

Searches for Ultra-Low-Mass Dark Matter

Yevgeny Stadnik

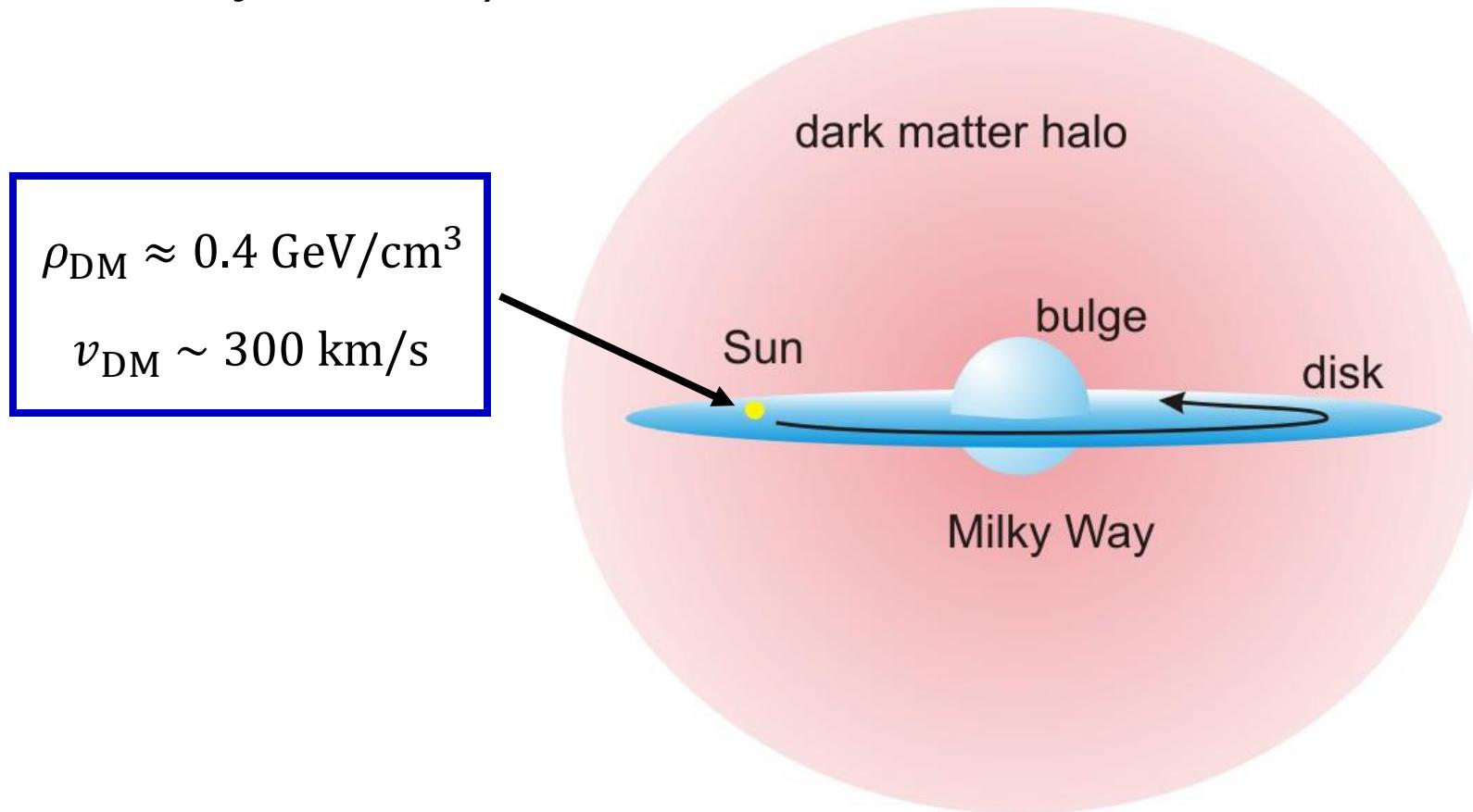
Kavli IPMU, University of Tokyo, Japan



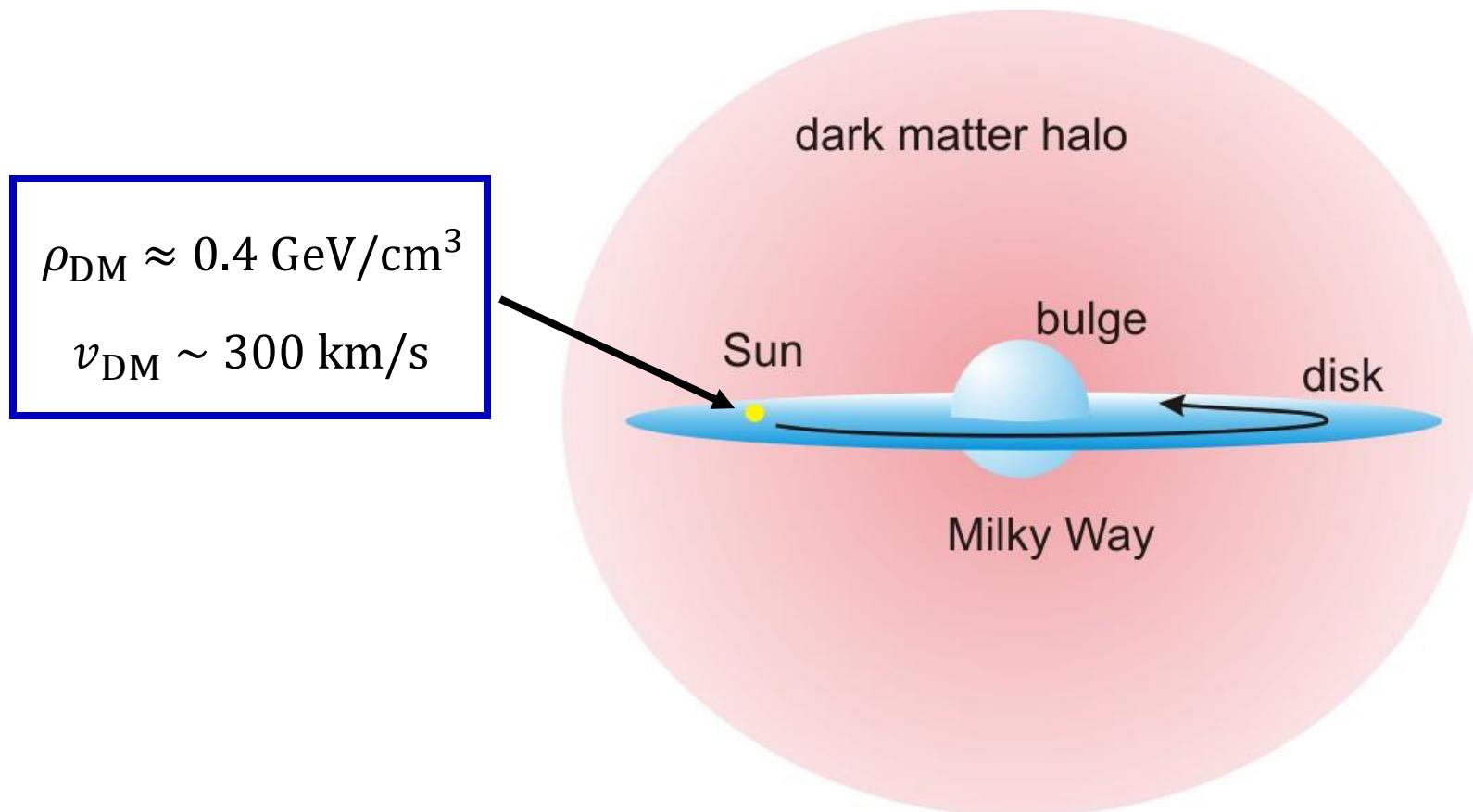
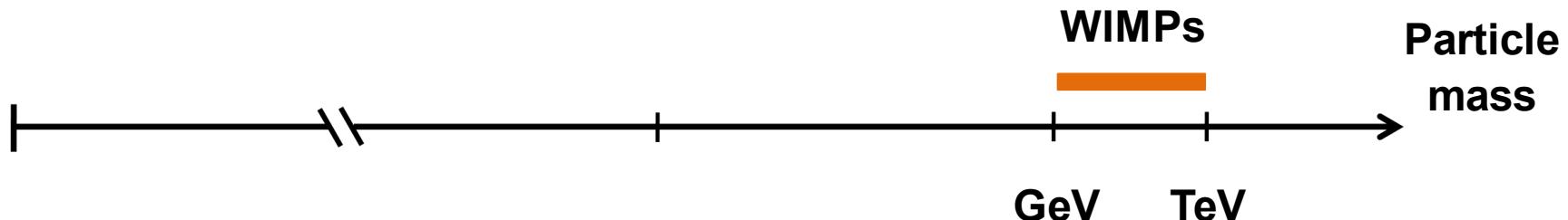
Virtual Colloquium, PSI, March 2021

