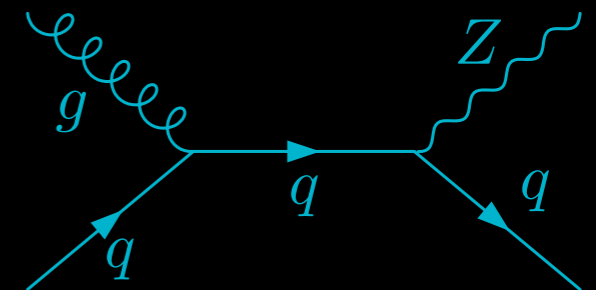
DY + Jets Production

- Z/W boson production purely EW $\sigma_{DY} \propto \alpha_{em}^2$
- Introduce hard jets associated with DY
 - introduce strong coupling at tree level $\sigma_{DY+n\text{ jets}} \propto \alpha_{em}^2 \alpha_s^n$
- Emitted parton can be collinear or soft
 - well defined jet has a p_T cutoff - regulates IR divergences
- Z/W + 1 jet helps constrain gluon PDF



DY + Jets Production

- Z/W boson production purely EW $\sigma_{DY} \propto \alpha_{em}^2$
- Introduce hard jets associated with DY
 - introduce strong coupling at tree level $\sigma_{DY+n\text{ jets}} \propto \alpha_{em}^2 \alpha_s^n$
- Emitted parton can be collinear or soft
 - well defined jet has a p_T cutoff - regulates IR divergences
- Z/W + 1 jet helps constrain gluon PDF
- Background to new physics searches



Z+Jet LO Process

- Focus on Z+jet production

Z+Jet LO Process

- Focus on Z+jet production
- LO process $\propto \alpha_{em}^2 \alpha_s$

Z+Jet LO Process

- Focus on Z+jet production
- LO process $\propto \alpha_{em}^2 \alpha_s$

3 partonic processes

$$q \bar{q} \rightarrow Z/\gamma^* g \rightarrow \ell^- \ell^+ g$$

$$g q \rightarrow Z/\gamma^* q \rightarrow \ell^- \ell^+ q$$

$$g \bar{q} \rightarrow Z/\gamma^* \bar{q} \rightarrow \ell^- \ell^+ \bar{q}$$

Z+Jet LO Process

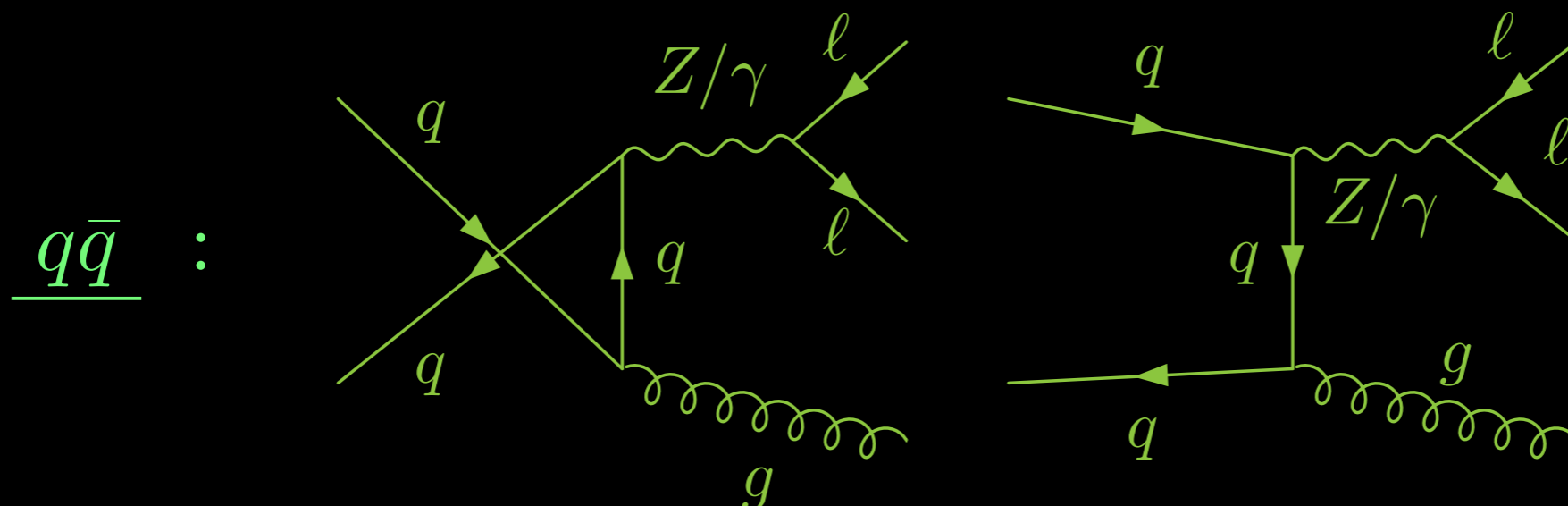
- Focus on Z+jet production
- LO process $\propto \alpha_{em}^2 \alpha_s$

3 partonic processes

$$q \bar{q} \rightarrow Z/\gamma^* g \rightarrow \ell^- \ell^+ g$$

$$g q \rightarrow Z/\gamma^* q \rightarrow \ell^- \ell^+ q$$

$$g \bar{q} \rightarrow Z/\gamma^* \bar{q} \rightarrow \ell^- \ell^+ \bar{q}$$



Z+Jet LO Process

- Focus on Z+jet production
- LO process $\propto \alpha_{em}^2 \alpha_s$

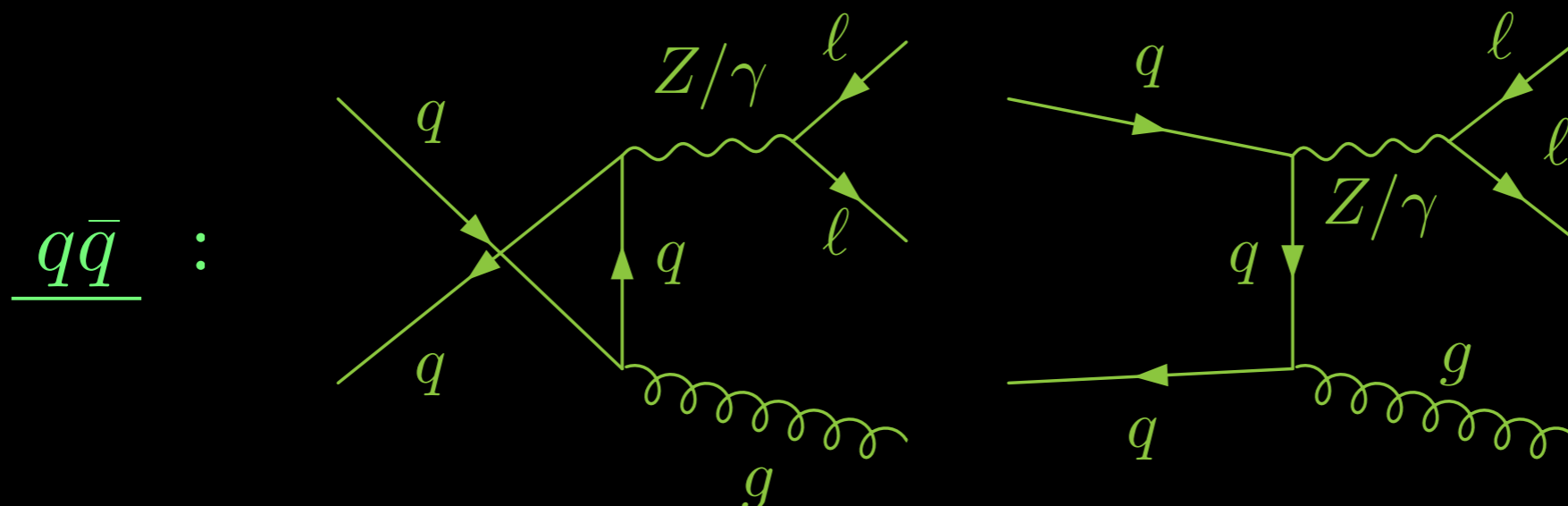
3 partonic processes

$$q \bar{q} \rightarrow Z/\gamma^* g \rightarrow \ell^- \ell^+ g$$

$$g q \rightarrow Z/\gamma^* q \rightarrow \ell^- \ell^+ q$$

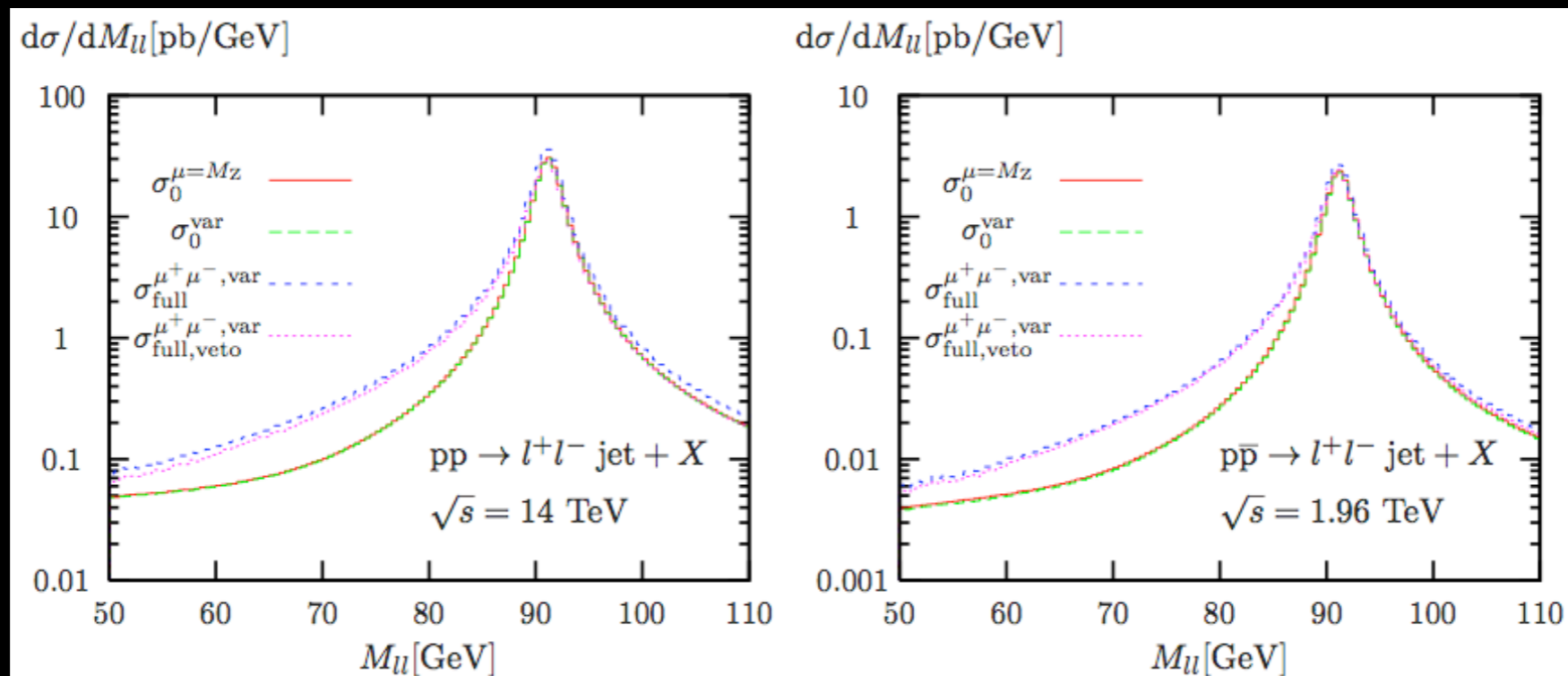
$$g \bar{q} \rightarrow Z/\gamma^* \bar{q} \rightarrow \ell^- \ell^+ \bar{q}$$

diagrams obtained by quark and antiquark crossing with gluon



Z+Jet at NLO

- Quantitative predictive power begins at NLO
 - Studies in EW & QCD NLO corrections Denner, Dittmaier, Kasprzik, Muck - arXiv:1103.0914
 - QCD: $\sigma_{\text{int}} \sim 35\%$
 - EW: $\sigma_{\text{int}} \sim -3$ to -5%
- } variation in differential distributions



New Physics Contributions

- Z+jet a critical process at hadron colliders

New Physics Contributions

- Z+jet a critical process at hadron colliders
- New physics can enter SM processes in the form of higher order corrections
 - how does new physics effect this standard candle?

New Physics Contributions

- Z+jet a critical process at hadron colliders
- New physics can enter SM processes in the form of higher order corrections
 - how does new physics effect this standard candle?
- Consider Z+jet in the supersymmetric framework
 - calculate supersymmetric QCD corrections to Z+jet

New Physics Contributions

- Z+jet a critical process at hadron colliders
- New physics can enter SM processes in the form of higher order corrections
 - how does new physics effect this standard candle?
- Consider Z+jet in the supersymmetric framework
 - calculate supersymmetric QCD corrections to Z+jet
- Goal: Investigate stability of Z+jet as a standard candle under sQCD corrections

Supersymmetry

- Symmetry relating fermions and bosons
- Solution to hierarchy problem
 - SM quadratic corrections to Higgs mass cancel against corrections from supersymmetric partners
- Gauge coupling unification (weak scale soft SUSY breaking)
- Dark matter candidate with R -parity conservation

$$P_R = (-1)^{3(B-L)+2s}$$

B - baryon #
L - lepton #
s - spin

SM particle: $P_R = 1$ SUSY particle: $P_R = -1$

The MSSM

- Minimal Supersymmetric extension of the Standard Model
 - only 1 superpartner for each SM particle
 - anomaly free
 - supersymmetry is softly broken
 - R -parity conserved

MSSM Particle Content

chiral superfields

Names		spin 0	spin 1/2	$SU(3)_C, SU(2)_L, U(1)_Y$
squarks, quarks ($\times 3$ families)	Q	$(\tilde{u}_L, \tilde{d}_L)$	(u_L, d_L)	$(\mathbf{3}, \mathbf{2}, +\frac{1}{6})$
	\bar{u}	\tilde{u}_R^*	u_R^\dagger	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})$
	\bar{d}	\tilde{d}_R^*	d_R^\dagger	$(\bar{\mathbf{3}}, \mathbf{1}, +\frac{1}{3})$
sleptons, leptons ($\times 3$ families)	L	$(\tilde{\nu}, \tilde{e}_L)$	(ν, e_L)	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$
	\bar{e}	\tilde{e}_R^*	e_R^\dagger	$(\bar{\mathbf{1}}, \mathbf{1}, +1)$
Higgs, higgsinos	H_u	(H_u^+, H_u^0)	$(\tilde{H}_u^+, \tilde{H}_u^0)$	$(\mathbf{1}, \mathbf{2}, +\frac{1}{2})$
	H_d	(H_d^0, H_d^-)	$(\tilde{H}_d^0, \tilde{H}_d^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$

vector superfields

Names	spin 1/2	spin 1	$SU(3)_C, SU(2)_L, U(1)_Y$
gluino, gluon	\tilde{g}	g	$(\mathbf{8}, \mathbf{1}, 0)$
Winos, W bosons	$\tilde{W}^\pm, \tilde{W}^0$	W^\pm, W^0	$(\mathbf{1}, \mathbf{3}, 0)$
Binos, B boson	\tilde{B}^0	B^0	$(\mathbf{1}, \mathbf{1}, 0)$

