CP violating asymmetries induced by supersymmetry

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PSI Particle Theory Seminar September 2008

Outline



- 2 Introduction
 - Supersymmetry (SUSY)
 - The Minimal Supersymmetric Standard Model (MSSM)

OP violating decay rate asymmetry

- Introduction
- Contributions
- Numerical results

Summary

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Baryon asymmetry of the universe

- Exists much more baryonic matter than anti-matter
- Standard Model (SM) cannot explain baryon asymmetry of the universe!
- Evidence from acoustic peaks (early universe baryon-photon plasma oscillations) deduced from Cosmic Microwave Background measurements

Baryon-to-photon ratio

$$\eta \equiv \frac{n_B}{s} \equiv \frac{n_b - n_{\bar{b}}}{s} = (6.1^{+0.3}_{-0.2}) \times 10^{-10}$$

 $s \dots$ entropy density (roughly photon density) $n_b(n_{\overline{b}}) \dots$ number densities of baryons (anti-baryons)

Baryogenesis

Problem

How obtains η this small value from initial condition $\eta = 0$?

Criteria of a solution

Three necessary conditions for baryogenesis: Sakharov requirements

- Baryon number violation
- 2 Departure from thermal equilibrium
- Oharge (C) and Charge-Parity (CP) violation

SM can meet Sakharov criteria but baryon asymmetry is *too small*!

Possible solution

- Supersymmetric extensions of SM can contain new sources of CP violation
- Lead to increase and thus possible explanation of baryon asymmetry
- Special case: Minimal Supersymmetric Standard Model (MSSM) introduces new parameters
- If some parameters are chosen *complex*, radiative corrections at one-loop can lead to *new CP violating asymmetries*

Supersymmetry (SUSY) The Minimal Supersymmetric Standard Model (MSSM)

Supersymmetry (SUSY)

Main idea

Symmetry between bosons and fermions

- Superpartner: same quantum numbers and mass but different spin
- Simplest case: superpartner for every particle of the SM
- Until now no superpartners found
 - Have a high mass
 - Explainable with spontaneous breaking of supersymmetry

Supersymmetry (SUSY) The Minimal Supersymmetric Standard Model (MSSM)

Particle content of the MSSM

- All particles obtain a superpartner
- Particle and superpartner form supermultiplets
- More than one Higgs particle
- Naming convention
 - Prefix 's' for spin = 0 superpartner (sleptons, squarks, ...)
 - suffix '-ino' for spin = ¹/₂ superpartner (gluino, Higgsinos ...)
- Superpartner mix because of electroweak symmetry breaking

Supersymmetry (SUSY) The Minimal Supersymmetric Standard Model (MSSM)

Particle content of the MSSM Chiral supermultiplets

Names Spin 0 Spin 1/2 $SU(3)_C, SU(2)_L, U(1)_Y$ Õ $(\widetilde{u}_L \ \widetilde{d}_L)$ $(3, 2, \frac{1}{6})$ squarks, quarks $(u_L \ d_L)$ $(\overline{3}, 1, -\frac{2}{2})$ u_{R}^{\dagger} $(\times 3 \text{ families})$ \widetilde{U}_{R}^{*} īī d \widetilde{d}_{R}^{*} $(\bar{\bf 3}, {\bf 1}, \frac{1}{3})$ d_{R}^{\dagger} ĩ $(\widetilde{\nu} \ \widetilde{e}_L)$ $(1, 2, -\frac{1}{2})$ sleptons, leptons (νe_L) (×3 families) ē $\widetilde{\boldsymbol{e}}_{\boldsymbol{B}}^{*}$ e_{P}^{\dagger} **(1, 1, 1)** Higgs, Higgsinos H_{μ} (H_2^1) H_2^2) (H_2^1) H_2^2) $(1, 2, +\frac{1}{2})$ H_{1}^{2}) (\widetilde{H}_1^1) \widetilde{H}_{1}^{2}) H_d (H_{1}^{1}) $(1, 2, -\frac{1}{2})$

Spin-0 fields are complex scalars Spin-1/2 fields are left-handed two-component Weyl fermions

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Supersymmetry (SUSY) The Minimal Supersymmetric Standard Model (MSSM)

Particle content of the MSSM

Gauge supermultiplets & mixing

Names	Spin 1/2	Spin 1	$SU(3)_{C}, SU(2)_{L}, U(1)_{Y}$	
gluino, gluon	ĝ	g	(8 , 1 , 0)	
winos, W bosons	$\widetilde{\lambda}^{\pm}$ $\widetilde{\lambda}^{3}$	$W^{\pm} W^{0}$	(1 , 3 , 0)	
bino, B boson	$\widetilde{\lambda}'$	B^0	(1 , 1 , 0)	

• Remember:

Interaction eigenstates no longer mass eigenstates (physical particles)!