Dark Matter

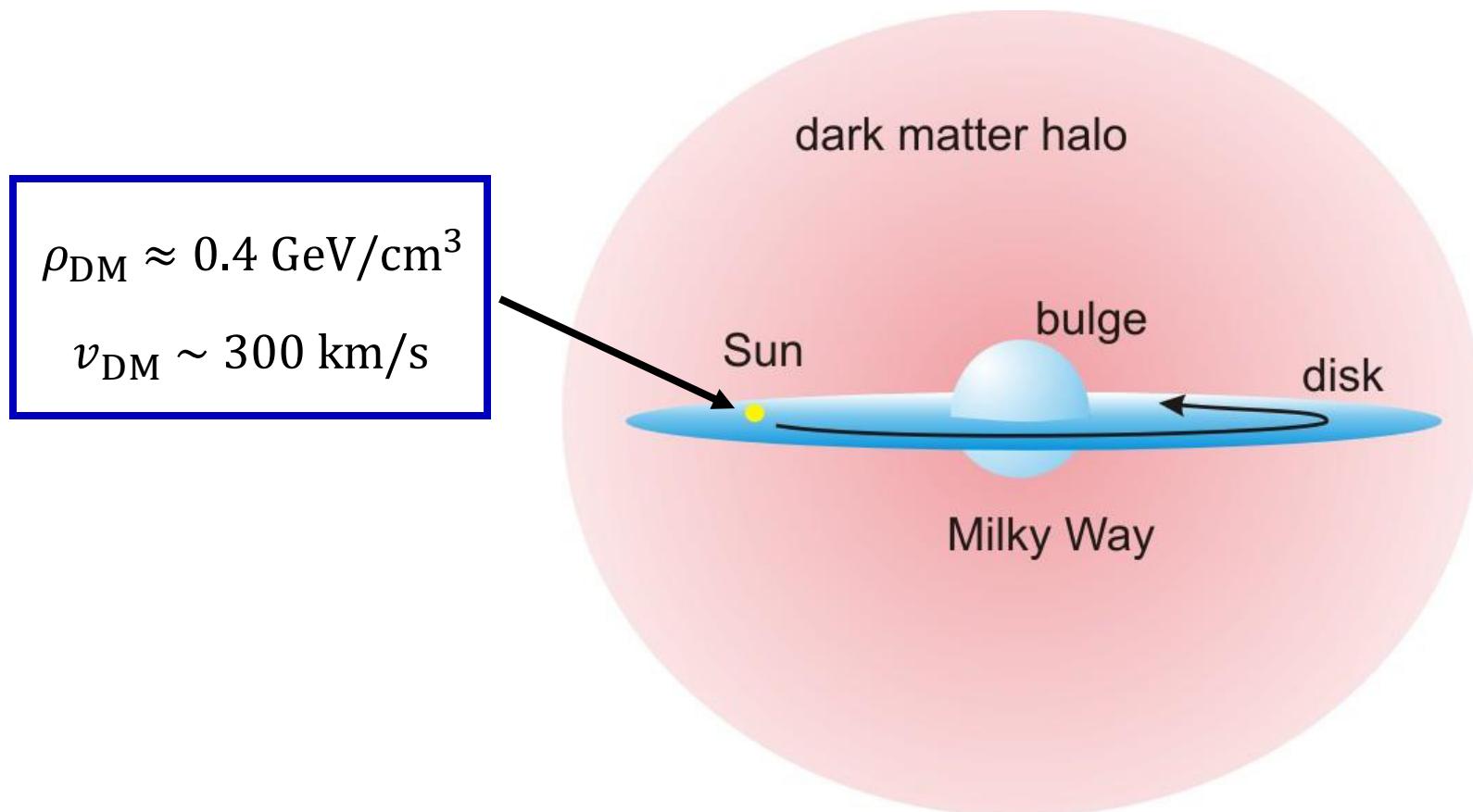
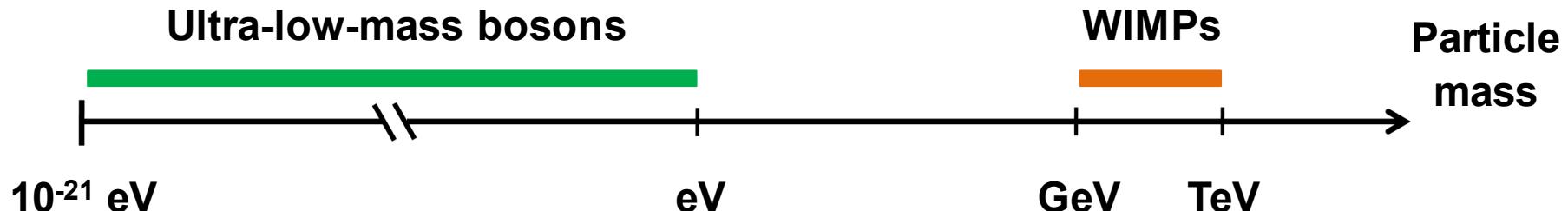
Strong astrophysical evidence for existence of **dark matter** (~5 times more dark matter than ordinary matter)



Dark Matter

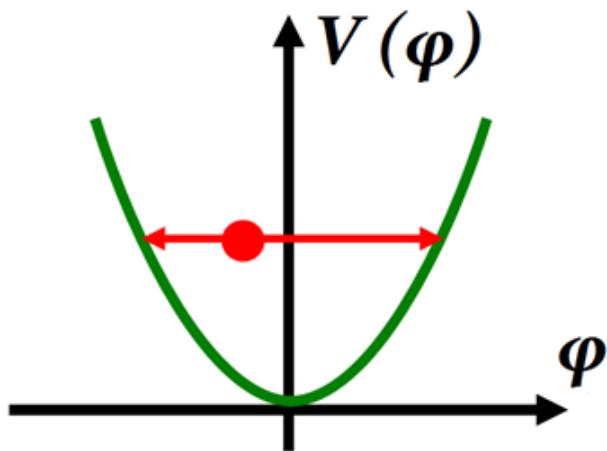


Dark Matter



Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)



$$V(\varphi) = \frac{m_\varphi^2 \varphi^2}{2}$$

$$\ddot{\varphi} + m_\varphi^2 \varphi \approx 0$$

Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- *Coherently* oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)

Low-mass Spin-0 Dark Matter

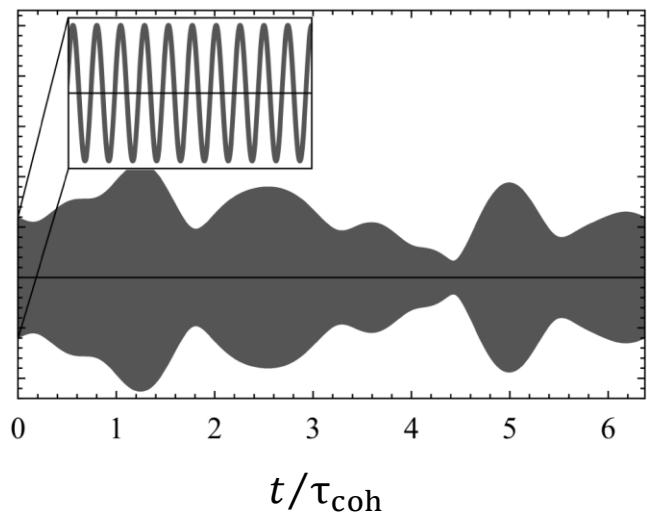
- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- *Coherently* oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- $\Delta E_\varphi / E_\varphi \sim \langle v_\varphi^2 \rangle / c^2 \sim 10^{-6} \Rightarrow \tau_{\text{coh}} \sim 2\pi / \Delta E_\varphi \sim 10^6 T_{\text{osc}}$

 $v_{\text{DM}} \sim 300 \text{ km/s}$

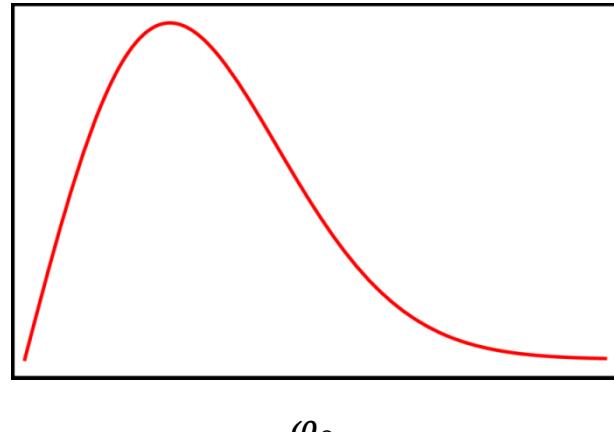
Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- *Coherently* oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- $\Delta E_\varphi / E_\varphi \sim \langle v_\varphi^2 \rangle / c^2 \sim 10^{-6} \Rightarrow \tau_{\text{coh}} \sim 2\pi / \Delta E_\varphi \sim 10^6 T_{\text{osc}}$