MSSM Particle Content

chiral superfields

Names		spin 0	spin 1/2	$SU(3)_C, SU(2)_L, U(1)_Y$
squarks, quarks ($\times 3$ families)	Q	$(\tilde{u}_L, \tilde{d}_L)$	(u_L, d_L)	$(\mathbf{3}, \mathbf{2}, +\frac{1}{6})$
	\bar{u}	\tilde{u}_R^*	u_R^\dagger	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})$
	\bar{d}	\tilde{d}_R^*	d_R^\dagger	$(\bar{\mathbf{3}}, \mathbf{1}, +\frac{1}{3})$
sleptons, leptons ($\times 3$ families)	L	$(\tilde{\nu}, \tilde{e}_L)$	(ν, e_L)	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$
	\bar{e}	\tilde{e}_R^*	e_R^\dagger	$(\bar{\mathbf{1}}, \mathbf{1}, +1)$
Higgs, higgsinos	H_u	(H_u^+, H_u^0)	$(\tilde{H}_u^+, \tilde{H}_u^0)$	$(\mathbf{1}, \mathbf{2}, +\frac{1}{2})$
	H_d	(H_d^0, H_d^-)	$(\tilde{H}_d^0, \tilde{H}_d^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$

interested in
QCD-like
corrections to
Z+l jet

what's charged
under $SU(3)_C$?

vector superfields

Names	spin 1/2	spin 1	$SU(3)_C, SU(2)_L, U(1)_Y$
gluino, gluon	\tilde{g}	g	$(\mathbf{8}, \mathbf{1}, 0)$
Winos, W bosons	$\tilde{W}^\pm, \tilde{W}^0$	W^\pm, W^0	$(\mathbf{1}, \mathbf{3}, 0)$
Binos, B boson	\tilde{B}^0	B^0	$(\mathbf{1}, \mathbf{1}, 0)$

MSSM Particle Content

chiral superfields

Names		spin 0	spin 1/2	$SU(3)_C, SU(2)_L, U(1)_Y$
squarks, quarks ($\times 3$ families)	Q	$(\tilde{u}_L, \tilde{d}_L)$	(u_L, d_L)	$(\mathbf{3}, \mathbf{2}, +\frac{1}{6})$
	\bar{u}	\tilde{u}_R^*	u_R^\dagger	$(\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})$
	\bar{d}	\tilde{d}_R^*	d_R^\dagger	$(\bar{\mathbf{3}}, \mathbf{1}, +\frac{1}{3})$
sleptons, leptons ($\times 3$ families)	L	$(\tilde{\nu}, \tilde{e}_L)$	(ν, e_L)	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$
	\bar{e}	\tilde{e}_R^*	e_R^\dagger	$(\bar{\mathbf{1}}, \mathbf{1}, +1)$
Higgs, higgsinos	H_u	(H_u^+, H_u^0)	$(\tilde{H}_u^+, \tilde{H}_u^0)$	$(\mathbf{1}, \mathbf{2}, +\frac{1}{2})$
	H_d	(H_d^0, H_d^-)	$(\tilde{H}_d^0, \tilde{H}_d^-)$	$(\mathbf{1}, \mathbf{2}, -\frac{1}{2})$

interested in
QCD-like
corrections to
Z+l jet

what's charged
under $SU(3)_C$?

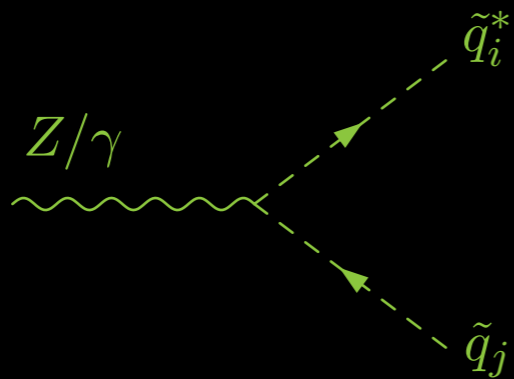
vector superfields

Names	spin 1/2	spin 1	$SU(3)_C, SU(2)_L, U(1)_Y$
gluino, gluon	\tilde{g}	g	$(\mathbf{8}, \mathbf{1}, 0)$
Winos, W bosons	$\tilde{W}^\pm, \tilde{W}^0$	W^\pm, W^0	$(\mathbf{1}, \mathbf{3}, 0)$
Binos, B boson	\tilde{B}^0	B^0	$(\mathbf{1}, \mathbf{1}, 0)$

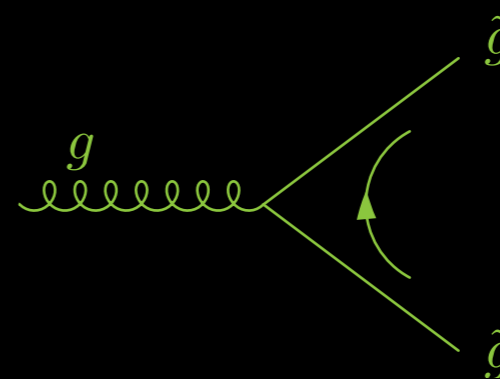
sQCD NLO Calculation

- supersymmetric QCD corrections
 - squarks, \tilde{q}_i & gluinos, \tilde{g}
- SM / SUSY interactions

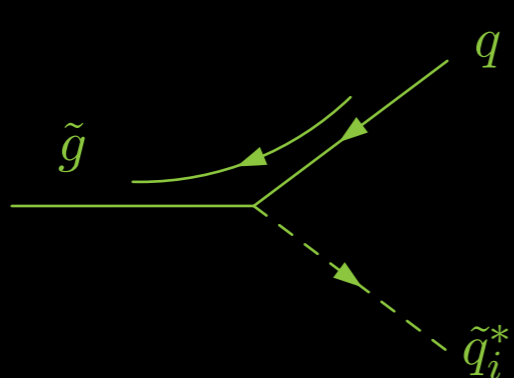
$Z \tilde{q} \tilde{q}$



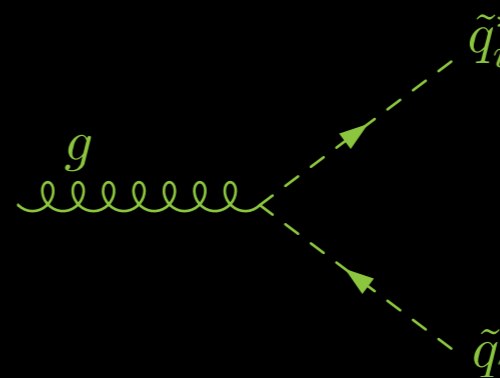
$g \tilde{g} \tilde{g}$



$q \tilde{q} \tilde{g}$



$g \tilde{q} \tilde{q}$



sQCD NLO Calculation

- Conventional diagrammatic approach used
- Two independent calculations performed

sQCD NLO Calculation

- Conventional diagrammatic approach used
- Two independent calculations performed

- FeynArts, FeynCalc, LoopTools

hep-ph/0012260, hep-ph/0105349,
hep-ph/9807565

sQCD NLO Calculation

- Conventional diagrammatic approach used
- Two independent calculations performed

- FeynArts, FeynCalc, LoopTools

hep-ph/0012260, hep-ph/0105349,
hep-ph/9807565

- QGRAF *P. Nogueira*

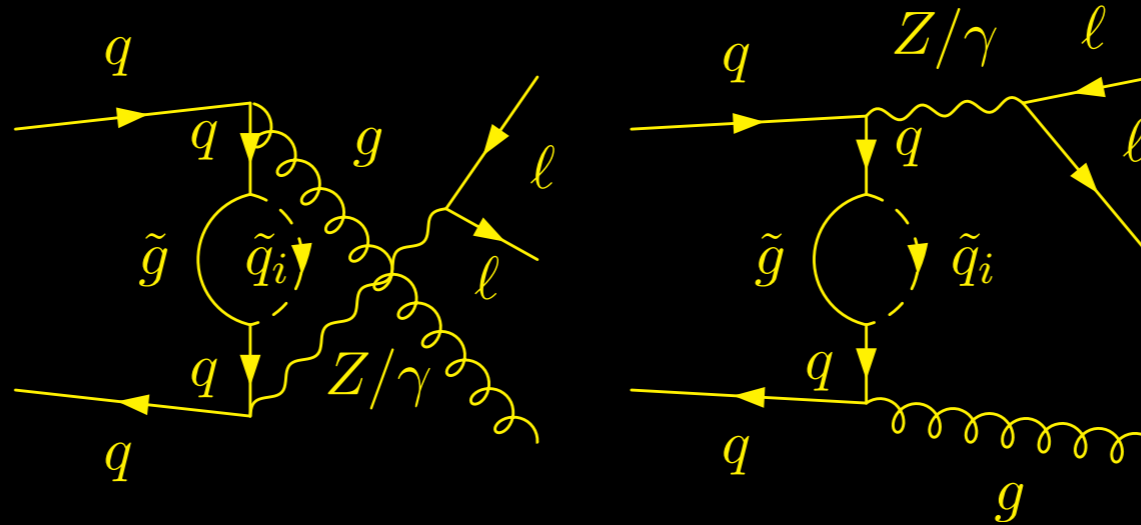
- in-house software using Form & Maple

- loops dealt with using AIR & QCDDLoop library

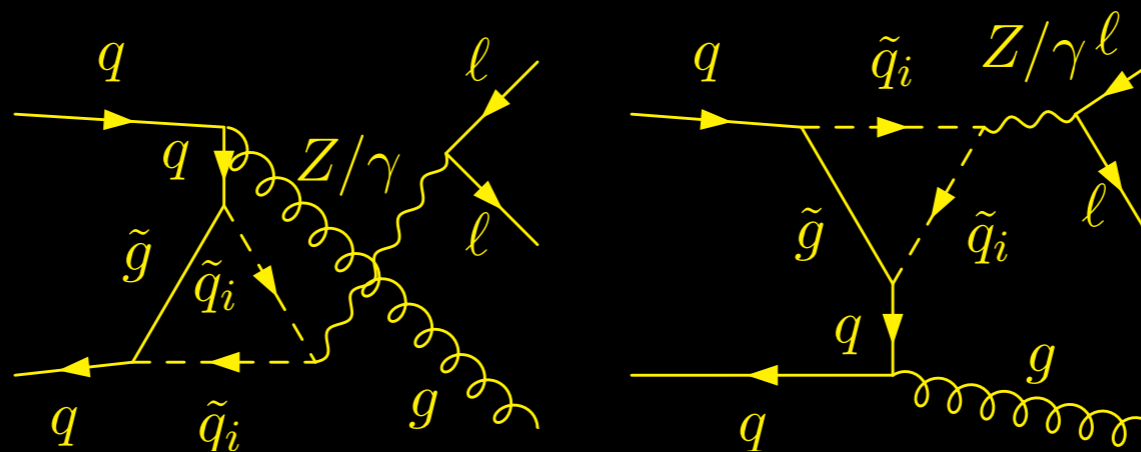
hep-ph/0404258, arXiv:0712.1851

Loop Diagrams

Self-energy Insertions

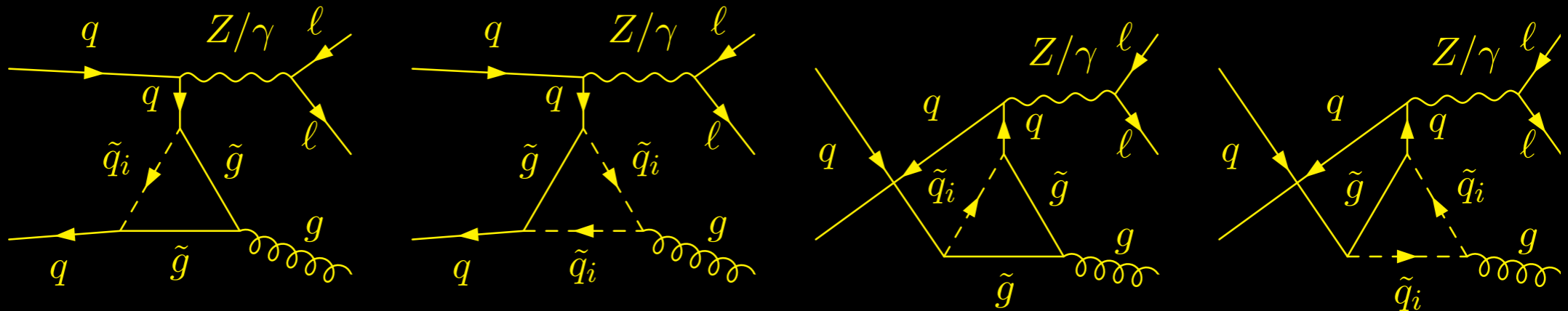


Z-Vertex Corrections

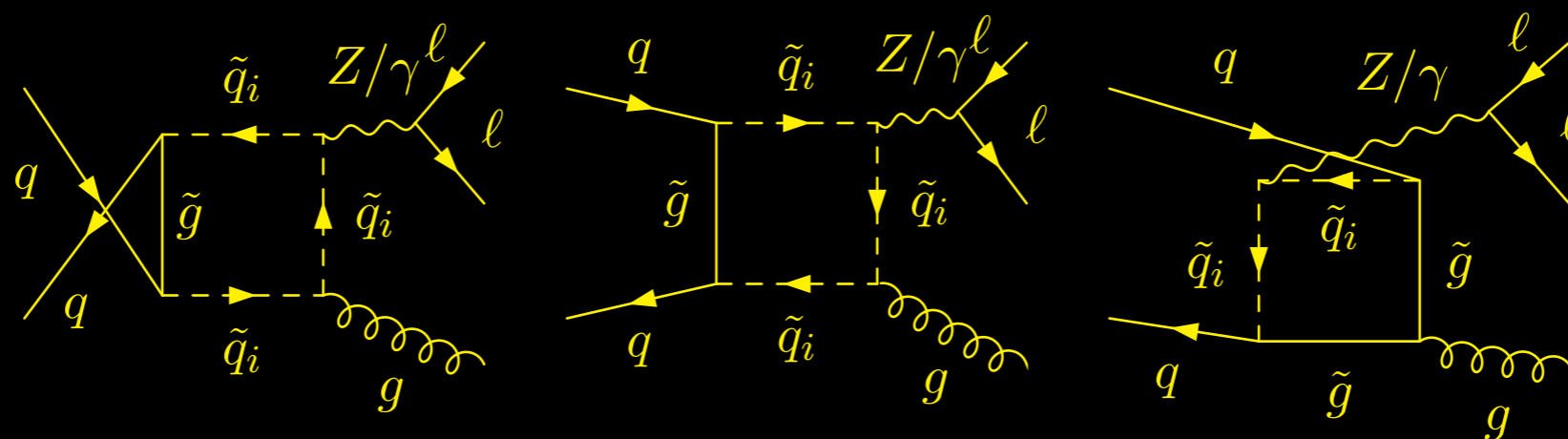


Loop Diagrams

Gluon-Vertex Corrections



Box Corrections

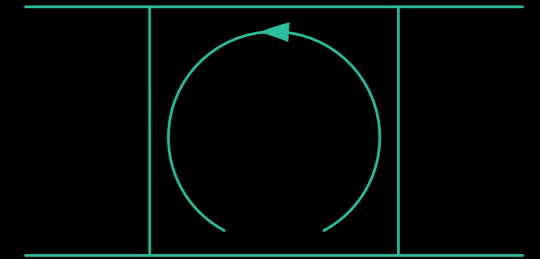


Loop Integrals

- Tensor and scalar integrals present in calculation
 - ▬ reduce tensor integrals into scalar integrals
 - Automated Integral Reduction (AIR) [hep-ph/0404258](#)
- All loop integrals are reduced to a set of master scalar integrals
- QCDloop used in numerical evaluation of scalar loop integrals [arXiv:0712.1851](#)