• Needs to consider mixing of interaction eigenstates to mass eigenstates

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Supersymmetry (SUSY) The Minimal Supersymmetric Standard Model (MSSM)

Particle content of the MSSM

Mixing of eigenstates

Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
Higgs bosons	0	+1	$H_2^0 H_1^0 H_1^+ H_1^-$	$h^0 H^0 A^0 H^{\pm}$
			$\widetilde{u}_L \ \widetilde{u}_R \ \widetilde{d}_L \ \widetilde{d}_R$	$\widetilde{u}_1 \ \widetilde{u}_2 \ \widetilde{d}_1 \ \widetilde{d}_2$
squarks	0	-1	$\widetilde{C}_L \ \widetilde{C}_R \ \widetilde{S}_L \ \widetilde{S}_R$	$\widetilde{C}_1 \ \widetilde{C}_2 \ \widetilde{S}_1 \ \widetilde{S}_2$
			$\widetilde{t}_L \widetilde{t}_R \widetilde{b}_L \widetilde{b}_R$	$\widetilde{t}_1 \ \widetilde{t}_2 \ \widetilde{b}_1 \ \widetilde{b}_2$
			$\widetilde{\mathbf{e}}_L \ \widetilde{\mathbf{e}}_R \ \widetilde{\mathbf{ u}}_e$	$\widetilde{e}_1 \ \widetilde{e}_2 \ \widetilde{\nu}_e$
sleptons	0	-1	$\widetilde{\mu}_L \; \widetilde{\mu}_{R} \; \widetilde{ u}_{\mu}$	$\widetilde{\mu}_1$ $\widetilde{\mu}_2$ $\widetilde{ u}_\mu$
			$\widetilde{ au}_L \ \widetilde{ au}_R \ \widetilde{ uu}_ au$	$\widetilde{ au}_1$ $\widetilde{ au}_2$ $\widetilde{ au}_{ au}$
neutralinos	1/2	-1	$\widetilde{\lambda}' \widetilde{\lambda}^3 \widetilde{H}_2^2 \widetilde{H}_1^1$	$\widetilde{\chi}^0_1 \ \widetilde{\chi}^0_2 \ \widetilde{\chi}^0_3 \ \widetilde{\chi}^0_4$
charginos	1/2	-1	$\widetilde{\lambda}^{\pm}$ \widetilde{H}_{2}^{1} \widetilde{H}_{1}^{2}	$\widetilde{\chi}_1^{\pm}$ $\widetilde{\chi}_2^{\pm}$
gluino	1/2	-1	\widetilde{g}	\widetilde{g}

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Introduction Contributions Numerical results

Definitions

- C transformation changes sign of charge, changes particle to its anti-particle
- P transformation changes sign of coordinate system, turns right-handed coordinate system into a left-handed one or vice versa $P: \psi(t, \vec{x}) \rightarrow \psi(t, -\vec{x})$
- CP transformation applied to Lagrangian changes
 - Signs of momentums and charges
 - Left- and right-handed parts
 - Conjugates coupling matrices
- If coupling matrices are *complex*, CP violation can occur!

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Introduction Contributions Numerical results

Definitions

- Decay rate asymmetry $\delta^{CP} = \frac{\Gamma^+ \Gamma^-}{\Gamma^+ + \Gamma^-}$
- Γ^{+,-} decay rate of a process (regular, CP transformed)
- Special case: $\Gamma^+ = \Gamma(\tilde{t}_i \to b \, \tilde{\chi}_k^+), \, \Gamma^- = \Gamma(\tilde{t}_i^* \to \bar{b} \, \tilde{\chi}_k^{+c})$
- Branching ratio $BR = \frac{\Gamma}{\Gamma_{\text{total}}}$
- δ^{CP} × BR probability how often certain decay is CP violated (compared to rest)

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Decay rate asymmetry δ^{CP}

Decay rate asymmetry δ^{CP} is \neq 0 only if

- including radiative corrections with at least one-loop (both CP transformation and calculating adjoint matrix M[↑] (in Γ ∝ ∑_s M[↑]M) conjugate tree-level couplings -> at tree-level net effect zero)
- coupling between particles complex (due to complex MSSM parameters) (CP transformation conjugates tree-level couplings)
- at least a second decay channel open (i.e. in addition to t
 *˜*_i → b X
 *˜*_k another like t
 *˜*_i → t *̃*g) (From CPT Theorem: only total decay width Γ_{total} CP invariant, but r necessarily partial decay widths!)

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Decay rate asymmetry δ^{CP}

Decay rate asymmetry δ^{CP} is \neq 0 only if

- including radiative corrections with at least one-loop (both CP transformation and calculating adjoint matrix \mathcal{M}^{\dagger} (in $\Gamma \propto \sum_{s} \mathcal{M}^{\dagger} \mathcal{M}$) conjugate tree-level couplings -> at tree-level net effect zero)
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- at least a second decay channel open
 (i.e. in addition to t
 *˜*_i → b χ
 *˜*_k another like t
 *˜*_i → t g
 (From CPT Theorem: only total decay width Γ_{total} CP invariant, but no necessarily partial decay widths!)