Evolution of φ_0 with time



Probability distribution function of φ_0
(e.g., Rayleigh distribution)



Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- *Coherently* oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- $\Delta E_\varphi / E_\varphi \sim \langle v_\varphi^2 \rangle / c^2 \sim 10^{-6} \Rightarrow \tau_{\text{coh}} \sim 2\pi / \Delta E_\varphi \sim 10^6 T_{\text{osc}}$
- *Classical* field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$

Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- *Coherently* oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- $\Delta E_\varphi / E_\varphi \sim \langle v_\varphi^2 \rangle / c^2 \sim 10^{-6} \Rightarrow \tau_{\text{coh}} \sim 2\pi / \Delta E_\varphi \sim 10^6 T_{\text{osc}}$
- *Classical* field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$
- $10^{-21} \text{ eV} \lesssim m_\varphi \lesssim 1 \text{ eV} \Leftrightarrow 10^{-7} \text{ Hz} \lesssim f_{\text{DM}} \lesssim 10^{14} \text{ eV}$
 $T_{\text{osc}} \sim 1 \text{ month}$ **IR frequencies**

Lyman- α forest measurements [suppression of structures for $L \lesssim \mathcal{O}(\lambda_{\text{dB},\varphi})$]

[Related figure-of-merit: $\lambda_{\text{dB},\varphi} / 2\pi \leq L_{\text{dwarf galaxy}} \sim 100 \text{ pc} \Rightarrow m_\varphi \gtrsim 10^{-21} \text{ eV}$]

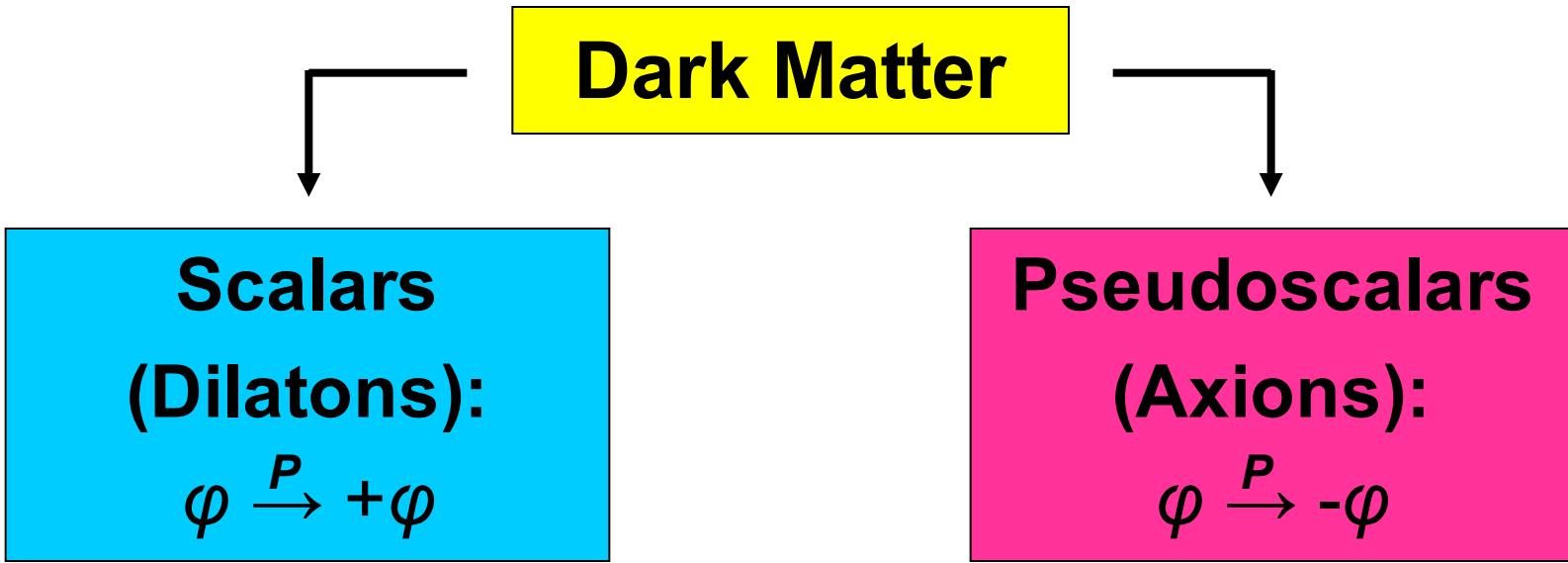
Low-mass Spin-0 Dark Matter

- Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- *Coherently* oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- $\Delta E_\varphi / E_\varphi \sim \langle v_\varphi^2 \rangle / c^2 \sim 10^{-6} \Rightarrow \tau_{\text{coh}} \sim 2\pi / \Delta E_\varphi \sim 10^6 T_{\text{osc}}$
- *Classical* field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$
- $10^{-21} \text{ eV} \lesssim m_\varphi \lesssim 1 \text{ eV} \Leftrightarrow 10^{-7} \text{ Hz} \lesssim f_{\text{DM}} \lesssim 10^{14} \text{ eV}$

- *Wave-like signatures* [cf. *particle-like* signatures of WIMP DM]

Lyman- α forest measurements [suppression of structures for $L \lesssim \mathcal{O}(\lambda_{\text{dB},\varphi})$]

Low-mass Spin-0 Dark Matter



→ Time-varying

fundamental constants

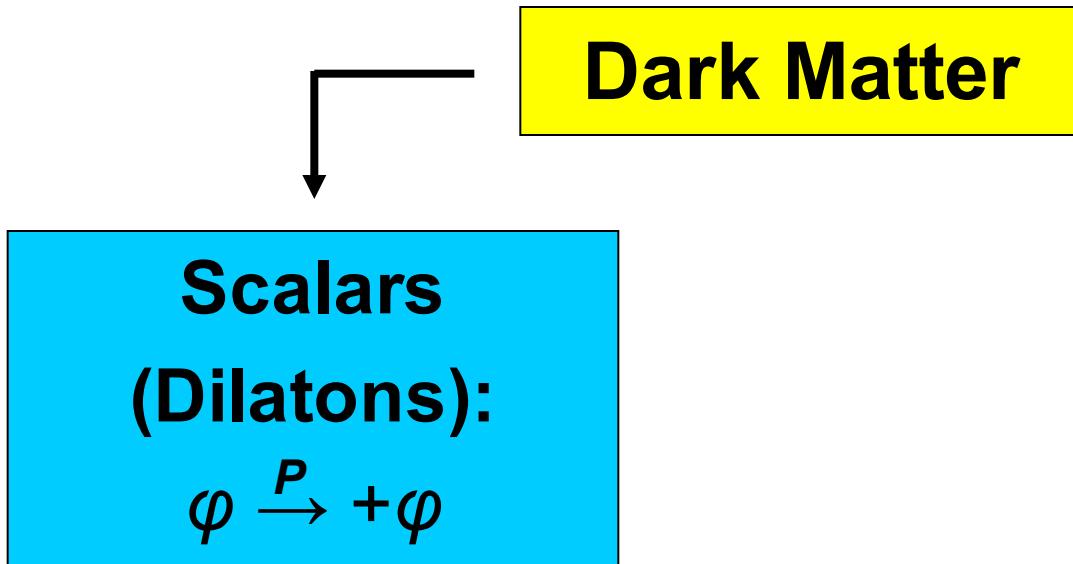
- Atomic clocks
- Cavities and interferometers
 - Torsion pendula
- Astrophysics (e.g., BBN)

→ Time-varying spin-

dependent effects

- Co-magnetometers
- Particle g -factors
- Spin-polarised torsion pendula
- Spin resonance (NMR, ESR)

Low-mass Spin-0 Dark Matter



→ Time-varying

fundamental constants

- Atomic clocks
- Cavities and interferometers
 - Torsion pendula
- Astrophysics (e.g., BBN)

Dark-Matter-Induced Variations of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],
[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_\gamma = \frac{\varphi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \frac{\delta\alpha}{\alpha} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma}$$

Dark-Matter-Induced Variations of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],
[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_\gamma = \frac{\varphi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \frac{\delta\alpha}{\alpha} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma}$$
$$\mathcal{L}_f = -\frac{\varphi}{\Lambda_f} m_f \bar{f} f \approx -\frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} m_f \bar{f} f \Rightarrow \frac{\delta m_f}{m_f} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f}$$

Dark-Matter-Induced Variations of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, PRD 98, 064051 (2018)]

$$\mathcal{L}_\gamma = \frac{\varphi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \frac{\delta \alpha}{\alpha} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma}$$

$$\mathcal{L}_f = -\frac{\varphi}{\Lambda_f} m_f \bar{f} f \approx -\frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} m_f \bar{f} f \Rightarrow \frac{\delta m_f}{m_f} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f}$$

$$\varphi = \varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x}) \Rightarrow \mathbf{F} \propto \mathbf{p}_\varphi \sin(m_\varphi t)$$

Lab frame

Solar System (and lab) move through stationary dark matter halo

Dark-Matter-Induced Variations of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],
 [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\begin{aligned} \mathcal{L}_\gamma &= \frac{\varphi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \frac{\delta\alpha}{\alpha} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \\ \mathcal{L}_f &= -\frac{\varphi}{\Lambda_f} m_f \bar{f} f \approx -\frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} m_f \bar{f} f \Rightarrow \frac{\delta m_f}{m_f} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} \end{aligned}$$

$$\varphi = \varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x}) \Rightarrow F \propto \mathbf{p}_\varphi \sin(m_\varphi t)$$

$$\left. \begin{aligned} \mathcal{L}'_\gamma &= \frac{\varphi^2}{(\Lambda'_\gamma)^2} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \\ \mathcal{L}'_f &= -\frac{\varphi^2}{(\Lambda'_f)^2} m_f \bar{f} f \end{aligned} \right\}$$

φ^2 interactions also exhibit the same oscillating-in-time signatures as above, as well as ...