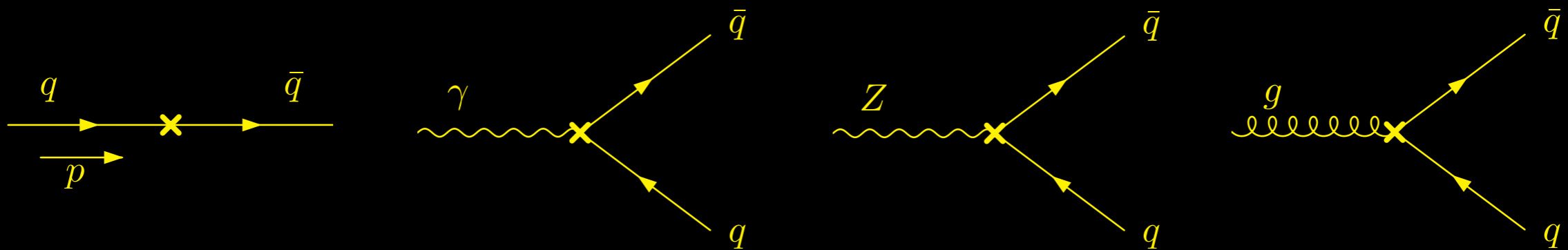
Singularities

- Massive particles propagating in the loop
 - no soft or collinear (IR) singularities
 - no real emission contribution
- Ultraviolet (UV) divergences do emerge from loop integrals
 - isolate UV singular structure using dimensional regularization
 - work in $d = 4 - 2\epsilon$ dimensions



Renormalization

- Self-energy insertions and vertex corrections divergent
 - boxes are finite
- redefine quark and gluon fields, and strong coupling to absorb UV divergences
 - add counterterm diagrams to cancel singularities



γ^5 in d-Dimensions

- γ^5 not well defined for $d \neq 4$
- Used non-anticommuting γ^5 scheme

Larin: hep-ph/9302240

$$\gamma^5 = \frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma$$

- Need to symmetrize axial current

$$\gamma^\mu \gamma^5 \rightarrow \frac{1}{2} (\gamma^\mu \gamma^5 - \gamma^5 \gamma^\mu)$$

- using identity above

$$\gamma^\mu \gamma^5 \rightarrow \frac{i}{3!} \epsilon_{\mu\nu\rho\sigma} \gamma^\nu \gamma^\rho \gamma^\sigma$$

γ^5 in d-Dimensions

- γ^5 not well defined for $d \neq 4$
- Used non-anticommuting γ^5 scheme

Larin: hep-ph/9302240

$$\gamma^5 = \frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma$$

- Need to symmetrize axial current

$$\gamma^\mu \gamma^5 \rightarrow \frac{1}{2} (\gamma^\mu \gamma^5 - \gamma^5 \gamma^\mu)$$

- using identity above

$$\gamma^\mu \gamma^5 \rightarrow \frac{i}{3!} \epsilon_{\mu\nu\rho\sigma} \gamma^\nu \gamma^\rho \gamma^\sigma$$

Note: Other calculation approach used naive anticommuting scheme

$$\{\gamma^\mu, \gamma^5\} = 0$$

where 4-dimensional

$$\text{Tr}[\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma^5] = -4i \epsilon^{\mu\nu\rho\sigma}$$

remains

Numerical results in agreement

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $M_{l^+l^-} > 50 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_Z$
- $\alpha_s(M_Z) = 0.118$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}, |\eta_{\text{jet}}| < 2.5$
- $M_{l^+l^-} > 50 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_Z$
- $\alpha_s(M_Z) = 0.118$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

$$\underline{\sqrt{s} = 7 \text{ TeV}}$$

$$\sigma_{\text{int}} = 124.1 \text{ pb}$$

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $M_{l^+l^-} > 50 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_Z$
- $\alpha_s(M_Z) = 0.118$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

$$\sqrt{s} = 7 \text{ TeV}$$

$$\sigma_{\text{int}} = 124.1 \text{ pb}$$

$$m_{sg} = 600 \text{ GeV}$$

$$m_{sq} = 500 \text{ GeV}$$

all squark masses equal
no L-R squark mixing

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $M_{l^+l^-} > 50 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_Z$
- $\alpha_s(M_Z) = 0.118$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

$$\sqrt{s} = 7 \text{ TeV}$$

$$\sigma_{\text{int}} = 124.1 \text{ pb}$$

$$s\text{QCD corr} = 8.784 \text{ fb}$$

$$\delta_{\text{rel}} = 0.007\%$$

$$\delta_{\text{rel}} = \frac{\mathcal{O}_{LO+s\text{QCD}} - \mathcal{O}_{LO}}{\mathcal{O}_{LO}}$$

