

NLO corrections to gauge-boson scattering at the LHC

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- 1 Motivation
 - Search for a light Higgs boson
 - Study of Electroweak Symmetry Breaking
- 2 What kinds of diagrams are involved?
- 3 What has been done so far?
 - Tree-level
 - NLO QCD corrections
- 4 My work
 - Why block structure?
 - Elements of calculation
- 5 Summary

SM Higgs Search

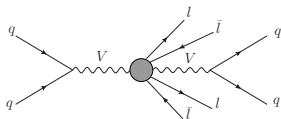
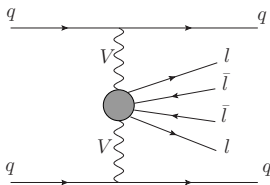


- vector boson fusion ($qq \rightarrow qqH$)
 - possible discovery mode for a light SM Higgs boson
 - second largest cross-section for the light Higgs
 - no color exchange between initial quarks
 - background can be reduced thanks to forward jet tagging
 - relatively low luminosity needed for discovery
- vector boson scattering ($qq \rightarrow qqVV$) - background to $H \rightarrow VV$ decay mode
- precise cross-section required to distinguish the Higgs signal

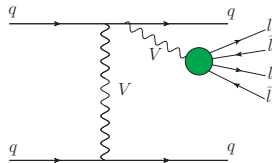
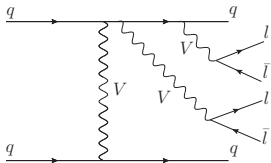
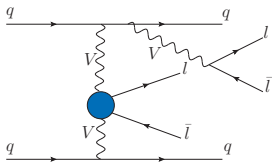
What if there is no light Higgs?

- new mechanism of EW symmetry breaking must be considered
- experiments impose constraints on the new strong dynamics (symmetry, mass scale, dynamics)
- reason to expect strong vector boson scattering
 - new effects preventing unitarity violation
- $qq \rightarrow qqWW$
 - forward jet tagging and energy cuts reduce other background

Diagrams - Tree Level



- following types of diagrams have to be included (to preserve gauge invariance)

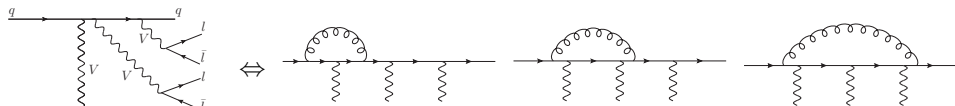
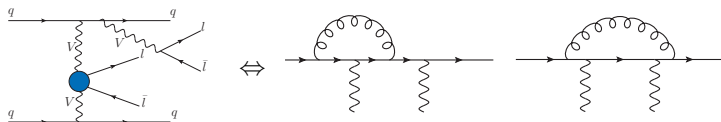


Tree-level studies

- first partial results [Cahn, Dawson \(1984\)](#)
- $pp \rightarrow qqWW$ in effective gauge boson approximation, only for longitudinal polarization [Duncan, Kane, Repko \(1986\)](#)
- exact calculation of $pp \rightarrow qqWW$, all polarizations [Dicus, Vega \(1986\)](#)
- $pp \rightarrow qqZZ$, effective gauge boson approximation [Abbasabadi, Repko \(1988\)](#)
- $pp \rightarrow qqZZ \rightarrow qqllll$, narrow width approximation [Baur, Glover \(1990\)](#)
- $pp \rightarrow (qqZW \rightarrow qqZW) + X$, effective gauge boson approximation, longitudinal polarization [Dobado, Herrero, Terron \(1991\)](#)
- $pp \rightarrow qqZW$, full tree-level, leptonic decay correlations [Barger, Cheung, Han, Stange, Zeppenfeld \(1992\)](#)
- $pp \rightarrow qqWW$ - electroweak chiral lagrangian formalism, semileptonic decay [Butterworth, Cox, Forshaw \(2002\)](#)

Diagrams - NLO QCD Contributions

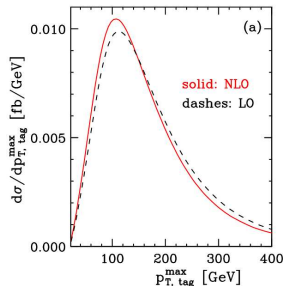
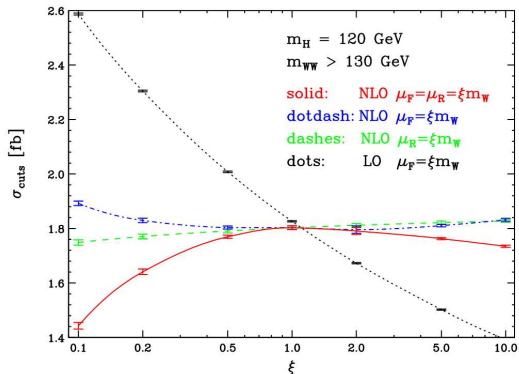
Virtual corrections