Dark-Matter-Induced Variations of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],
 [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\begin{aligned} \mathcal{L}_\gamma &= \frac{\varphi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \frac{\delta\alpha}{\alpha} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_\gamma} \\ \mathcal{L}_f &= -\frac{\varphi}{\Lambda_f} m_f \bar{f} f \approx -\frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} m_f \bar{f} f \Rightarrow \frac{\delta m_f}{m_f} \approx \frac{\varphi_0 \cos(m_\varphi t)}{\Lambda_f} \end{aligned}$$

$$\varphi = \varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x}) \Rightarrow \mathbf{F} \propto \mathbf{p}_\varphi \sin(m_\varphi t)$$

$$\left. \begin{aligned} \mathcal{L}'_\gamma &= \frac{\varphi^2}{(\Lambda'_\gamma)^2} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \\ \mathcal{L}'_f &= -\frac{\varphi^2}{(\Lambda'_f)^2} m_f \bar{f} f \end{aligned} \right\} \Rightarrow \left\{ \begin{array}{l} \frac{\delta\alpha}{\alpha} \propto \frac{\delta m_f}{m_f} \propto \Delta\rho_\varphi \\ \mathbf{F} \propto \nabla\rho_\varphi \end{array} \right.$$

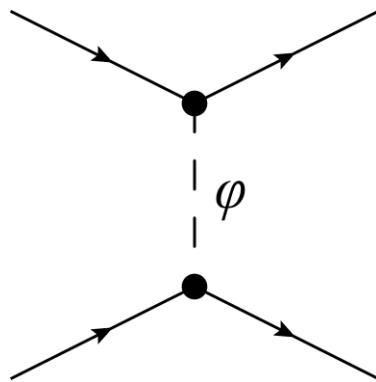
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Linear couplings ($\varphi \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



$$\varphi = \varphi_0 \cos(m_\varphi t) \pm A \frac{e^{-m_\varphi r}}{r}$$

Profile outside of a spherical body

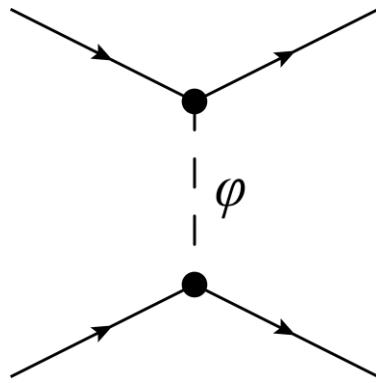
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

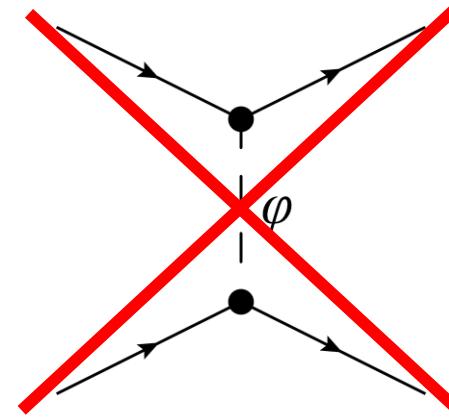
Linear couplings ($\varphi \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



Quadratic couplings ($\varphi^2 \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa' \rho \varphi \quad \text{Potential term}$$



$$\varphi = \varphi_0 \cos(m_\varphi t) \pm A \frac{e^{-m_\varphi r}}{r}$$

Profile outside of a spherical body

$$m_{\text{eff}}^2(\rho) = m_\varphi^2 \mp \kappa' \rho$$

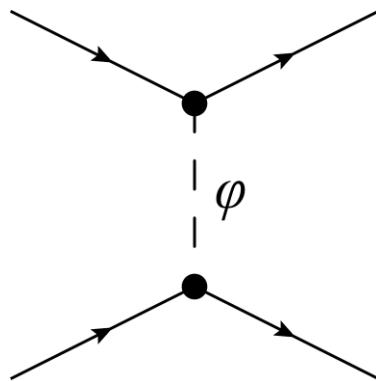
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

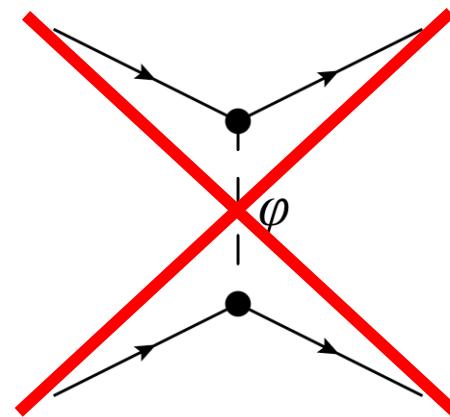
Linear couplings ($\varphi \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



Quadratic couplings ($\varphi^2 \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa' \rho \varphi \quad \text{Potential term}$$

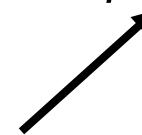


$$\varphi = \varphi_0 \cos(m_\varphi t) \pm A \frac{e^{-m_\varphi r}}{r}$$



Profile outside of a spherical body

$$\varphi = \varphi_0 \cos(m_\varphi t) \left(1 \pm \frac{B}{r} \right)$$



Gradients + amplification/screening

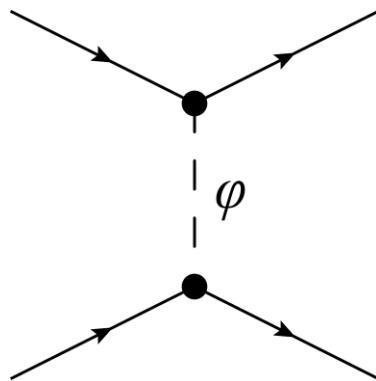
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

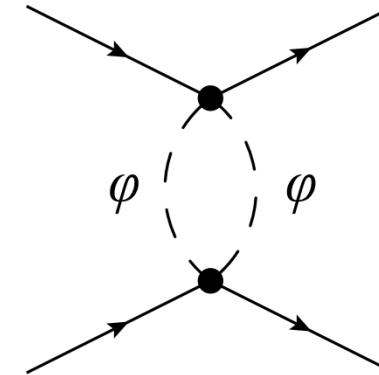
Linear couplings ($\varphi \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



Quadratic couplings ($\varphi^2 \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa' \rho \varphi \quad \text{Potential term}$$

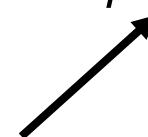


$$\varphi = \varphi_0 \cos(m_\varphi t) \pm A \frac{e^{-m_\varphi r}}{r}$$



Profile outside of a spherical body

$$\varphi = \varphi_0 \cos(m_\varphi t) \left(1 \pm \frac{B}{r} \right) - C \frac{e^{-2m_\varphi r}}{r^3}$$



Gradients + amplification/screening

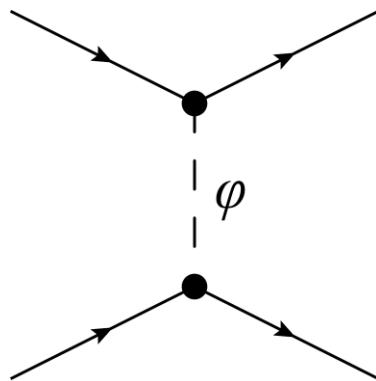
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

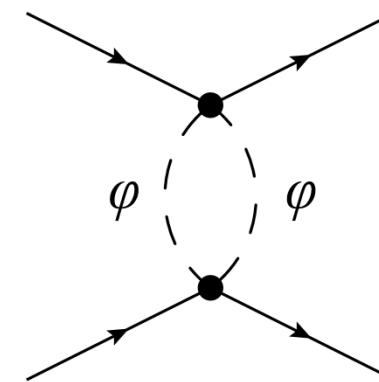
Linear couplings ($\varphi \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