$$m_{sg} = 600 \text{ GeV}$$

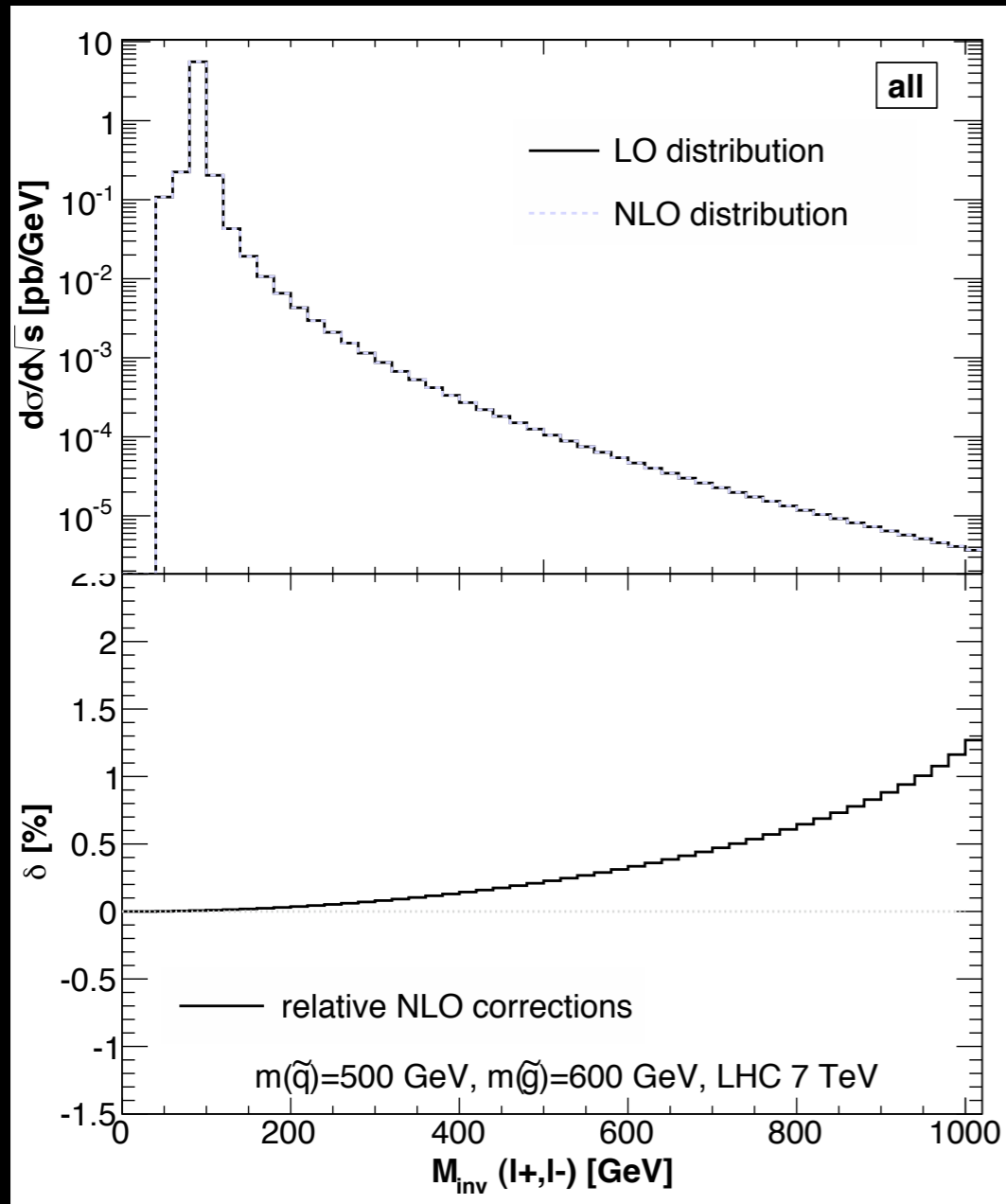
$$m_{sq} = 500 \text{ GeV}$$

all squark masses equal
no L-R squark mixing

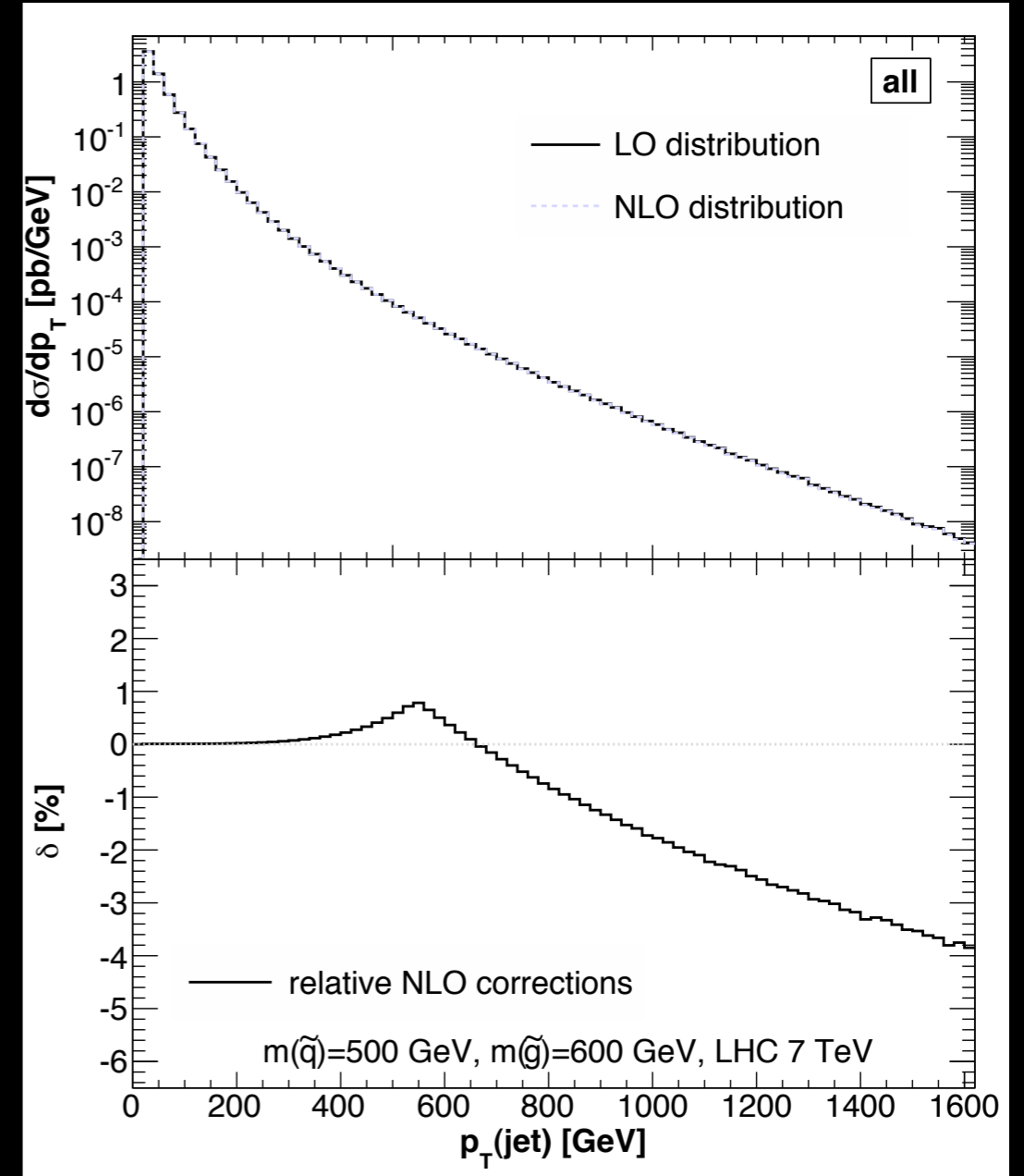
Kinematic Distributions

Preliminary

M_{\parallel}



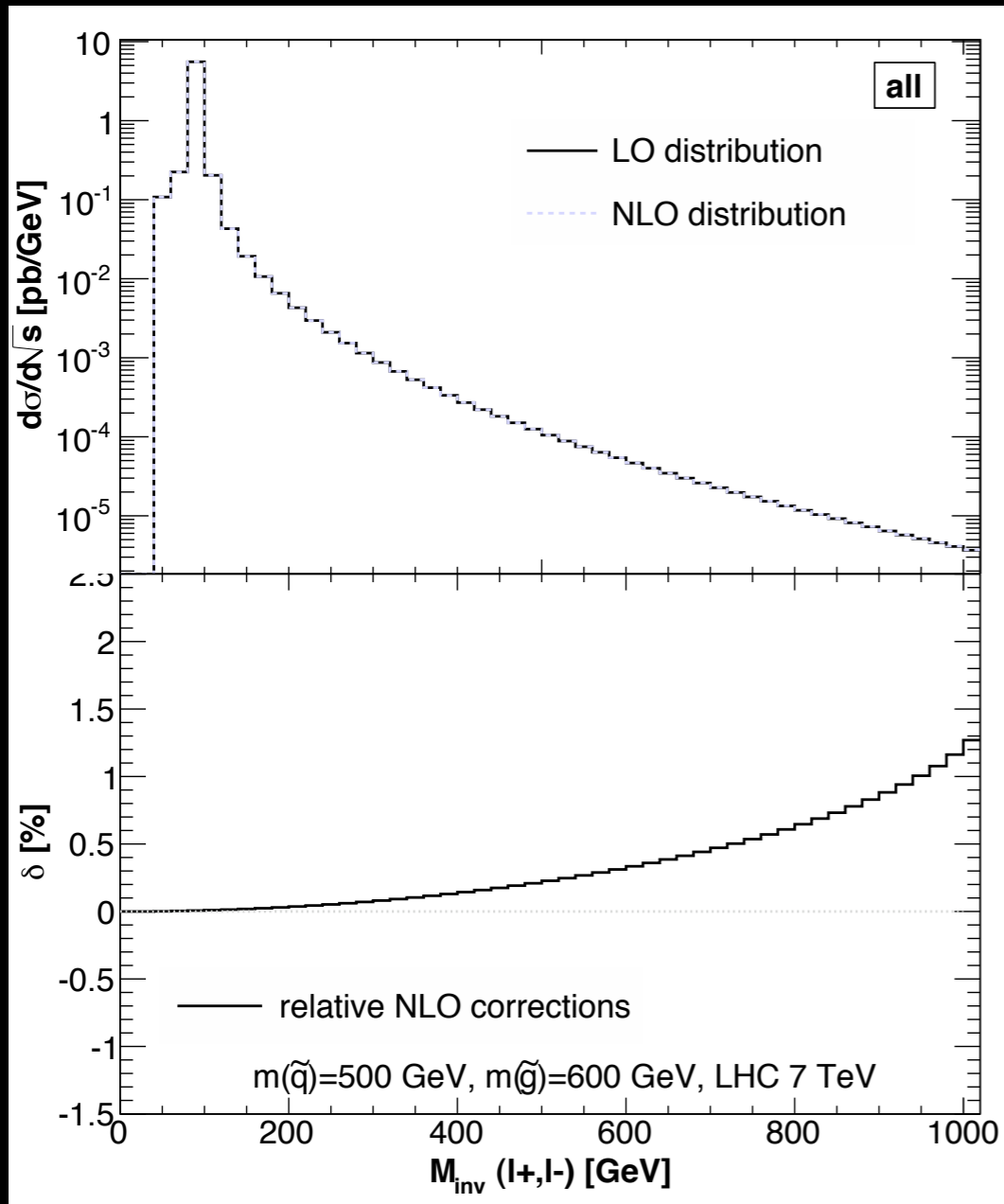
Jet p_T



Kinematic Distributions

Preliminary

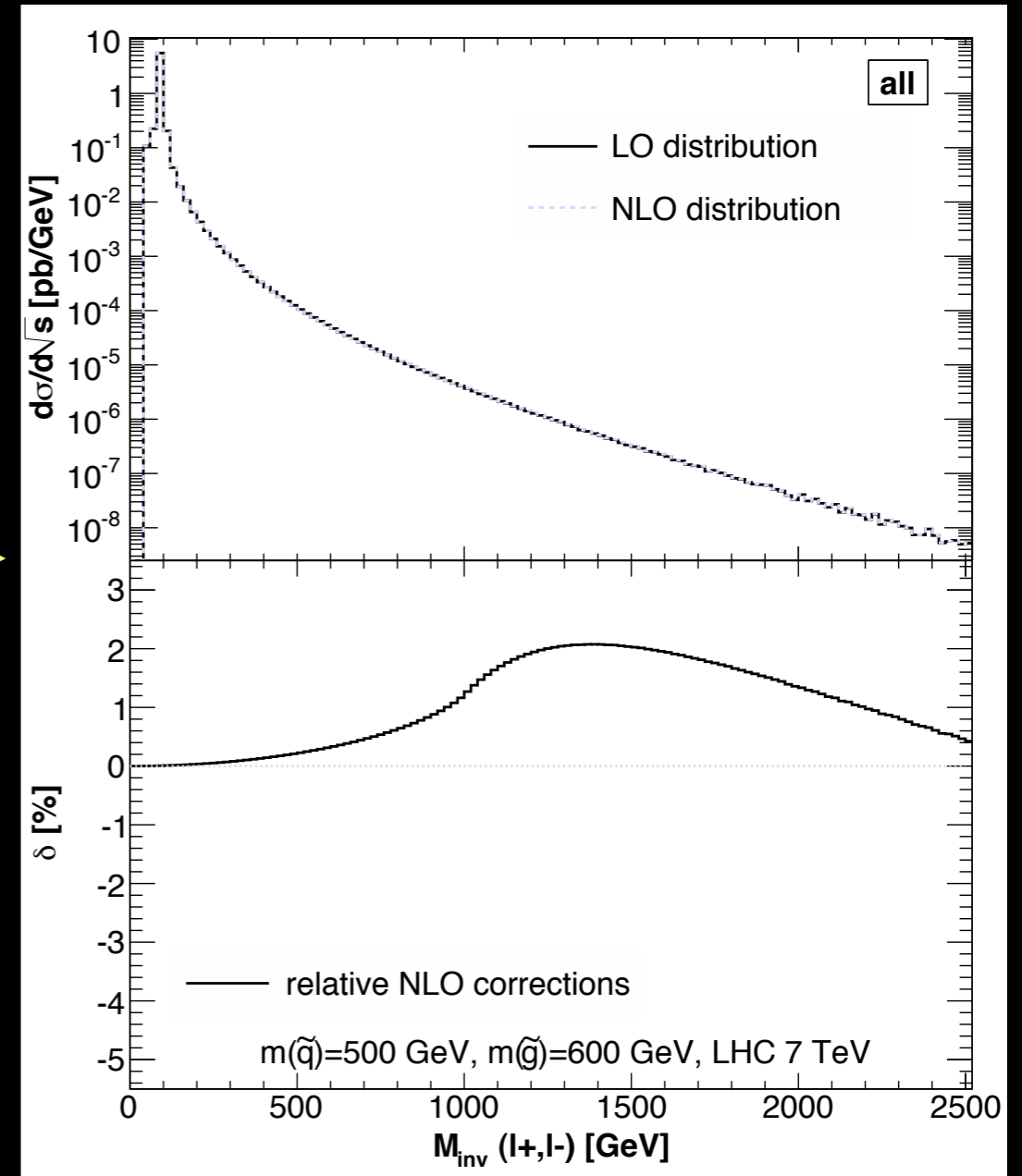
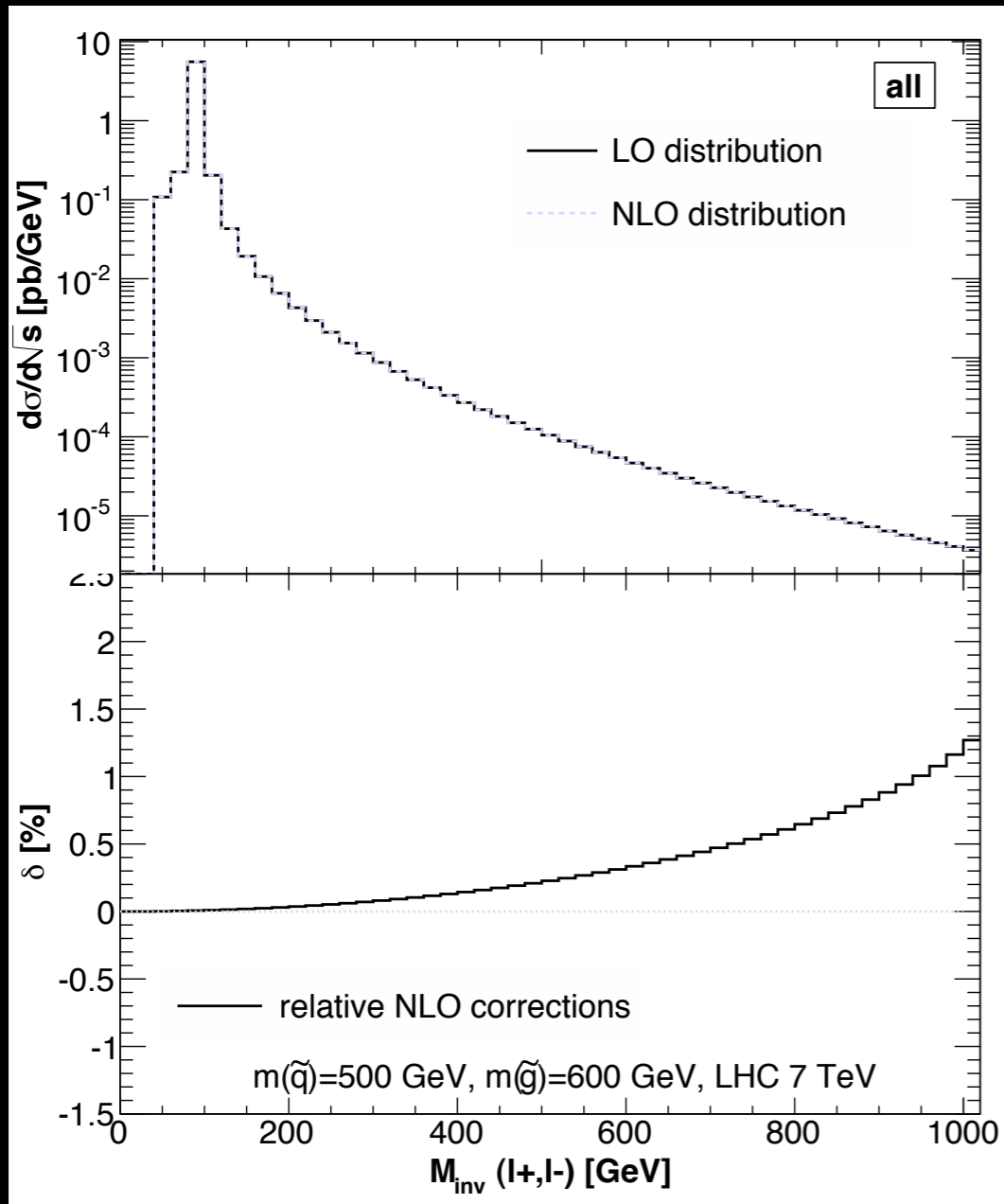
M_{\parallel}