NLO QCD calculation

- full tree-level calculation and NLO QCD corrections (real and virtual contributions)
 - 2006 - Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjWW \rightarrow jjllll$
 - 2006 - Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjZZ \rightarrow jjllll$
 - 2007 - Bozzi, Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjWZ \rightarrow jjllll$
- use of 'leptonic tensors'
 - faster code,
 - possibility to include new EW effects
- s-channel and interference neglected (due to VBF kinematics)
- analytical code in Maple, numerical in Fortran
- Monte Carlo program very fast thanks to modular structure

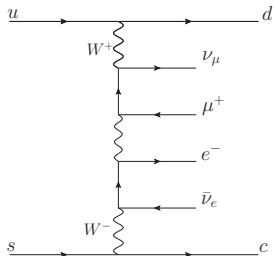
NLO QCD corrections - results for $qq \rightarrow qqWW$



Jäger, Oleari, Zeppenfeld (2006)

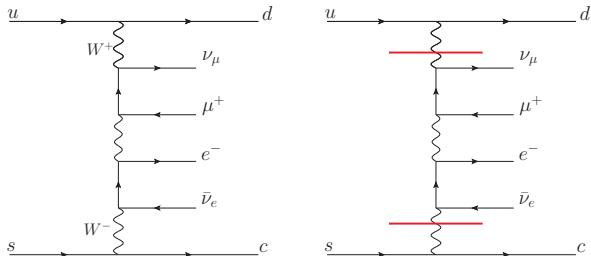
Block structure

- EW parts are not affected by QCD corrections and only have to be evaluated once
- separating QCD and EW blocks
 - simplifies calculation
 - speeds up Monte Carlo simulations



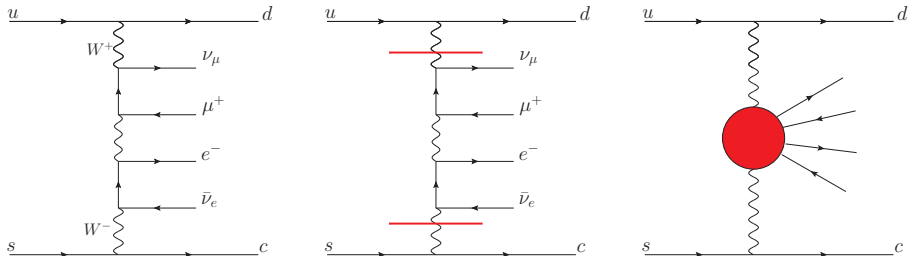
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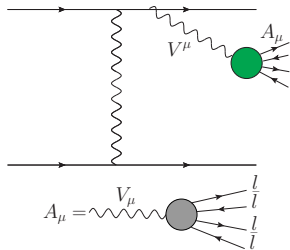


How it works

- example - \mathcal{A} -block:

$$\mathcal{M} = \mathcal{M}_{QCD\mu} A^\mu = \mathcal{M}_{QCD\mu} g^{\mu\nu} A_\nu \quad \text{and} \quad g_{\mu\nu} = - \sum_i \varepsilon(k)_{i\mu} \varepsilon(k)_{i\nu} + \frac{k_\mu k_\nu}{k^2}$$

$$\begin{aligned} \Rightarrow \quad \mathcal{M} &= -(\mathcal{M}_{QCD} \cdot \varepsilon_+) (\mathcal{A} \cdot \varepsilon_+) - (\mathcal{M}_{QCD} \cdot \varepsilon_-) (\mathcal{A} \cdot \varepsilon_-) \\ &\quad - (\mathcal{M}_{QCD} \cdot \varepsilon_0) (\mathcal{A} \cdot \varepsilon_0) + \frac{1}{k^2} (\mathcal{M}_{QCD} \cdot k) (\mathcal{A} \cdot k) \end{aligned}$$