Quadratic couplings ($\varphi^2 \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa' \rho \varphi \quad \text{Potential term}$$



$$\varphi = \frac{\varphi_0 \cos(m_\varphi t)}{\pm A} \frac{e^{-m_\varphi r}}{r}$$

Motional gradients: $\varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x})$

$$\varphi = \frac{\varphi_0 \cos(m_\varphi t)}{\pm B} \left(1 - \frac{C}{r^3} e^{-2m_\varphi r} \right)$$



Gradients + amplification/screening

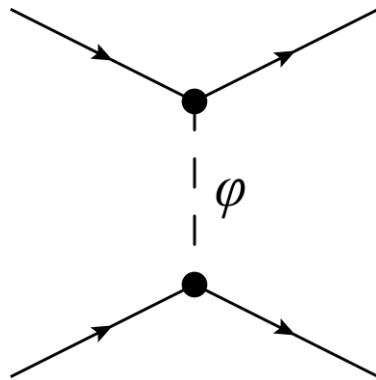
Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

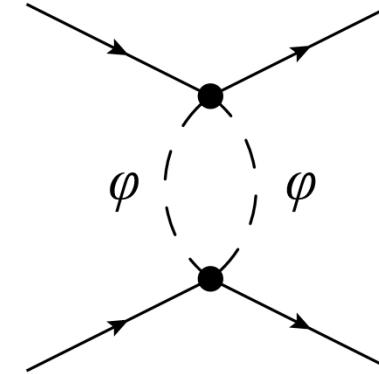
Linear couplings ($\varphi \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa \rho \quad \text{Source term}$$



Quadratic couplings ($\varphi^2 \bar{X} X$)

$$\square\varphi + m_\varphi^2 \varphi = \pm \kappa' \rho \varphi \quad \text{Potential term}$$



$$\varphi = \frac{\varphi_0 \cos(m_\varphi t)}{\pm A} \frac{e^{-m_\varphi r}}{r}$$

Motional gradients: $\varphi_0 \cos(m_\varphi t - \mathbf{p}_\varphi \cdot \mathbf{x})$

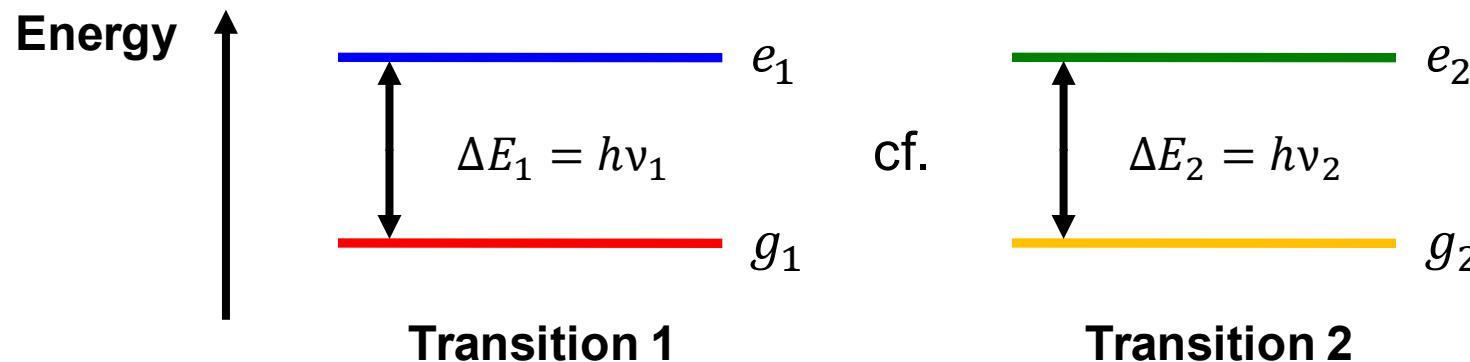
“Fifth-force” experiments: torsion pendula, atom interferometry

$$\varphi = \frac{\varphi_0 \cos(m_\varphi t)}{\pm B} \left(1 - \frac{C}{r^3} \right) e^{-2m_\varphi r}$$



Gradients + amplification/screening

Atomic Spectroscopy Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter



$$\frac{\delta(v_1/v_2)}{v_1/v_2} = (K_{X,1} - K_{X,2}) \frac{\delta X}{X} ; \quad X = \alpha, m_e/m_N, \dots$$

Atomic spectroscopy (including clocks) has been used for decades to search for “slow drifts” in fundamental constants

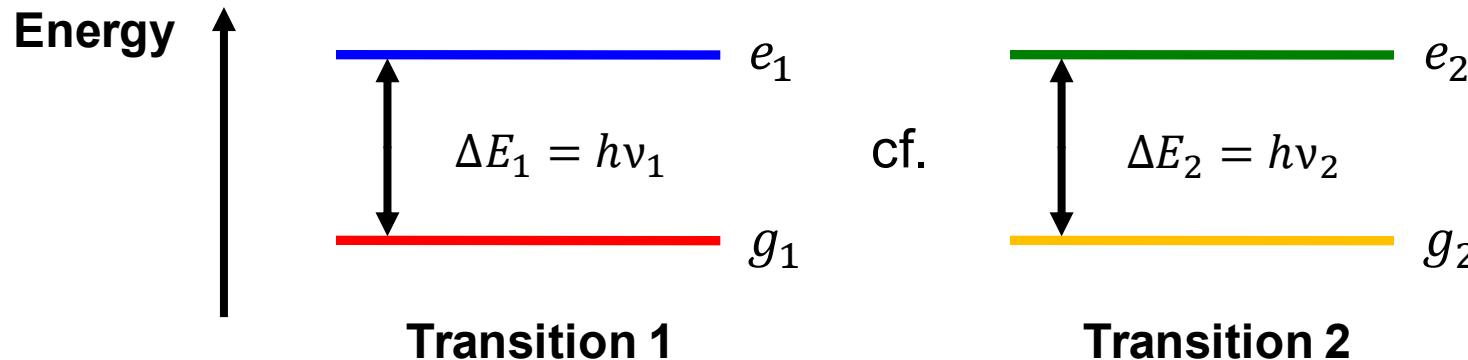
Recent overview: [Ludlow, Boyd, Ye, Peik, Schmidt, *Rev. Mod. Phys.* **87**, 637 (2015)]

“Sensitivity coefficients” K_X required for the interpretation of experimental data have been calculated extensively by Flambaum group

Reviews: [Flambaum, Dzuba, *Can. J. Phys.* **87**, 25 (2009); *Hyperfine Interac.* **236**, 79 (2015)]

Atomic Spectroscopy Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015)], [Arvanitaki, Huang, Van Tilburg, *PRD* **91**, 015015 (2015)]



$$\frac{\delta(\nu_1/\nu_2)}{\nu_1/\nu_2} \propto \sum_{X=\alpha, m_e/m_N, \dots} (K_{X,1} - K_{X,2}) \cos(2\pi f_{\text{DM}} t) ; 2\pi f_{\text{DM}} = m_\varphi \text{ or } 2m_\varphi$$

- **Dy/Cs [Mainz]:** [Van Tilburg *et al.*, *PRL* **115**, 011802 (2015)],
[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]
- **Rb/Cs [SYRTE]:** [Hees *et al.*, *PRL* **117**, 061301 (2016)],
[Stadnik, Flambaum, *PRA* **94**, 022111 (2016)]
- **Yb⁺(E3)/Sr [PTB]:** [Huntemann, Peik *et al.*, Ongoing]
- **Al⁺/Yb, Yb/Sr, Al⁺/Hg⁺ [NIST + JILA]:** [Hume, Leibrandt *et al.*, Ongoing]

Cavity-Based Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]

Solid material



$$L_{\text{solid}} \propto a_B = 1/(m_e \alpha)$$

$$\Rightarrow v_{\text{solid}} \propto 1/L_{\text{solid}} \propto m_e \alpha$$

(adiabatic regime)

Cavity-Based Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

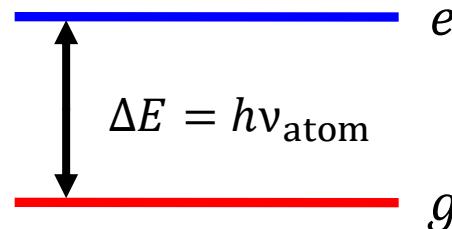
[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]

Solid material



cf.