Kinematic Distributions

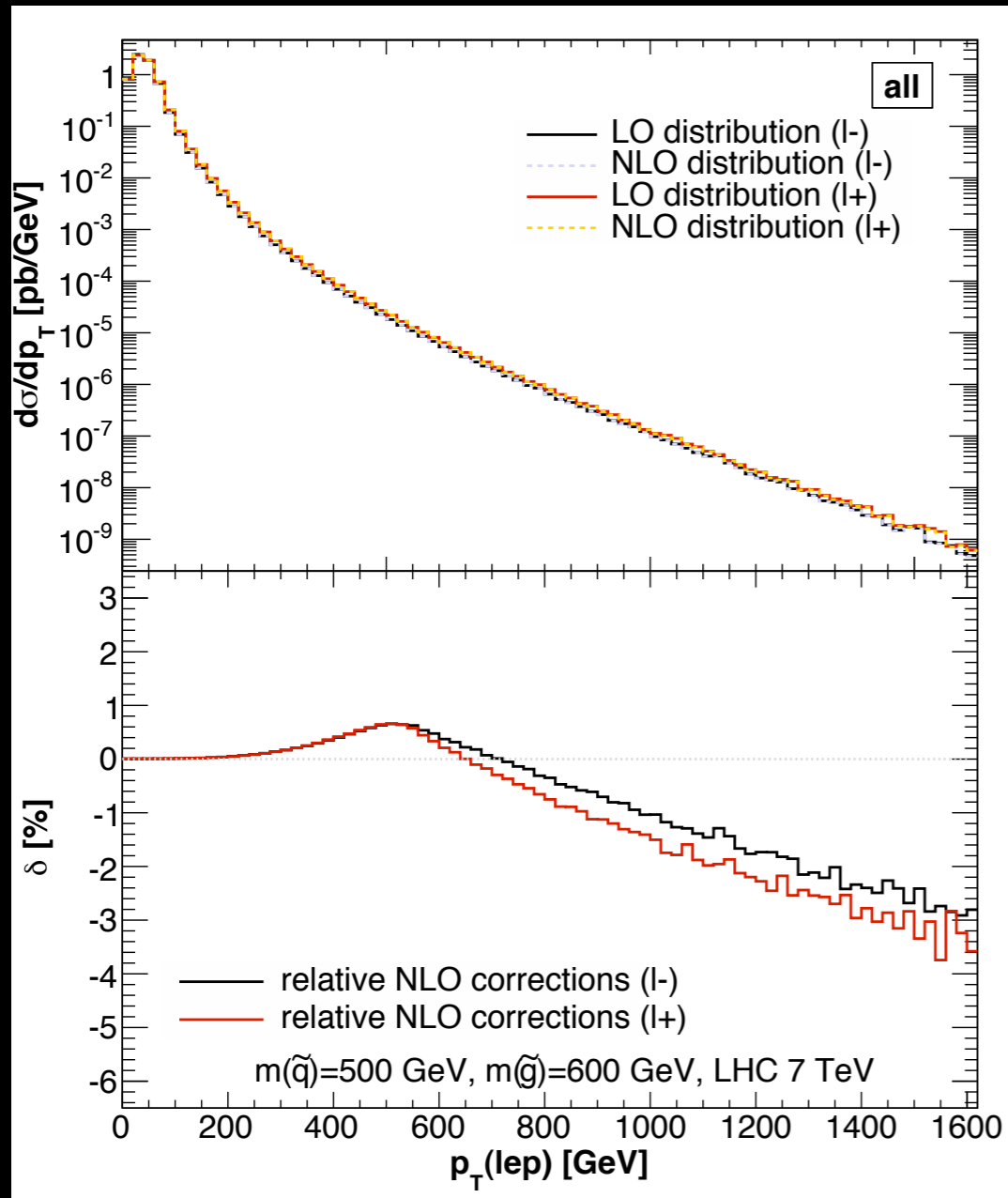
Preliminary

M_{ll}

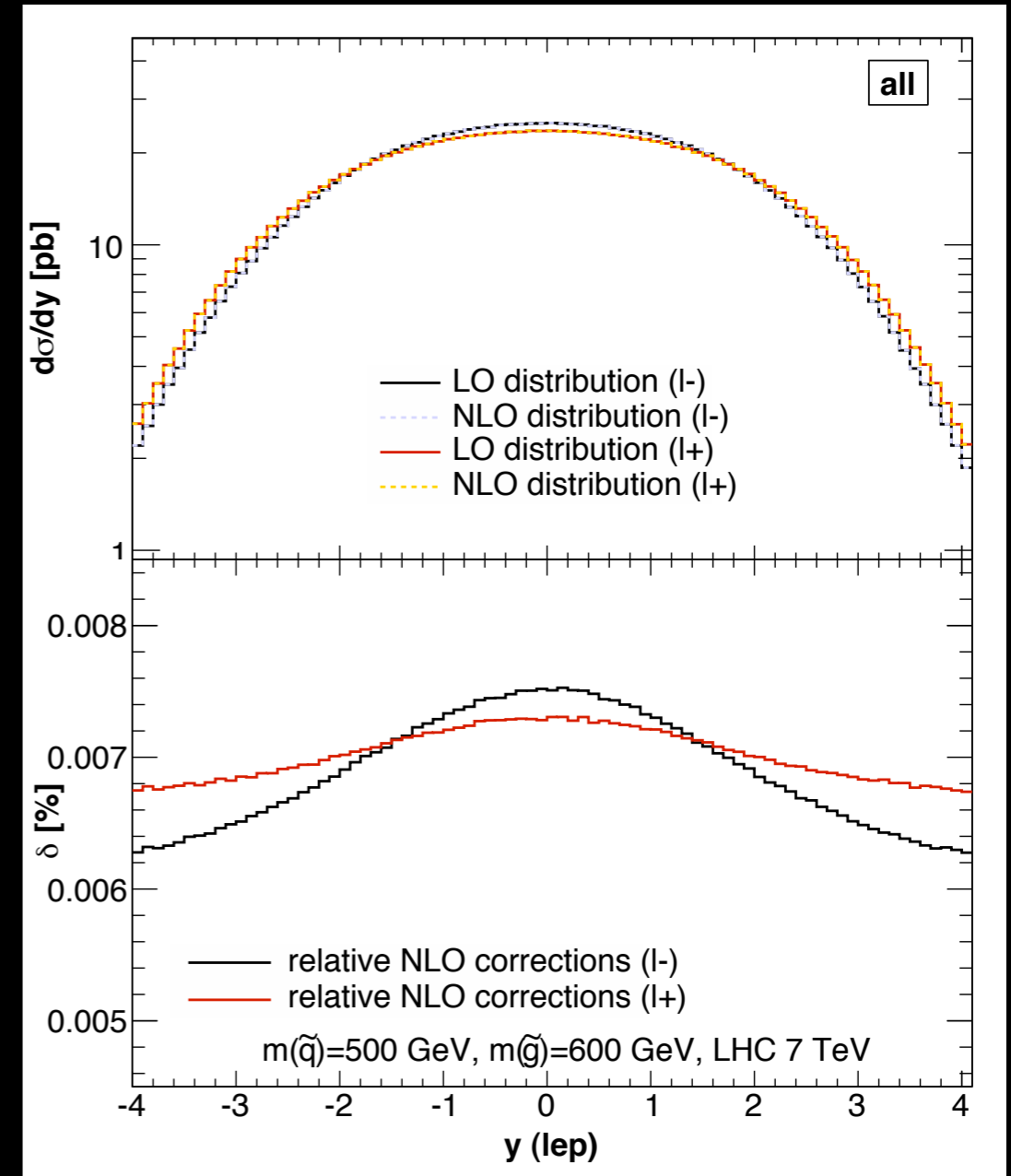


Kinematic Distributions Preliminary

lepton p_T



lepton rapidity



Numerical Results

Preliminary

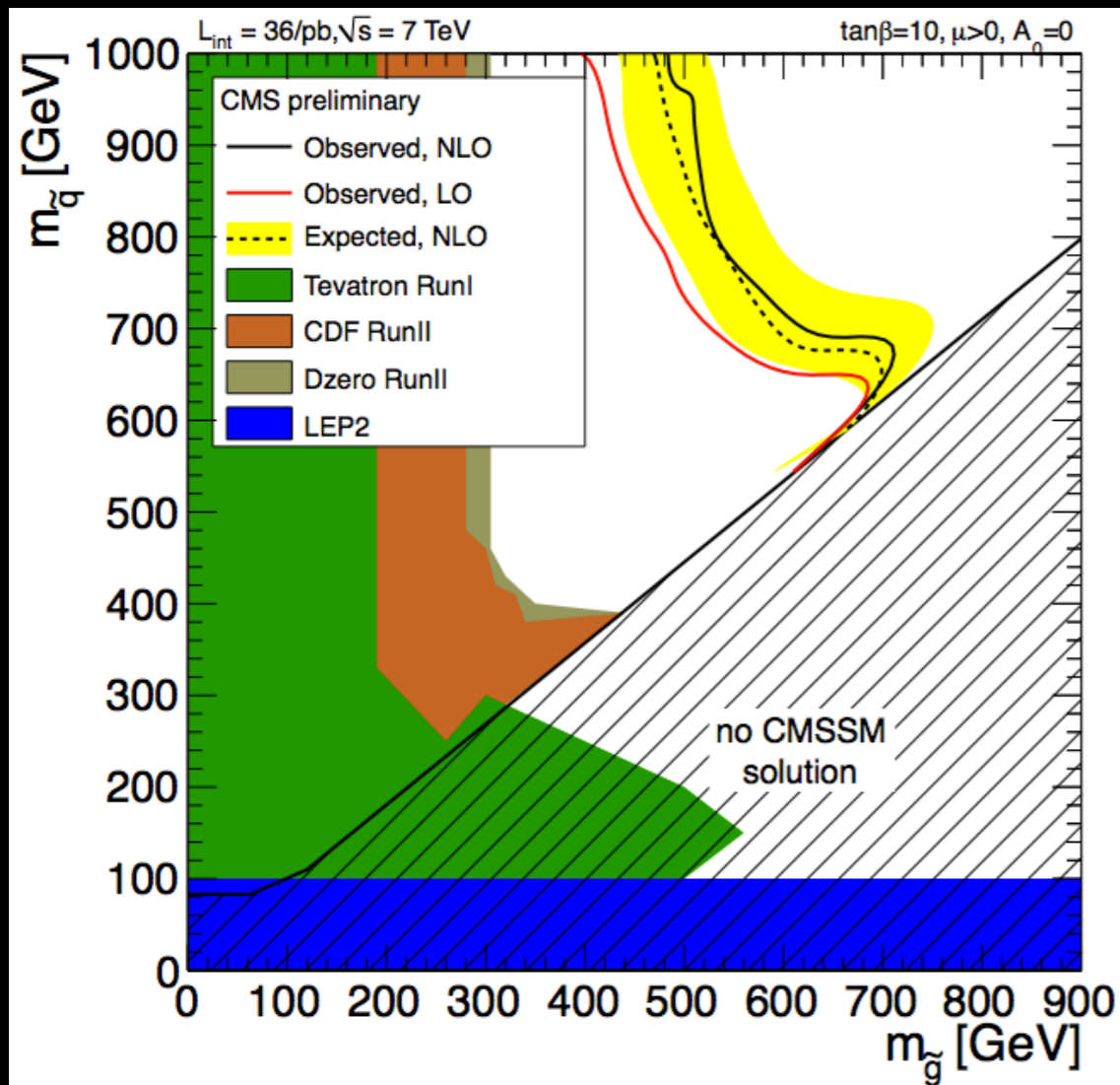
- Different squark and gluino masses?
- Suspect lower masses \rightarrow greater contribution

Numerical Results

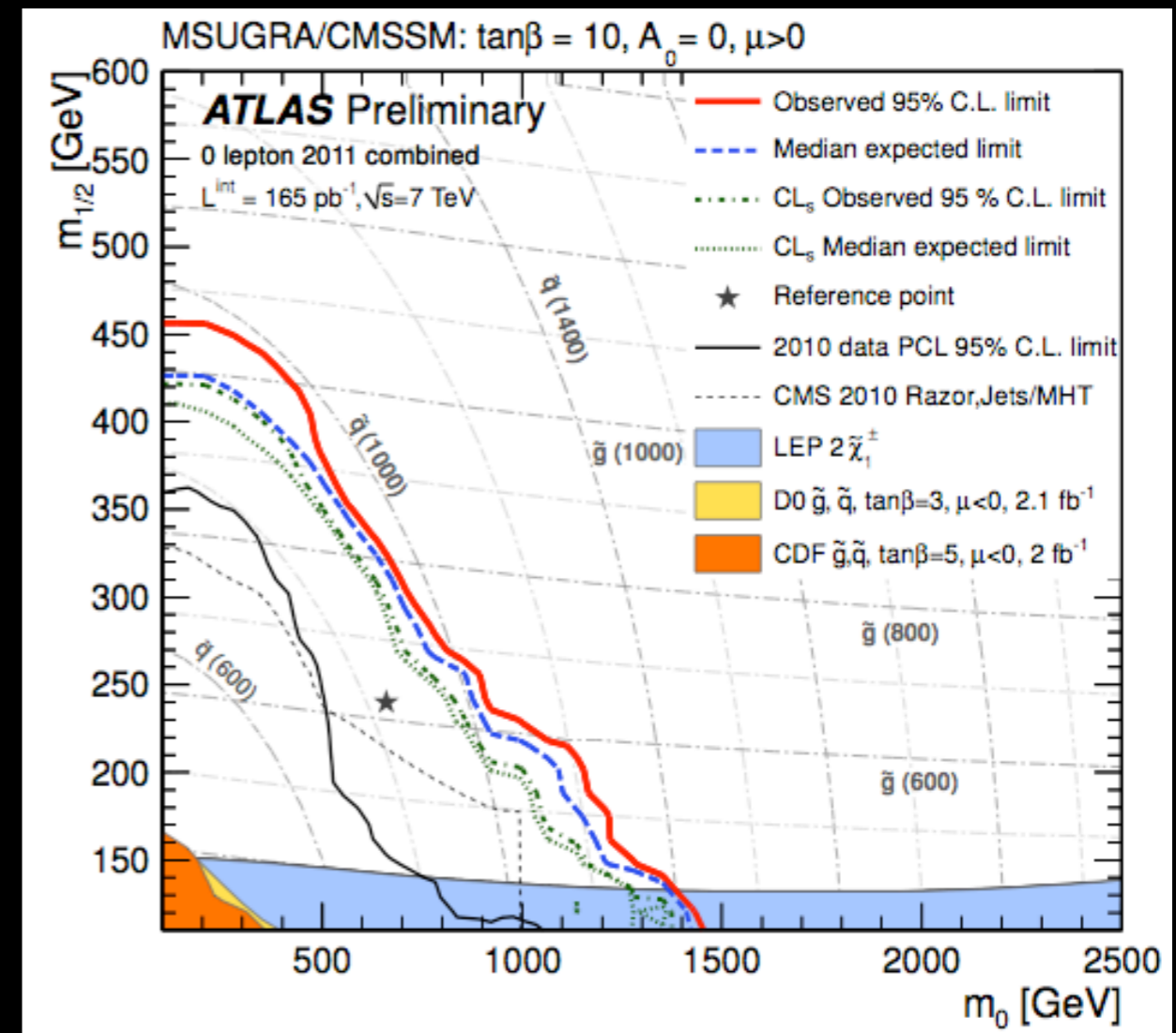
Preliminary

- Different squark and gluino masses?
- Suspect lower masses → greater contribution

CMS PAS SUS-10-005



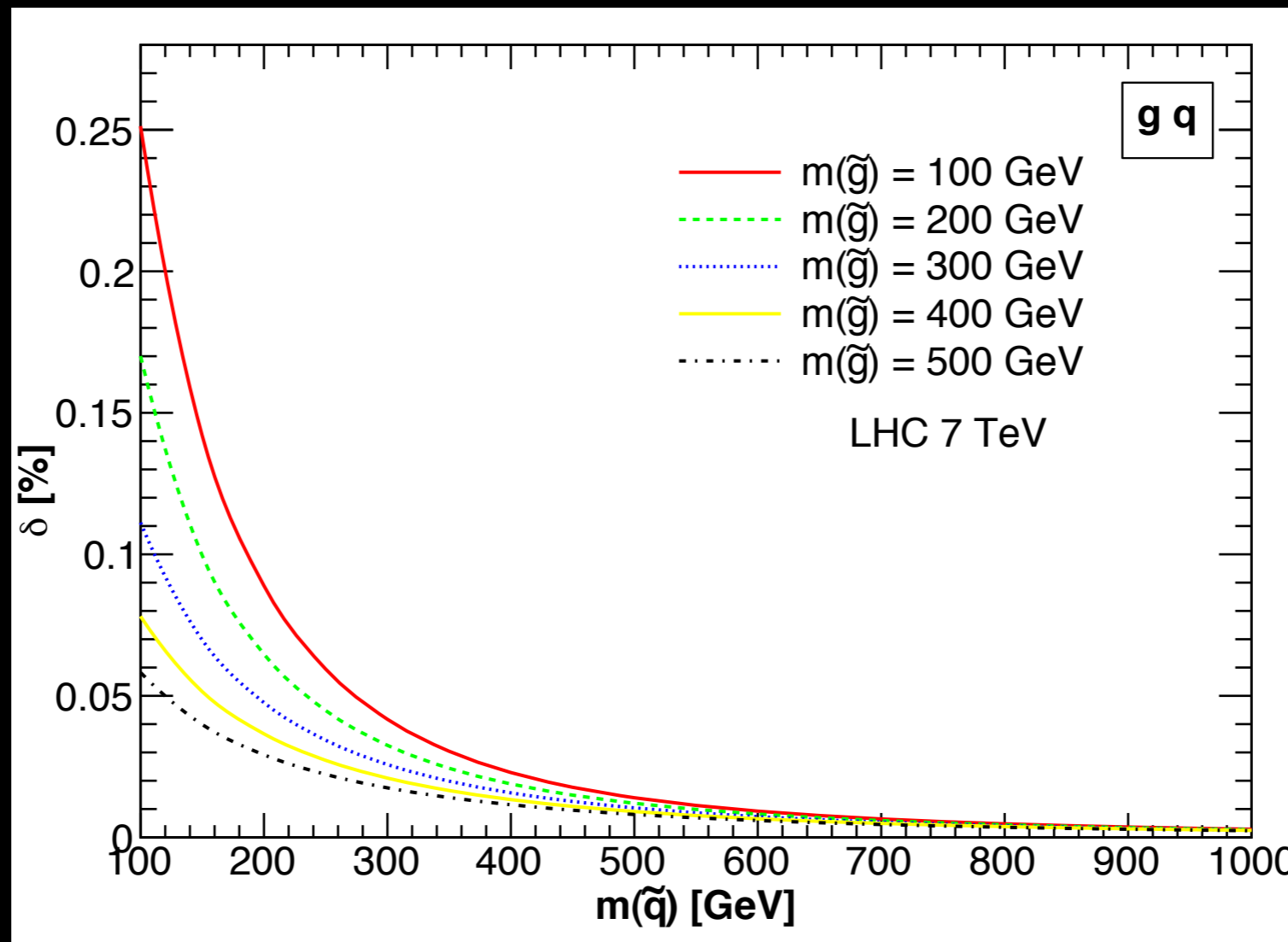
ATLAS-CONF-2011-086



Numerical Results

Preliminary

- Different squark and gluino masses?
- Suspect lower masses \rightarrow greater contribution



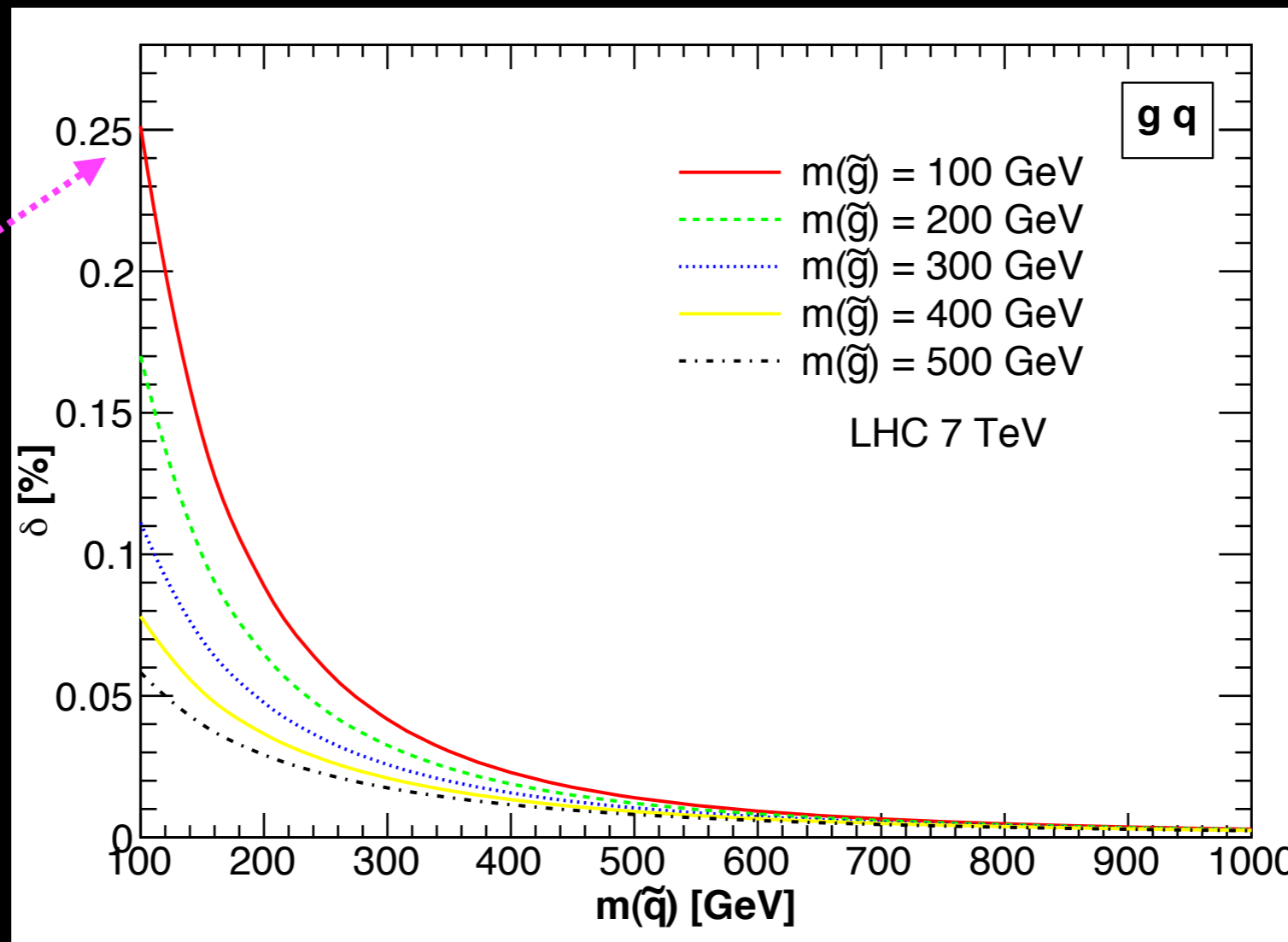
integrated
cross section

gluon-quark
channel only

Numerical Results

Preliminary

- Different squark and gluino masses?
- Suspect lower masses \rightarrow greater contribution