How it works

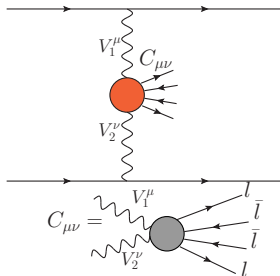
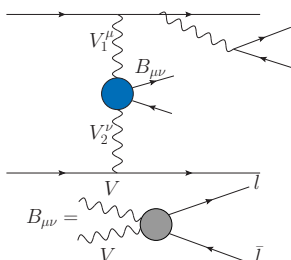
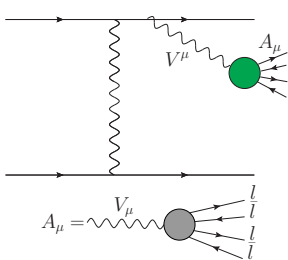
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$$\mathcal{M} = -(\mathcal{M}_{QCD} \cdot \varepsilon_+) (\mathcal{A} \cdot \varepsilon_+) - (\mathcal{M}_{QCD} \cdot \varepsilon_-) (\mathcal{A} \cdot \varepsilon_-)$$

\Rightarrow

$$-(\mathcal{M}_{QCD} \cdot \varepsilon_0) (\mathcal{A} \cdot \varepsilon_0) + \frac{1}{k^2} (\mathcal{M}_{QCD} \cdot k) (\mathcal{A} \cdot k)$$



Elements of Calculation

- diagrams generated with FeynArts
 - EW and QCD blocks are generated independently, modified model used to accommodate the modular structure
- analytical expressions generated with FormCalc and modified within Mathematica
- Weyl-van-der-Waerden formalism - translates all kinematic objects into two-component WvdW spinors in chiral representation

$$\Psi = \begin{pmatrix} \phi_A \\ \psi^{\dot{A}} \end{pmatrix} \quad \psi_A \phi^A = (\psi \phi) \quad \psi_{\dot{A}} \phi^{\dot{A}} = \langle \psi \phi \rangle$$

$$k_\mu \sigma_{\dot{A}B}^\mu = \sum_{i=1,2} \lambda_i n_{i\dot{A}} n_{iB} \quad 2k_\mu p^\mu = (k p) \langle k p \rangle \quad \epsilon_{\mu+}(k) \sigma_{\dot{A}B}^\mu = \sqrt{2} n_{2,\dot{A}} n_{1,B}$$

1. introducing WvdW spinors

$$\bar{\Psi}(k_2) P_R \gamma^\mu \Psi(k_3) \mapsto k_2^a \sigma_{ab}^\mu k_3^b$$

2. open Lorentz indices contracted (polarization sum)

$$k_2^a \sigma_{ab}^\mu k_3^b \mapsto k_2^a \varepsilon_{\mu+}(k_1) \sigma_{ab}^\mu k_3^b$$

3. scalar products of individual spinors

$$k_2^a \varepsilon_{\mu+}(k_1) \sigma_{ab}^\mu k_3^b \mapsto 2\sqrt{2k_0(2)k_0(3)} \langle n_2(1)n_1(3) \rangle (n_1(1)n_1(2))$$

4. full forms of spinors are introduced

$$(n_1(1)n_1(2)) = \left(\sin \frac{\theta(1)}{2}, -e^{-i\Phi(1)} \cos \frac{\theta(1)}{2} \right) \cdot \begin{pmatrix} e^{-i\Phi(2)} \cos \frac{\theta(2)}{2} \\ \sin \frac{\theta(2)}{2} \end{pmatrix}$$

5. introducing abbreviations reduces the size of expressions

6. a fully self-contained Fortran program is generated

- Fortran program
 - modular approach (using abbreviations)
 - produces numerical values for $\mathcal{M}_{QCD}^\mu \varepsilon_{i\mu}$ and $(\mathcal{A}.\varepsilon_i)$, $(\mathcal{B}.\varepsilon_i)$, $(\mathcal{C}.\varepsilon_i)$
 - independent EW and QCD blocks are "sewn" together
 - input - external momenta and helicities
- outlook
 - verification of the Fortran code by comparison with FormCalc and MadGraph results
 - full QCD NLO corrections with FormCalc - first independent check of the calculation by Zeppenfeld and his group
- development of Monte Carlo program in progress

Summary

- vector boson scattering is essential for
 - search for a light Higgs at the LHC
 - probing for new strong effects in case of no Higgs is found
- long and rich history of studies
- NLO corrections are fairly complicated
 - requires lot of computer power
 - new methods and strategies to cope with the large number of diagrams
- my work
 - constructed the building blocks in Mathematica + FormCalc
 - implemented modular structure in Fortran
 - verification of the code in progress
- future plan - EW corrections