Electronic transition



$$L_{\text{solid}} \propto a_B = 1/(m_e \alpha)$$

$$\Rightarrow v_{\text{solid}} \propto 1/L_{\text{solid}} \propto m_e \alpha$$

$$v_{\text{atom}} \propto Ry \propto m_e \alpha^2$$

$$\frac{v_{\text{atom}}}{v_{\text{solid}}} \propto \alpha$$

- **Sr vs Glass cavity [Torun]:** [[Wcislo et al., Nature Astronomy 1, 0009 \(2016\)](#)]
- **Various combinations [Worldwide]:** [[Wcislo et al., Science Advances 4, eaau4869 \(2018\)](#)
 - **Cs vs Steel cavity [Mainz]:** [[Antypas et al., PRL 123, 141102 \(2019\)](#)]
 - **Sr⁺ vs Glass cavity [Weizmann]:** [[Aharony et al., arXiv:1902.02788](#)]
 - **Sr/H vs Silicon cavity [JILA + PTB]:** [[Kennedy et al., PRL 125, 201302 \(2020\)](#)]
 - **H vs Sapphire/Quartz cavities [UWA]:** [[Campbell et al., arXiv:2010.08107](#)]

Cavity-Based Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]

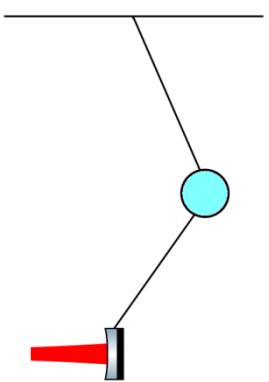
Solid material



Freely-suspended mirrors



Double-pendulum suspensions



$$L_{\text{solid}} \propto a_B = 1/(m_e \alpha)$$

$$\Rightarrow v_{\text{solid}} \propto 1/L_{\text{solid}} \propto m_e \alpha$$

$$L_{\text{free}} \approx \text{const. for } f_{\text{DM}} > f_{\text{natural}}$$

$$\Rightarrow v_{\text{free}} \approx \text{constant}$$

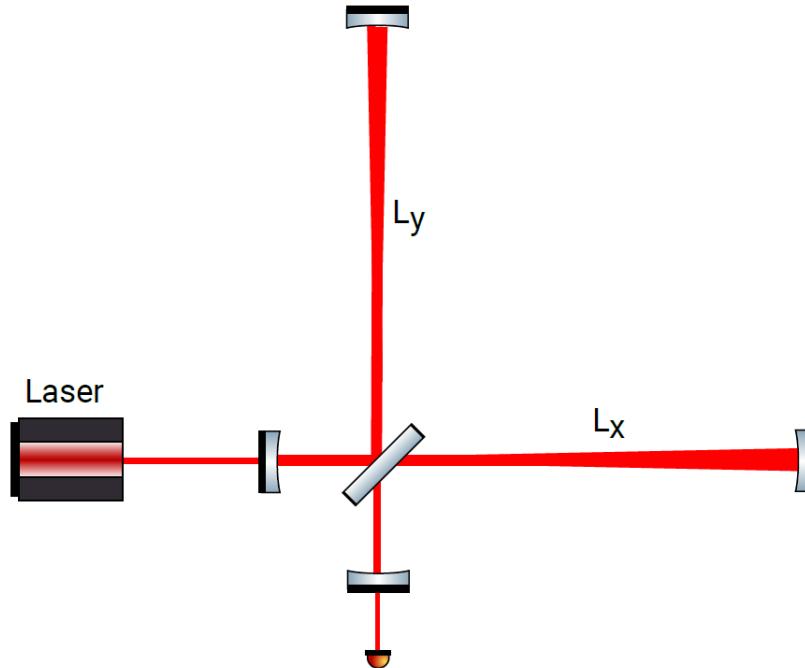
$$\frac{v_{\text{solid}}}{v_{\text{free}}} \propto m_e \alpha$$

$$\text{cf. } \frac{v_{\text{atom}}}{v_{\text{solid}}} \propto \alpha$$

Small-scale experiment currently under development at Northwestern University

Laser Interferometry Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

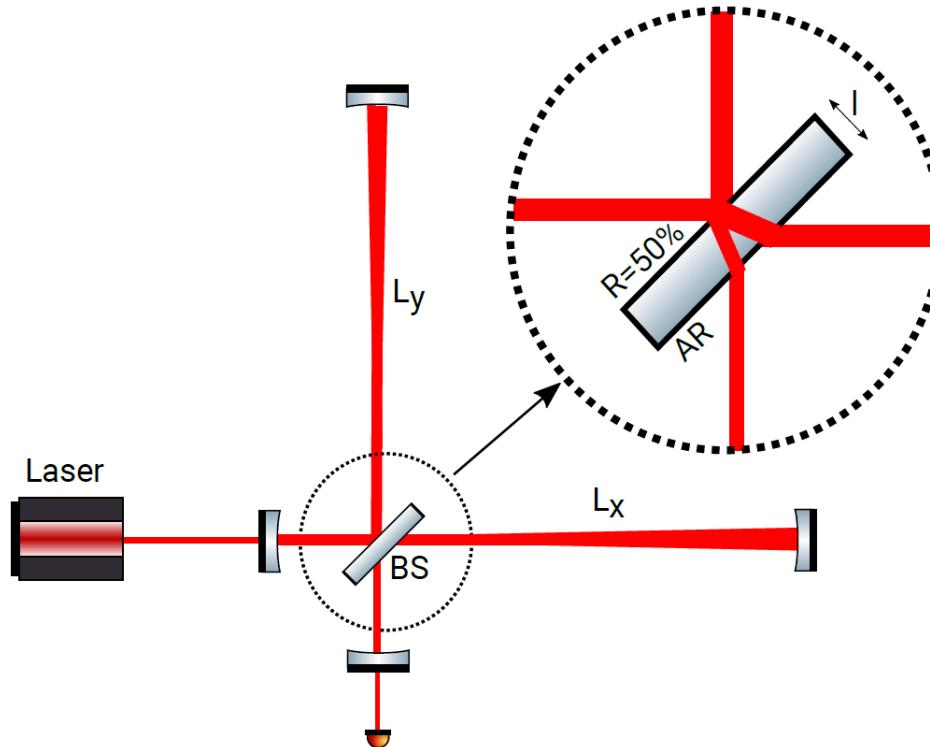
[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]



Michelson interferometer (GEO600)

Laser Interferometry Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

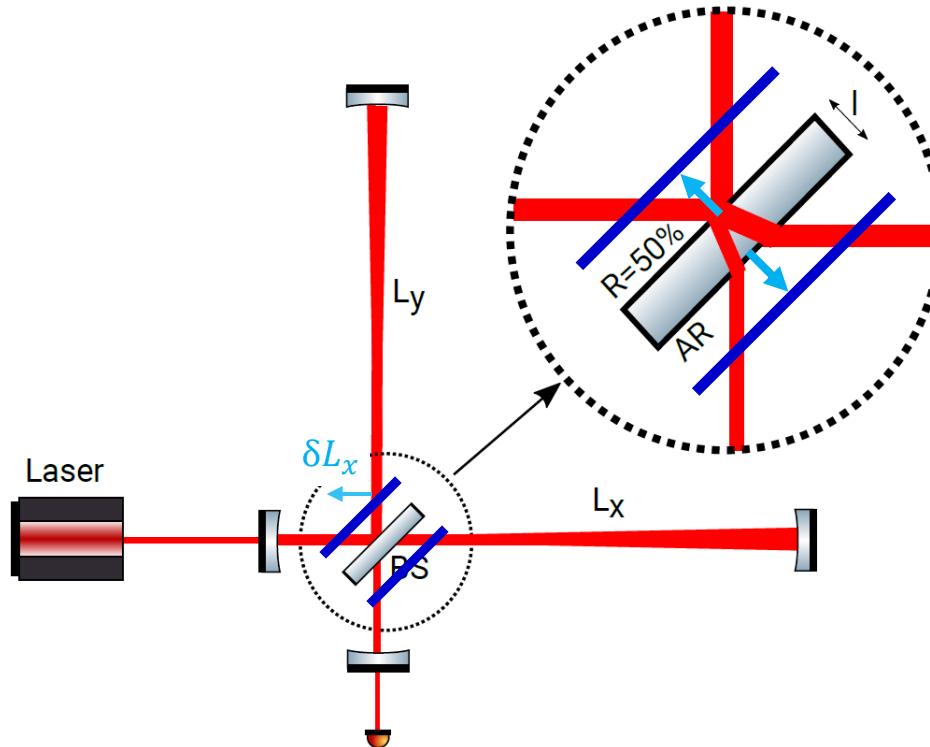
[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]