maximal
 $\sim 0.25\%$

integrated
cross section

gluon-quark
channel only

Kinematic Distributions Preliminary

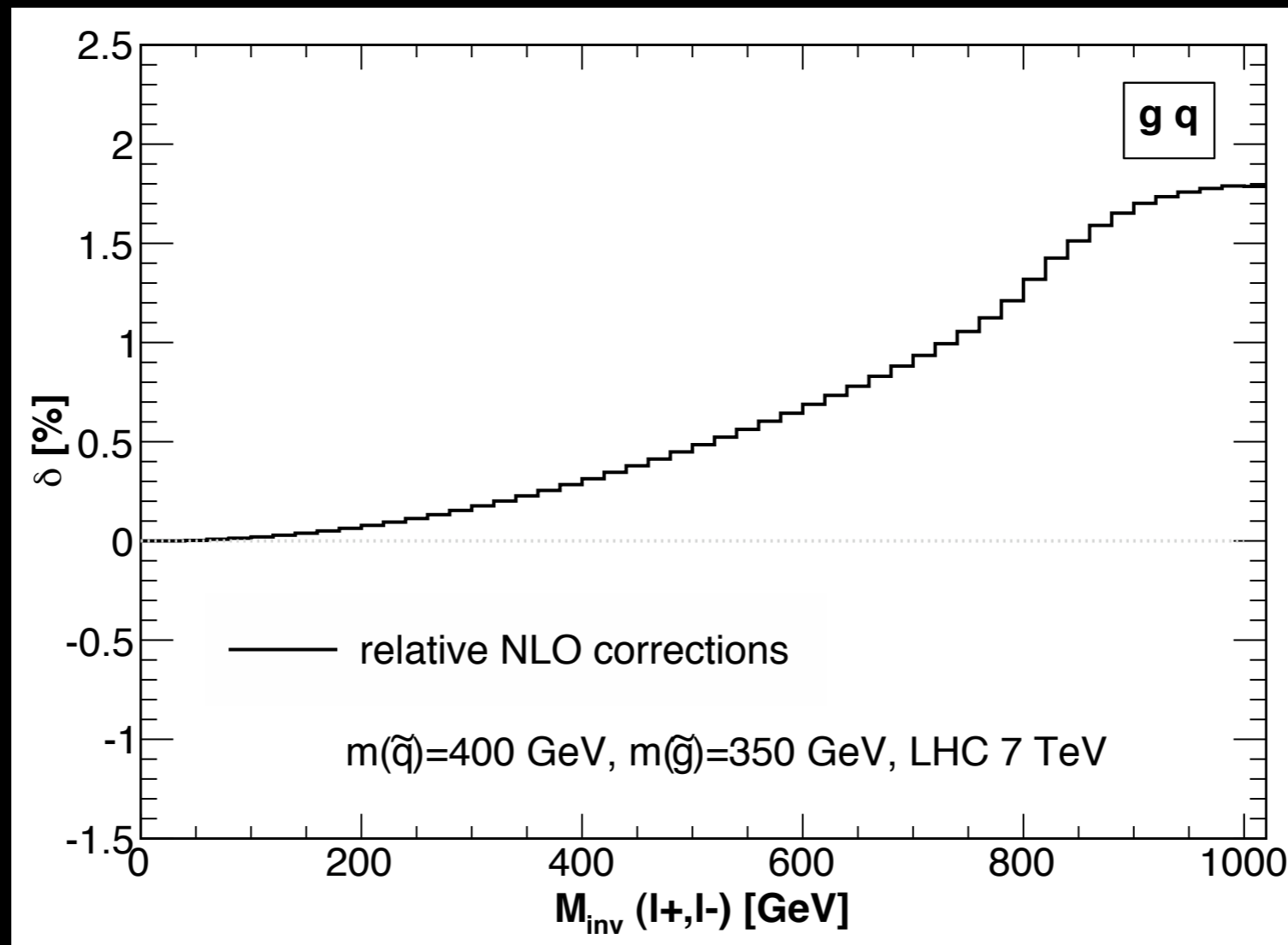
- Different masses, shift in kinematic distributions?
 - gluino mass = 350 GeV, squark mass = 400 GeV

Kinematic Distributions

Preliminary

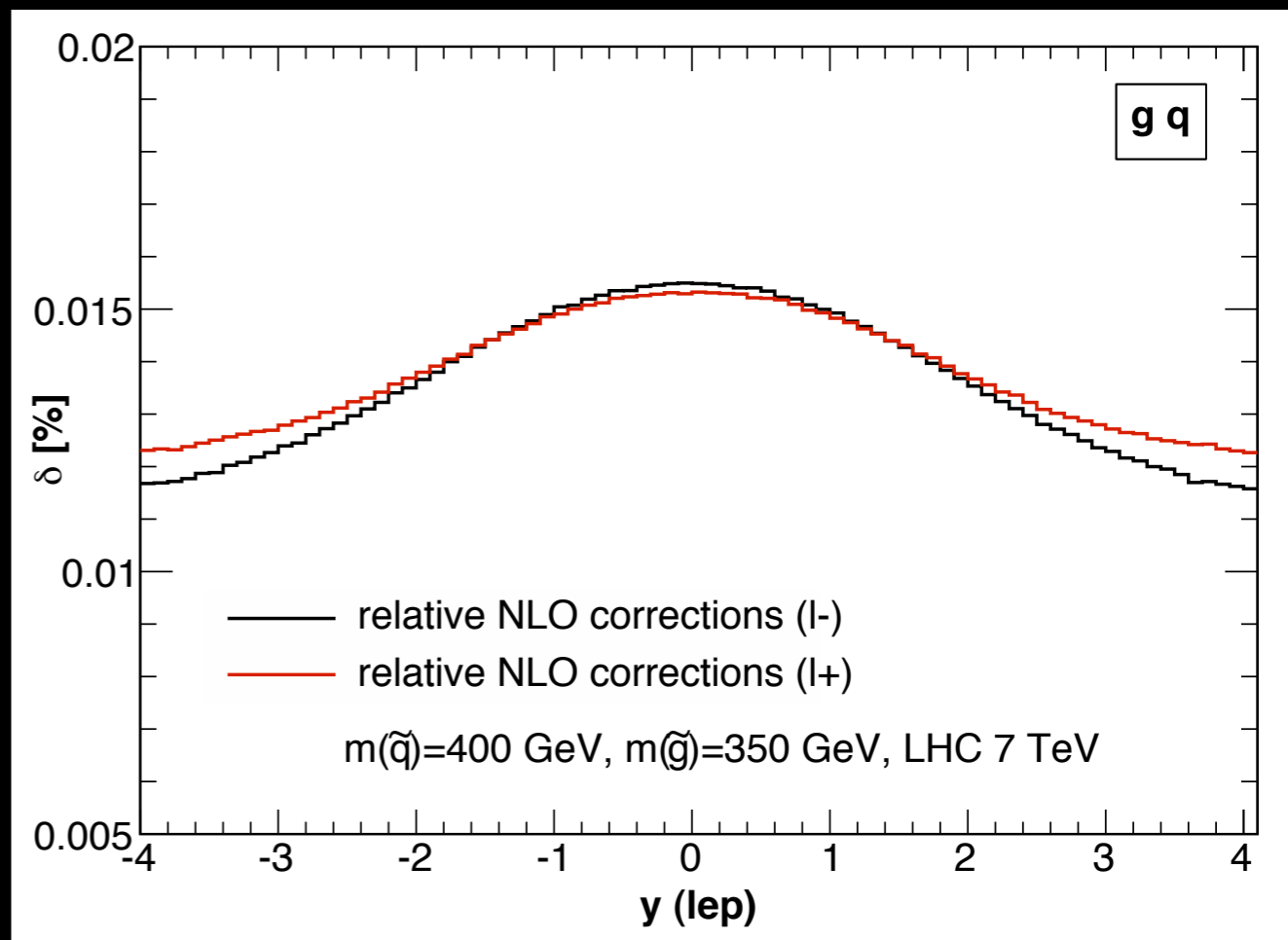
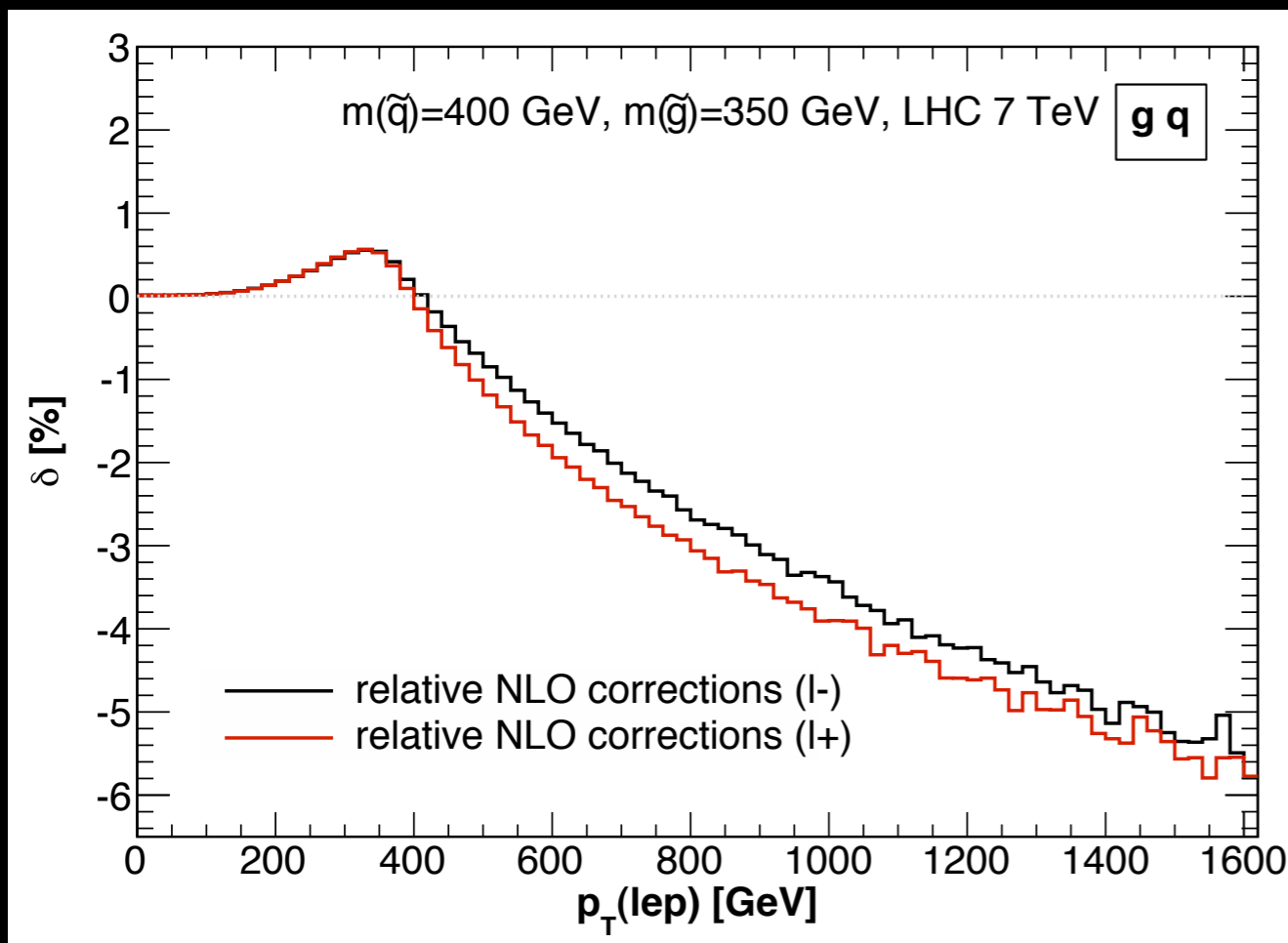
- Different masses, shift in kinematic distributions?

- gluino mass = 350 GeV, squark mass = 400 GeV **gluon-quark channel only**



Kinematic Distributions Preliminary

- Different masses, shift in kinematic distributions?
 - gluino mass = 350 GeV, squark mass = 400 GeV gluon-quark channel only



W+Jet Process

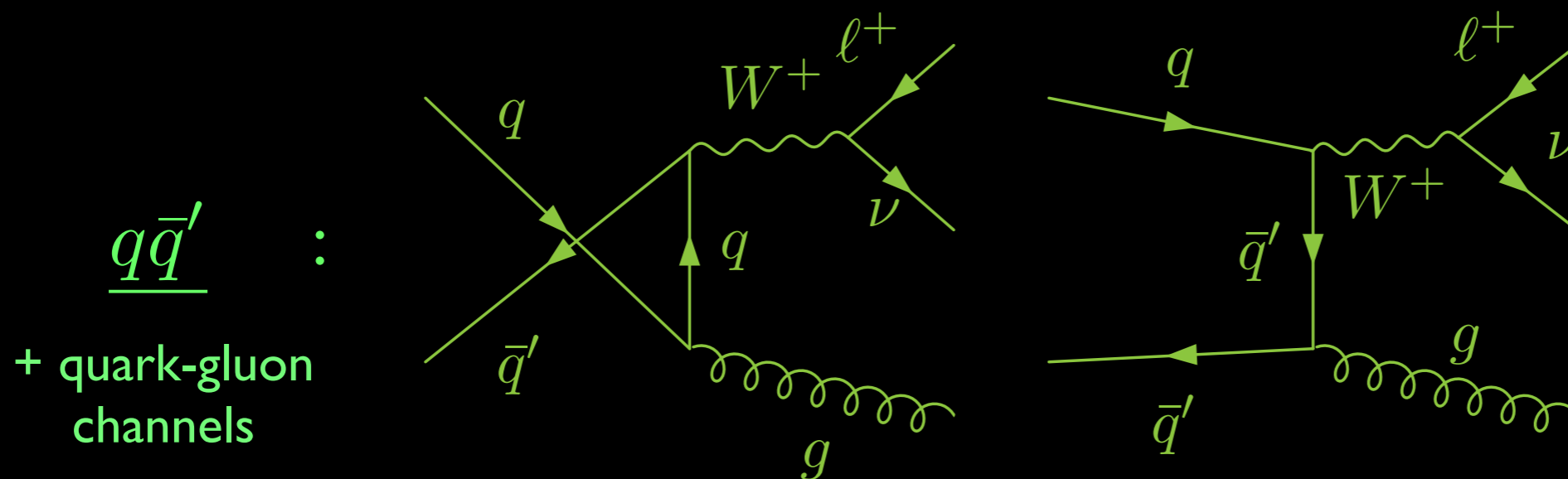
- W+jet production

W+Jet Process

- W+jet production
 - how do sQCD corrections effect this standard candle?