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- coupling between particles complex (due to complex MSSM parameters) (CP transformation conjugates tree-level couplings)

3 at least a second decay channel open (i.e. in addition to $\tilde{t}_i \rightarrow b \, \tilde{\chi}_k^+$ another like $\tilde{t}_i \rightarrow t \, \tilde{g}$) (From CPT Theorem: only total decay width Γ_{total} CP invariant, but

necessarily partial decay widths!)

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Introduction Contributions Numerical results

Decay rate asymmetry δ^{CP}

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Introduction Contributions Numerical results

Most important contributions (expected) All gluino contributions



- Gluino \tilde{g} couples like gluon g with strong interaction force
- If decay t̃_i → t g̃ kinematically possible, these contributions should dominate over all others

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All vertex contributions



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All stop-selfenergy contributions



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All chargino-selfenergy contributions



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Calculation

- Derived most important contributions analytically and numerically by myself (including derivation of mass matrices, couplings, generic structures etc)
- Calculation of full one-loop corrections with FeynArts/FormCalc/LoopTools (successfully checked with own computations)
- Input parameters of SM and MSSM for a typical scenario
- Complex parameters can violate experimental limit of electric dipole moment of electron
 - Checked automatically with self-written routine

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Some input parameters of SM and MSSM

- Coupling of strong interaction force α_s is taken running in dimensional reduction regularization scheme DR, renormalized at scale of decaying particle (SPA convention)
- Gluino mass $m_{\tilde{g}}$ calculated from α_s via GUT relations
- Yukawa couplings of third generation (s)quarks (*h_t*, *h_b*) are taken running
- SUSY breaking mass parameters M_{Q̃}, M_ũ, M_{d̃}, M_{l̃} and M_ē set equal in all generations
- Trilinear breaking parameters of 1st and 2nd generation set to zero $(A_{u,d,e} = A_{c,s,\mu} = 0)$
- We further simplify and set |M₁| = M₂/2 (GUT relation),
 M_{Q̃} = M_ũ = M_{d̃}, M_{L̃} = M_ẽ, |A_t| = |A_b| = |A_τ| and φ_{At} = φ_{Ab} = φ_{Aτ}
- Due to stringent experimental constraints, we set $\varphi_{\mu} = 0$ (phase of Higgsino mass) and focus just on phase of A_f (the phase of SUSY breaking gaugino mass M_1 is negligible in our case)

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Some numerical results $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^+$ (All contributions)



- Actually vary input parameter M_{Q̃} (SUSY breaking mass) from 500 to 1500 GeV but show output parameter m_{t̃1} for convenience
- $\tan \beta$ is ratio of the two VEVs of Higgs fields
- Threshold of decay $ilde{t}_1
 ightarrow t\, ilde{g}$ at $m_{ ilde{t}_1}\sim$ 708 GeV
- Dominance of gluino contributions over all insignificant others

•
$$\delta^{CP}_{
m max}\sim$$
 22 %, $(\delta^{CP} imes BR)_{
m max}\sim$ 2.4 %

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Some numerical results $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^+$ (Gluino-to-all ratio)



• After threshold at $m_{\tilde{t}_1} \sim$ 708 GeV gluino processes account for \sim 98 % of all processes

• Kink at $m_{\tilde{t}_1} \sim 1175$ GeV comes from threshold of $\tilde{t}_1 \rightarrow t \, \tilde{\chi}_3^0$

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Some numerical results $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^+$ (Comparison of gluino contributions)

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- (1) Gluino in selfenergy-loop,
 (2) Gluino in vertex correction
- Contrary to expectation only one gluino contribution dominates!
- Reason lies in coupling *b_j* − *t* − *x̃*⁺₁ embedded in vertex correction, however no simple explanation possible *c* ≥ *x* ∈

Introduction Contributions Numerical results

Some numerical results $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^+$ (All contributions)



- Complex phase of breaking parameter A_t highly influences δ^{CP} (only source of CP violation in our scenario)
- Overall maximum of scenario: $\delta^{CP}_{max} \sim$ 24 %

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Introduction Contributions Numerical results

Some numerical results $\tilde{t}_1 \rightarrow b \tilde{\chi}_1^+$ (All contributions)



- Closure of dominating decay channel t
 ₁ → t g
 _i, because SUSY breaking gaugino mass M₂ is related to gluino mass m<sub>g
 _j
 </sub>
- The higher mass $M_{\tilde{Q}}(m_{\tilde{t}_1})$, the later this closure happens
- Overall maximum of scenario: $\delta^{CP} imes BR \sim 3.5$ %

Summary

- Baryogenesis needs CP violation exceeding the one in SM
- Minimal Supersymmetric Standard Model (MSSM) with complex parameters leads to new CP violation
- One-loop corrections to *t
 _i*→ *b χ*⁺_k leads to CP violating decay rate asymmetry δ^{CP}
- Detailed numerical analysis of δ^{CP} and $\delta^{CP} \times BR$
- δ^{CP} rises up to \sim 24 %, $\delta^{CP} \times BR$ up to \sim 3.5 %
- Asymmetry δ^{CP} will be measurable at LHC

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