- Geometric asymmetry from beam-splitter

Laser Interferometry Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

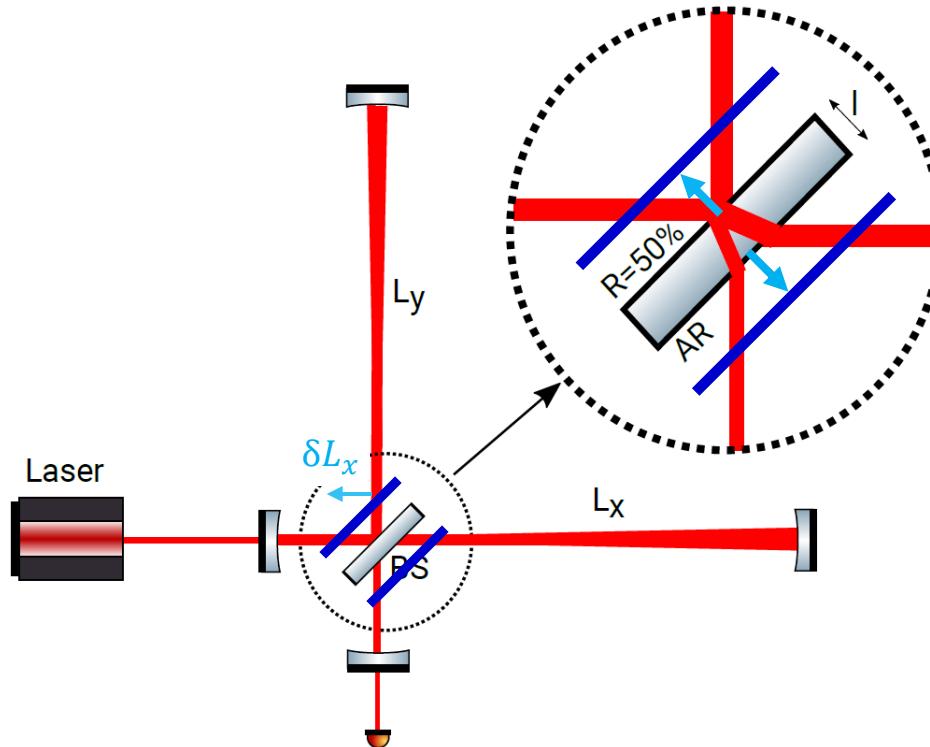
[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]



- Geometric asymmetry from beam-splitter: $\delta(L_x - L_y) \sim \delta(nl)$

Laser Interferometry Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]



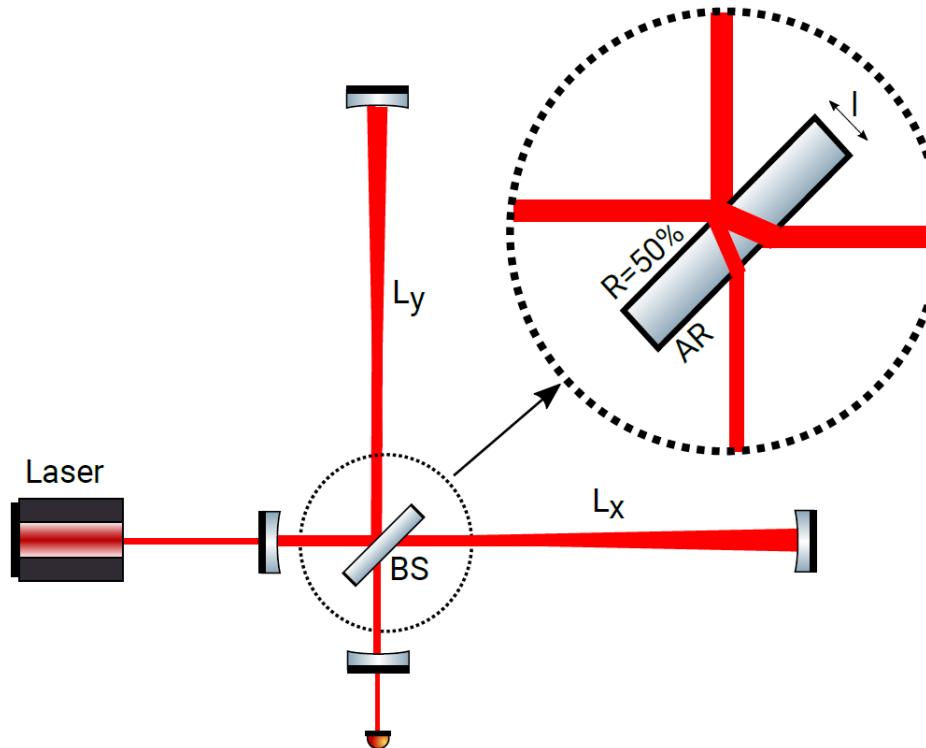
- Geometric asymmetry from beam-splitter: $\delta(L_x - L_y) \sim \delta(nl)$

First results recently reported using GEO600 data:

[Vermeulen *et al.*, arXiv:2103.03783]

Laser Interferometry Searches for Oscillating Variations of Fundamental Constants induced by Dark Matter

[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]

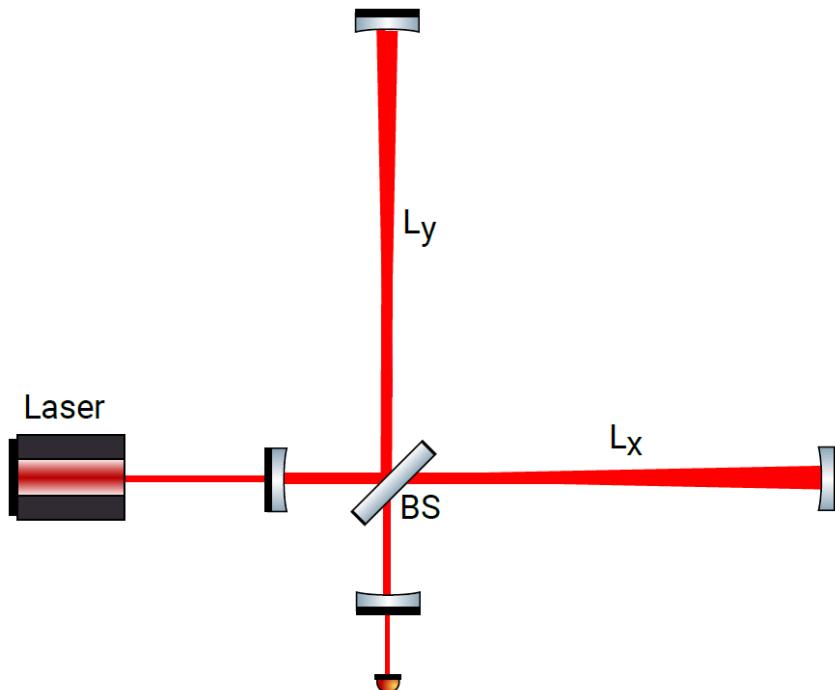


- Geometric asymmetry from beam-splitter: $\delta(L_x - L_y) \sim \delta(nl)$
- Both broadband and resonant narrowband searches possible:
$$f_{\text{DM}} \approx f_{\text{vibr,BS}}(T) \sim v_{\text{sound}}/l \Rightarrow Q \sim 10^6 \text{ enhancement}$$

Michelson vs Fabry-Perot-Michelson Interferometers

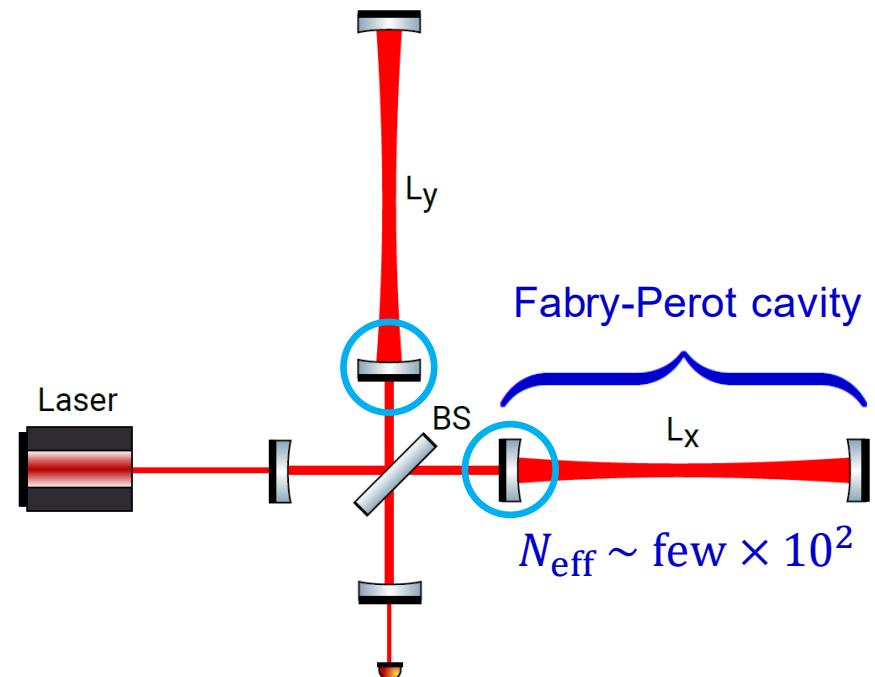
[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]