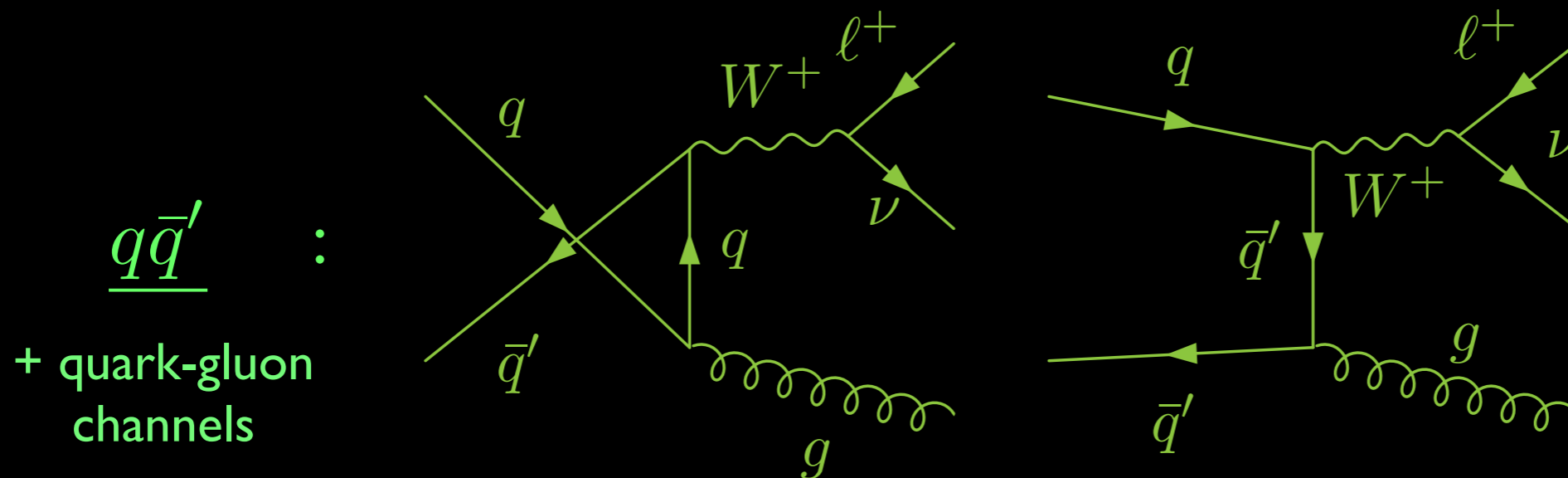
W+Jet Process

- W+jet production
 - how do sQCD corrections effect this standard candle?
- W+jet process similar to Z+jet



W+Jet Process

- W+jet production
 - how do sQCD corrections effect this standard candle?
- W+jet process similar to Z+jet



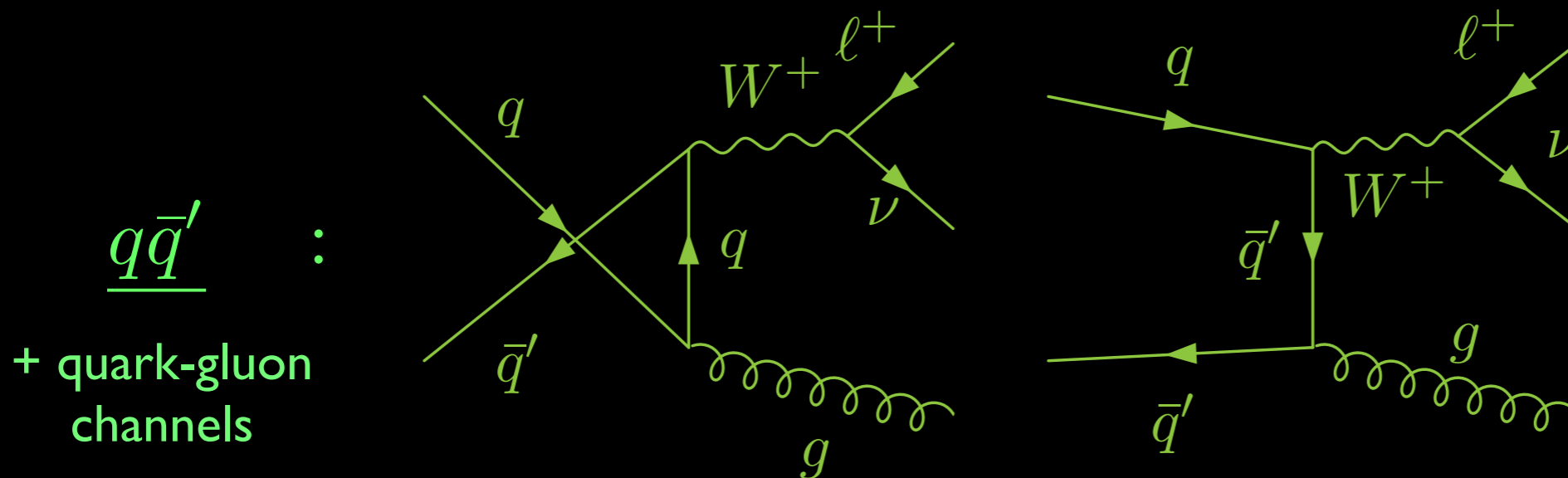
- same NLO structure: self-energy insertion, gluon and W vertex corrections, box terms

W+Jet Process

- W+jet production
 - how do sQCD corrections effect this standard candle?

- W+jet process similar to Z+jet

only “left-handed” squarks contribute to W+jet with no L-R squark mixing



- same NLO structure: self-energy insertion, gluon and W vertex corrections, box terms

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $p_{T,\text{lep}} > 25 \text{ GeV}$, $E_{T,\text{miss}} > 25 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_W$
- $\alpha_s(M_W) = 0.1203$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $p_{T,\text{lep}} > 25 \text{ GeV}$, $E_{T,\text{miss}} > 25 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_W$
- $\alpha_s(M_W) = 0.1203$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

$$\sqrt{s} = 7 \text{ TeV}$$

$$\sigma_{\text{all,int}} = 635.94 \text{ pb}$$

all = W^+ & W^-

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $p_{T,\text{lep}} > 25 \text{ GeV}$, $E_{T,\text{miss}} > 25 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_W$
- $\alpha_s(M_W) = 0.1203$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

$$\sqrt{s} = 7 \text{ TeV}$$

$$\sigma_{\text{all,int}} = 635.94 \text{ pb}$$

all = W^+ & W^-

$$m_{sg} = 600 \text{ GeV}$$

$$m_{sq} = 500 \text{ GeV}$$

all squark masses
equal, no L-R squark
mixing

Numerical Results

Preliminary

- $p_{T,\text{jet}} > 20 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.5$
- $p_{T,\text{lep}} > 25 \text{ GeV}$, $E_{T,\text{miss}} > 25 \text{ GeV}$
- CTEQ6M PDF set
- $\mu_R = \mu_F = M_W$
- $\alpha_s(M_W) = 0.1203$
- $\alpha \sim 1/132$
- u,d,c,s initial state quarks

$$\sqrt{s} = 7 \text{ TeV}$$

$$\sigma_{\text{all,int}} = 635.94 \text{ pb}$$

all = W^+ & W^-

$$s\text{QCD corr} = 36.94 \text{ fb}$$

$$\delta_{\text{rel}} = 0.006\%$$

$$m_{sg} = 600 \text{ GeV}$$

$$m_{sq} = 500 \text{ GeV}$$

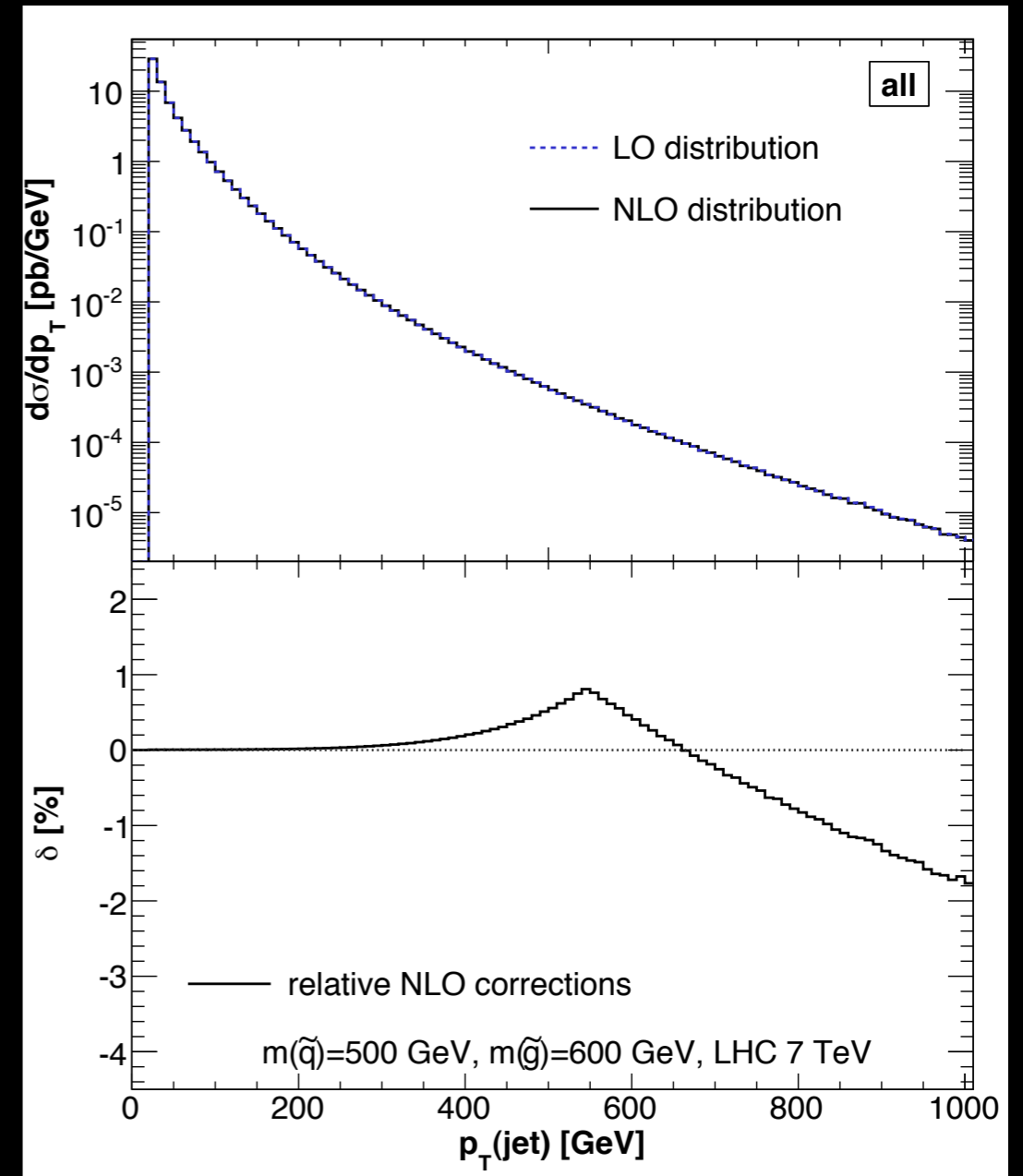
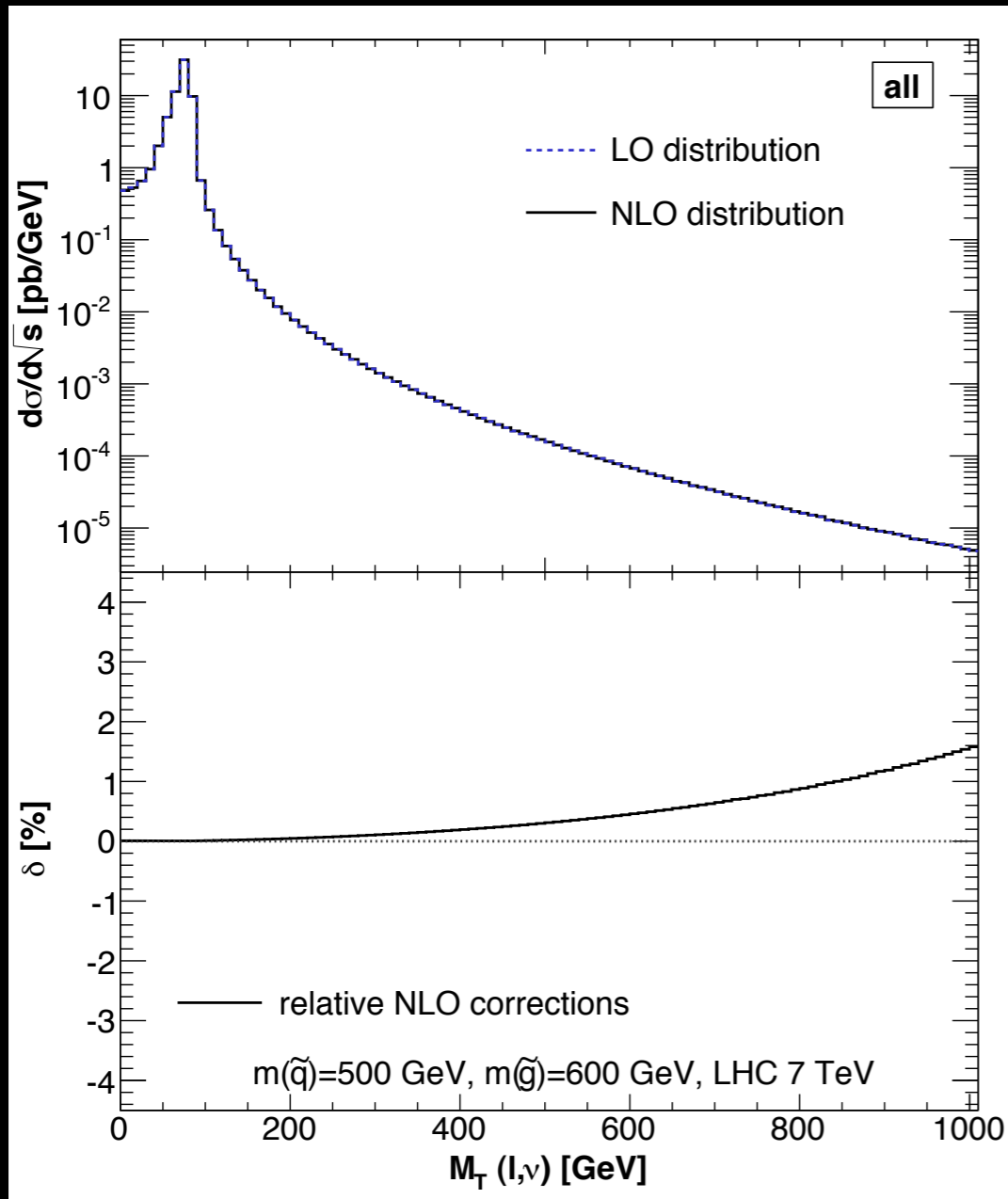
all squark masses
equal, no L-R squark
mixing

Kinematic Distributions

Preliminary

M_T

Jet p_T

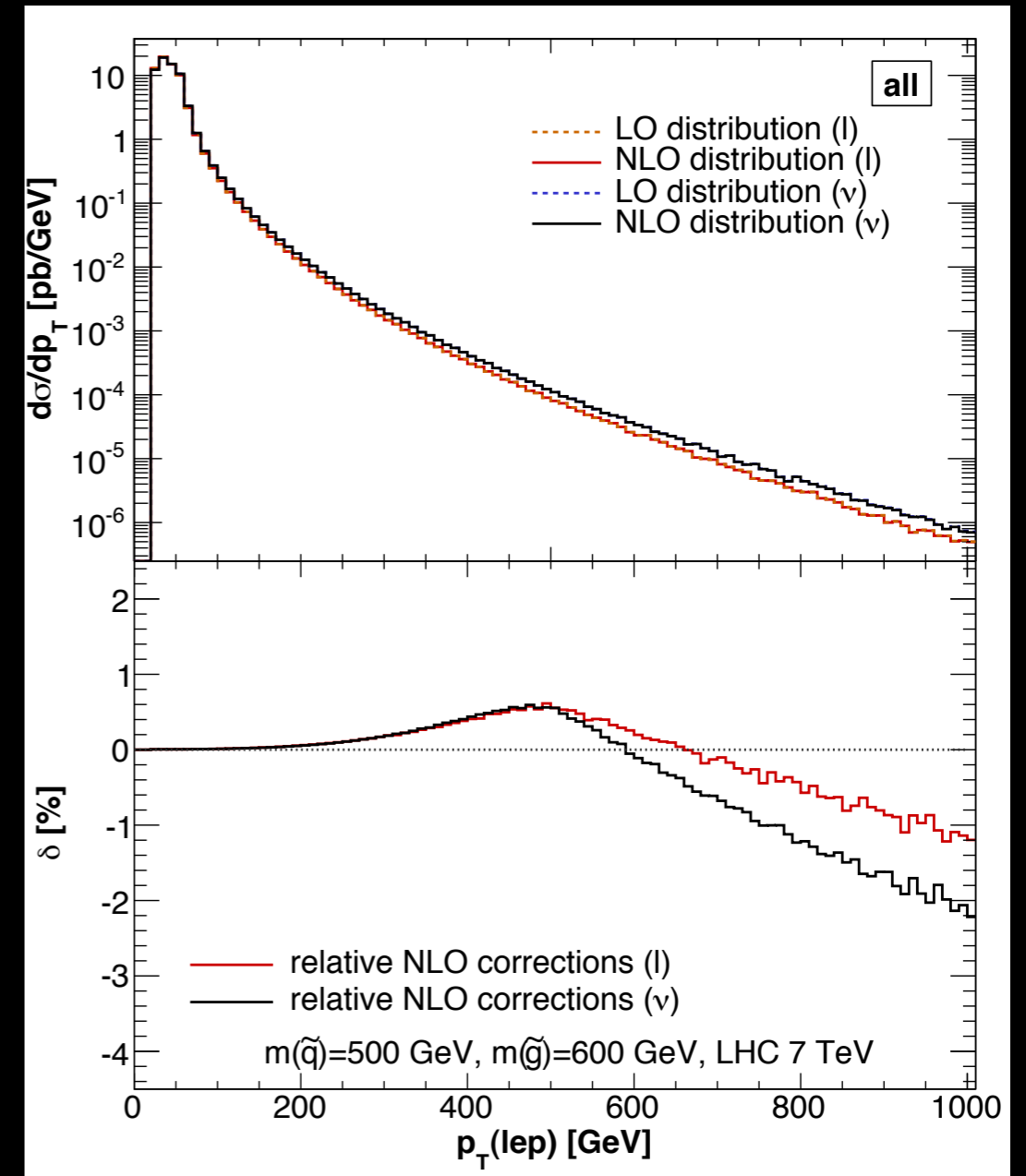
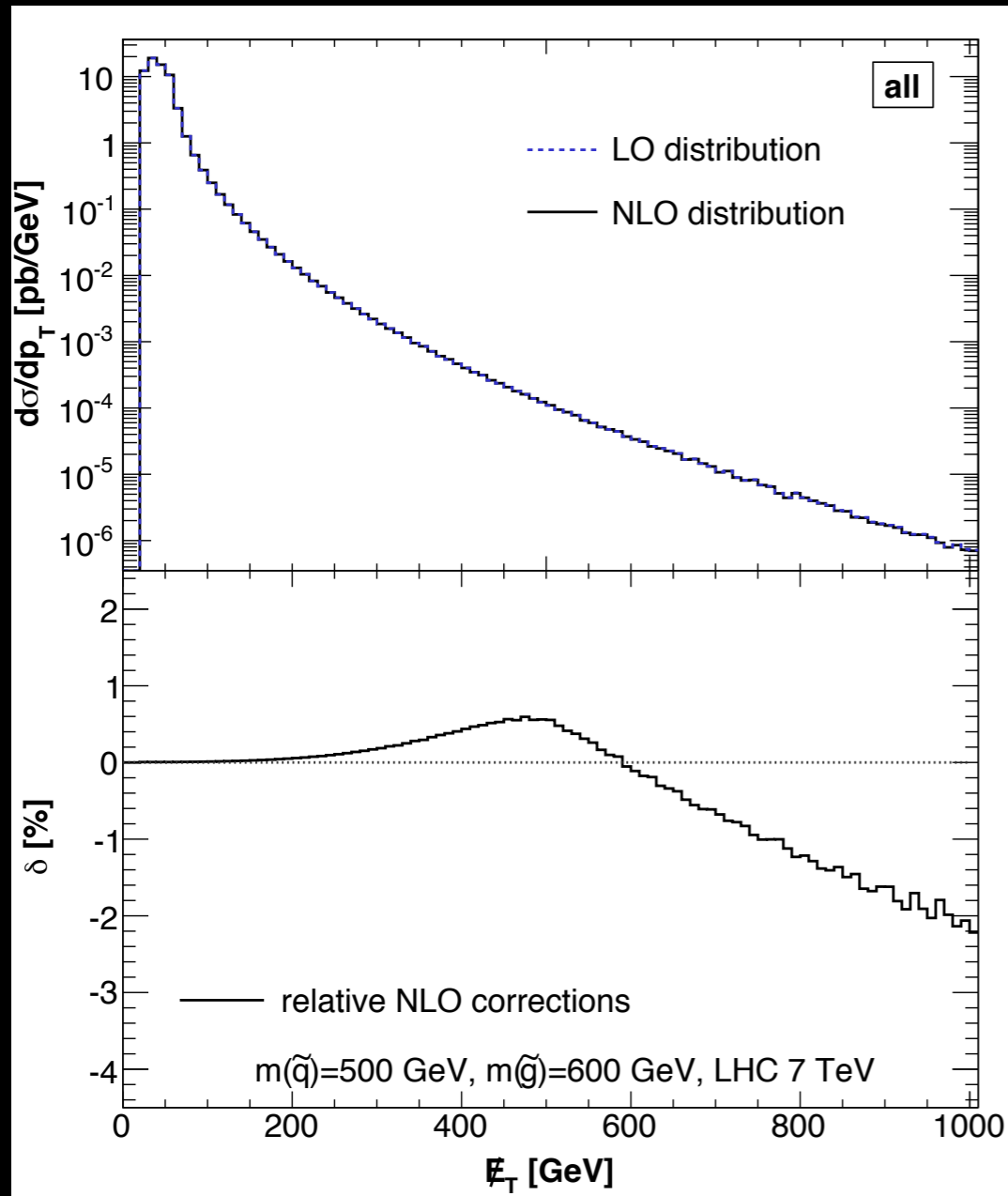


Kinematic Distributions

Preliminary

$E_{T,miss}$ p_T

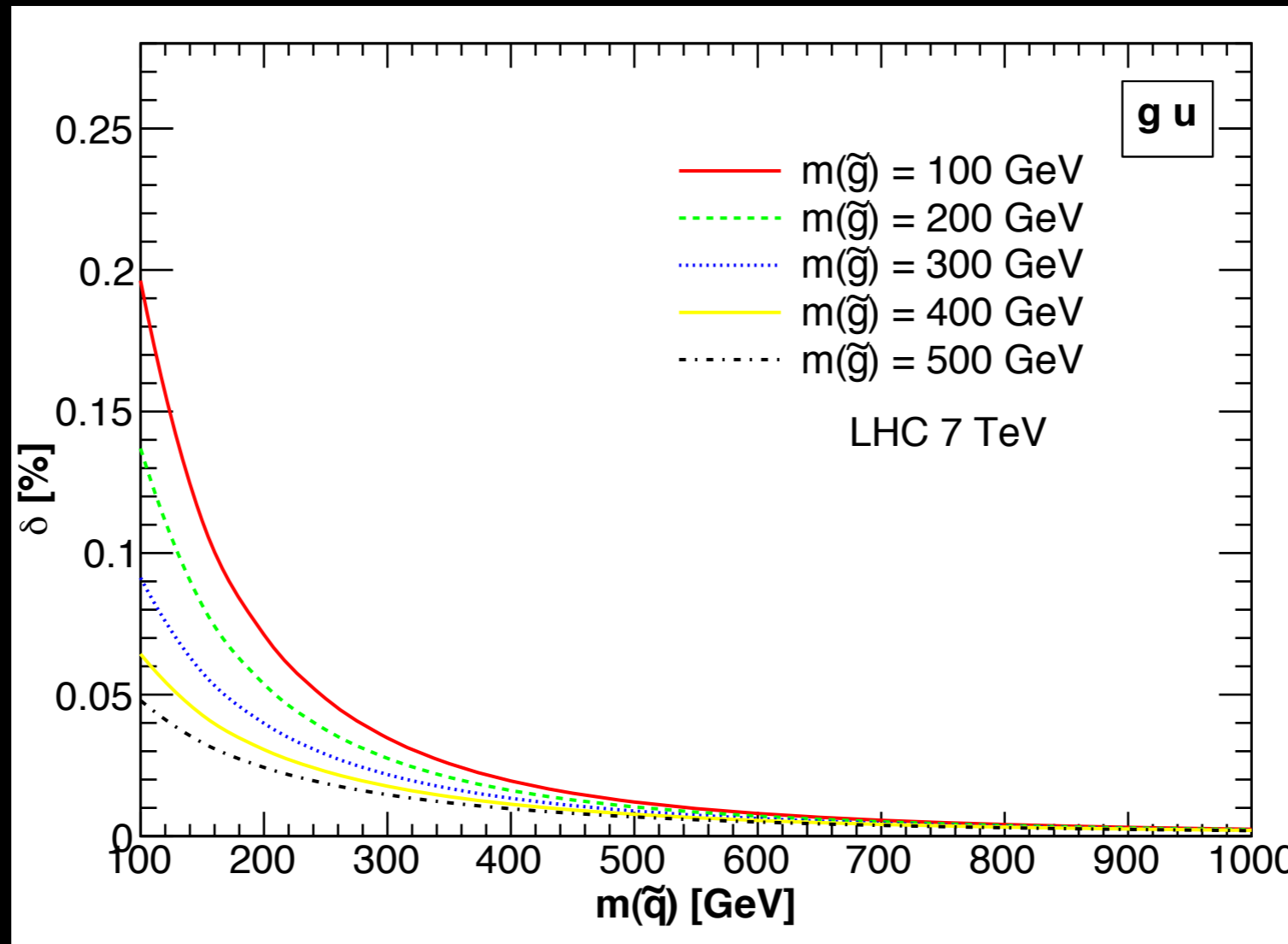
lepton p_T



Numerical Results

Preliminary

- Different squark and gluino masses?



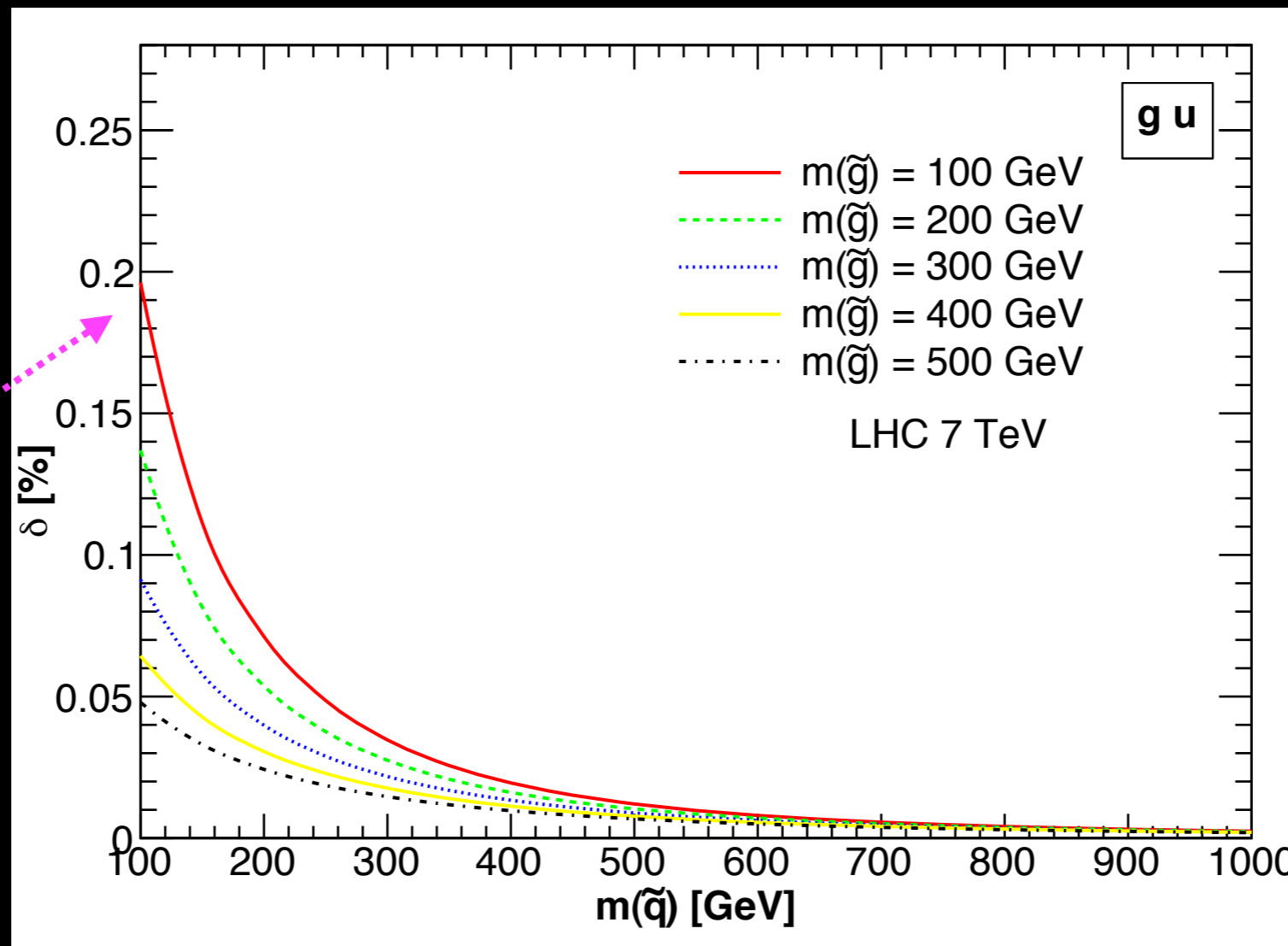
integrated
cross section

gluon-quark
channel only

Numerical Results

Preliminary

- Different squark and gluino masses?



maximal
~0.20%

integrated
cross section

gluon-quark
channel only

sQCD to Drell-Yan

sQCD to Drell-Yan

- Drell-Yan at NLO in sQCD performed previously (no jet)

sQCD to Drell-Yan

- Drell-Yan at NLO in sQCD preformed previously (no jet)
 - neutral current [arXiv:0911.2329](https://arxiv.org/abs/0911.2329)
 - low invariant mass range: $\sim 0.1\%$ high (TeV): $\sim 1-2\%$

sQCD to Drell-Yan

- Drell-Yan at NLO in sQCD preformed previously (no jet)
 - neutral current [arXiv:0911.2329](https://arxiv.org/abs/0911.2329)
 - low invariant mass range: $\sim 0.1\%$ high (TeV): $\sim 1-2\%$
 - charged current [arXiv:0710.3309](https://arxiv.org/abs/0710.3309)
 - low invariant mass range: $< 0.1\%$ high (TeV): $\sim 1\%$

sQCD to Drell-Yan

- Drell-Yan at NLO in sQCD preformed previously (no jet)
 - neutral current [arXiv:0911.2329](https://arxiv.org/abs/0911.2329)
 - low invariant mass range: $\sim 0.1\%$ high (TeV): $\sim 1-2\%$
 - charged current [arXiv:0710.3309](https://arxiv.org/abs/0710.3309)
 - low invariant mass range: $< 0.1\%$ high (TeV): $\sim 1\%$
- W/Z + jet production at NLO in sQCD is in good agreement with DY scenario

Conclusions

- EW gauge boson production is still a very important process at the LHC
 - standard candles
 - key processes for EW precision physics & PDFs
 - implications for BSM physics

Conclusions

- EW gauge boson production is still a very important process at the LHC
 - standard candles
 - key processes for EW precision physics & PDFs
 - implications for BSM physics
- Investigate Z/W +jet at NLO in sQCD
 - small K-factor
 - percent level relative corrections in large p_T kinematic regions
 - size of effect in agreement with earlier $DY+0$ jet studies

Conclusions

- EW gauge boson production is still a very important process at the LHC
 - standard candles
 - key processes for EW precision physics & PDFs
 - implications for BSM physics
- Investigate Z/W +jet at NLO in sQCD
 - small K-factor
 - percent level relative corrections in large p_T kinematic regions
 - size of effect in agreement with earlier $DY+0$ jet studies
- Outlook
 - Other BSM effects to Drell-Yan processes?
 - are they always smaller than theoretical and PDF uncertainties?

Conclusions

- EW gauge boson production is still a very important process at the LHC
 - standard candles
 - key processes for EW precision physics & PDFs
 - implications for BSM physics
- Investigate Z/W +jet at NLO in sQCD
 - small K-factor
 - percent level relative corrections in large p_T kinematic regions
 - size of effect in agreement with earlier $DY+0$ jet studies
- Outlook
 - Other BSM effects to Drell-Yan processes?
 - are they always smaller than theoretical and PDF uncertainties?
- Z/W +jet a very stable standard candle under sQCD corrections