**Michelson interferometer
(GEO 600)**



$$\delta(L_x - L_y)_{\text{BS}} \sim \delta(nl)$$

**Fabry-Perot-Michelson
interferometer (LIGO)**

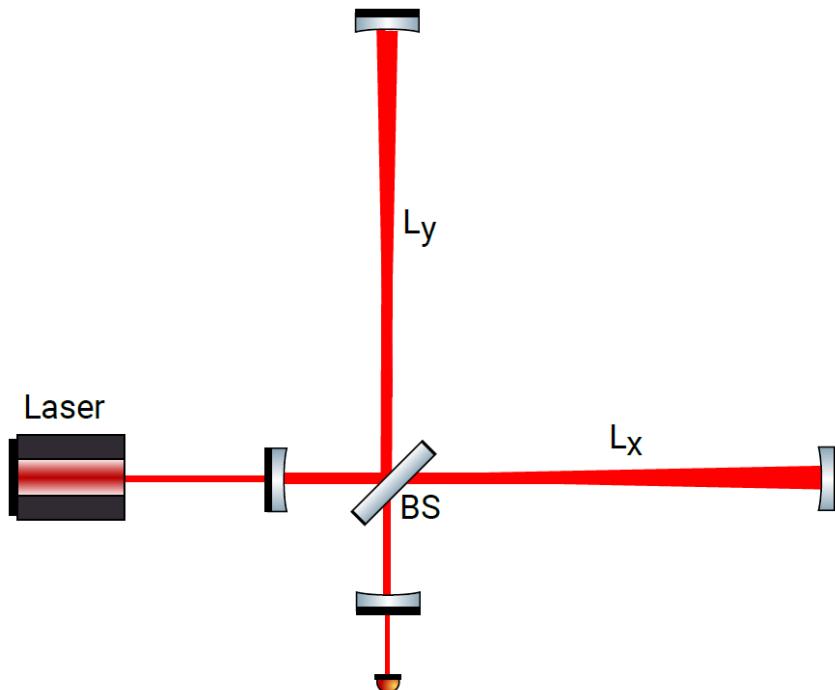


$$\delta(L_x - L_y)_{\text{BS}} \sim \delta(nl)/N_{\text{eff}}$$

Michelson vs Fabry-Perot-Michelson Interferometers

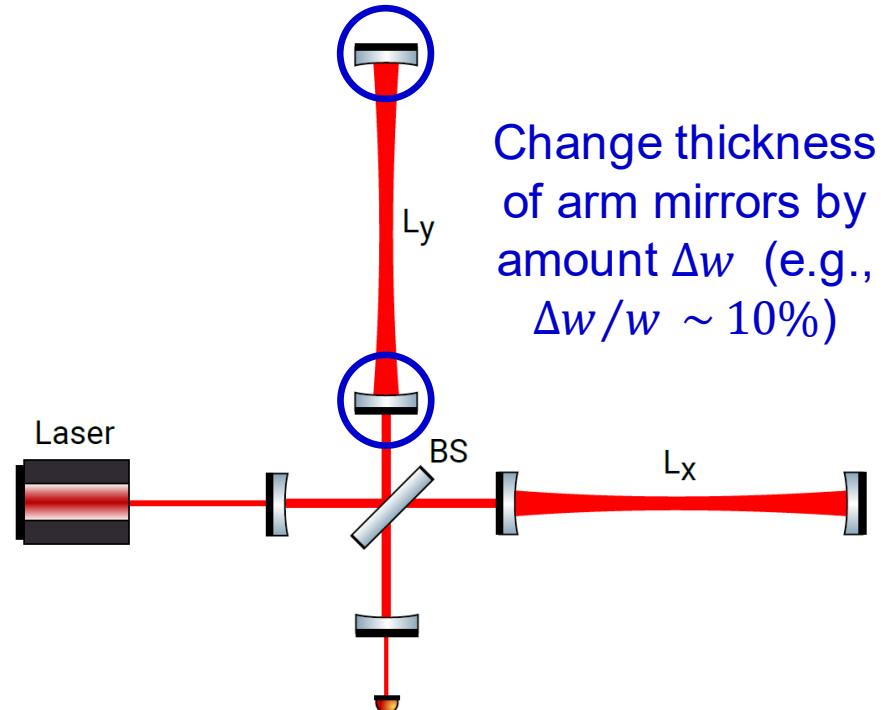
[Grote, Stadnik, *Phys. Rev. Research* 1, 033187 (2019)]

**Michelson interferometer
(GEO 600)**



$$\delta(L_x - L_y)_{\text{BS}} \sim \delta(nl)$$

**Fabry-Perot-Michelson
interferometer (LIGO)**



Change thickness
of arm mirrors by
amount Δw (e.g.,
 $\Delta w/w \sim 10\%$)

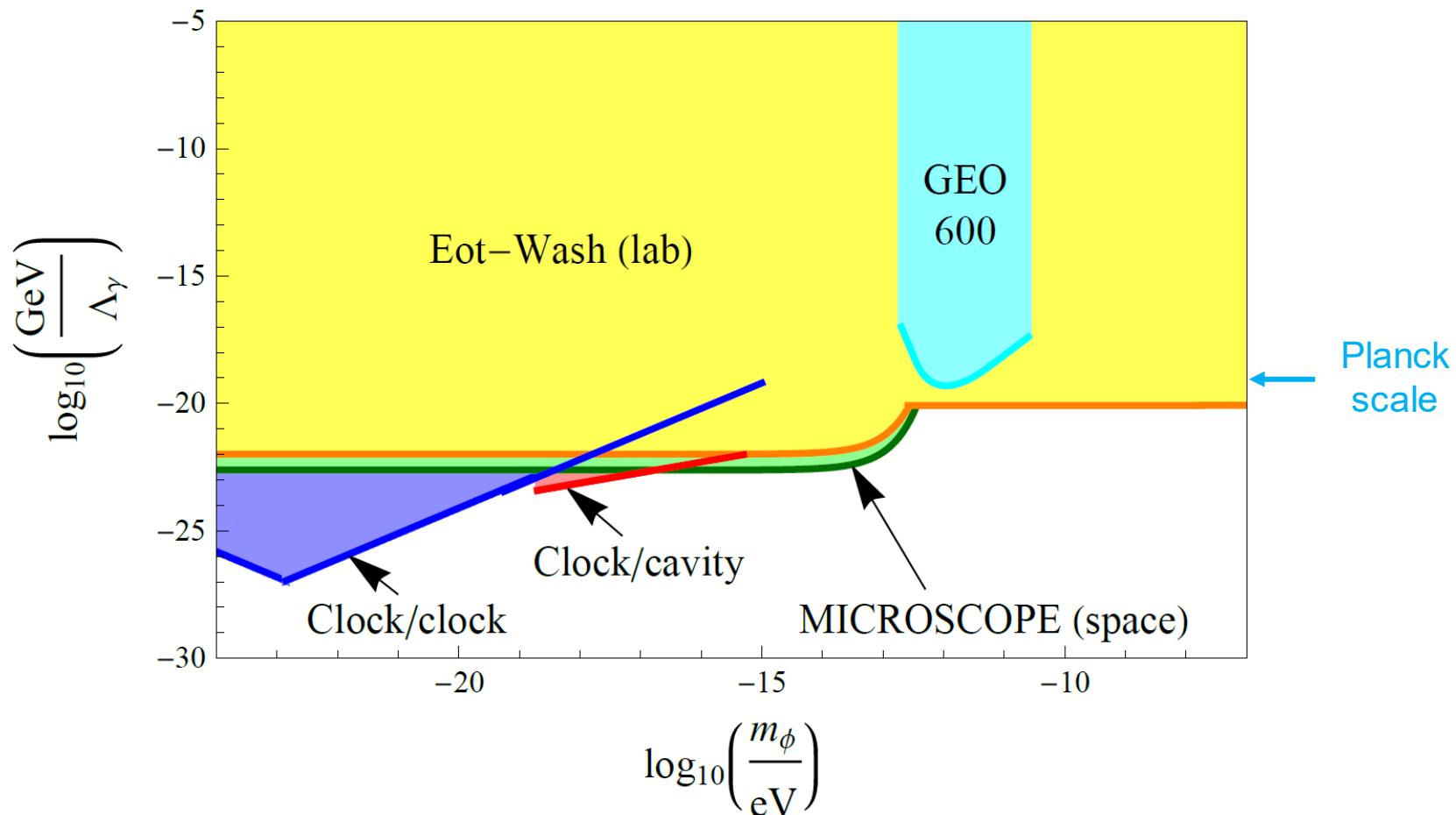
$$\delta(L_x - L_y) \approx \delta(\Delta w)$$

Constraints on Linear Interaction of Scalar Dark Matter with the Photon

Clock/clock: [Van Tilburg et al., *PRL* **115**, 011802 (2015)], [Hees et al., *PRL* **117**, 061301 (2016)];

Clock/cavity: [Kennedy et al., *PRL* **125**, 201302 (2020)]; **GEO600**: [Vermeulen et al., arXiv:2103.03783]

4 orders of magnitude improvement!

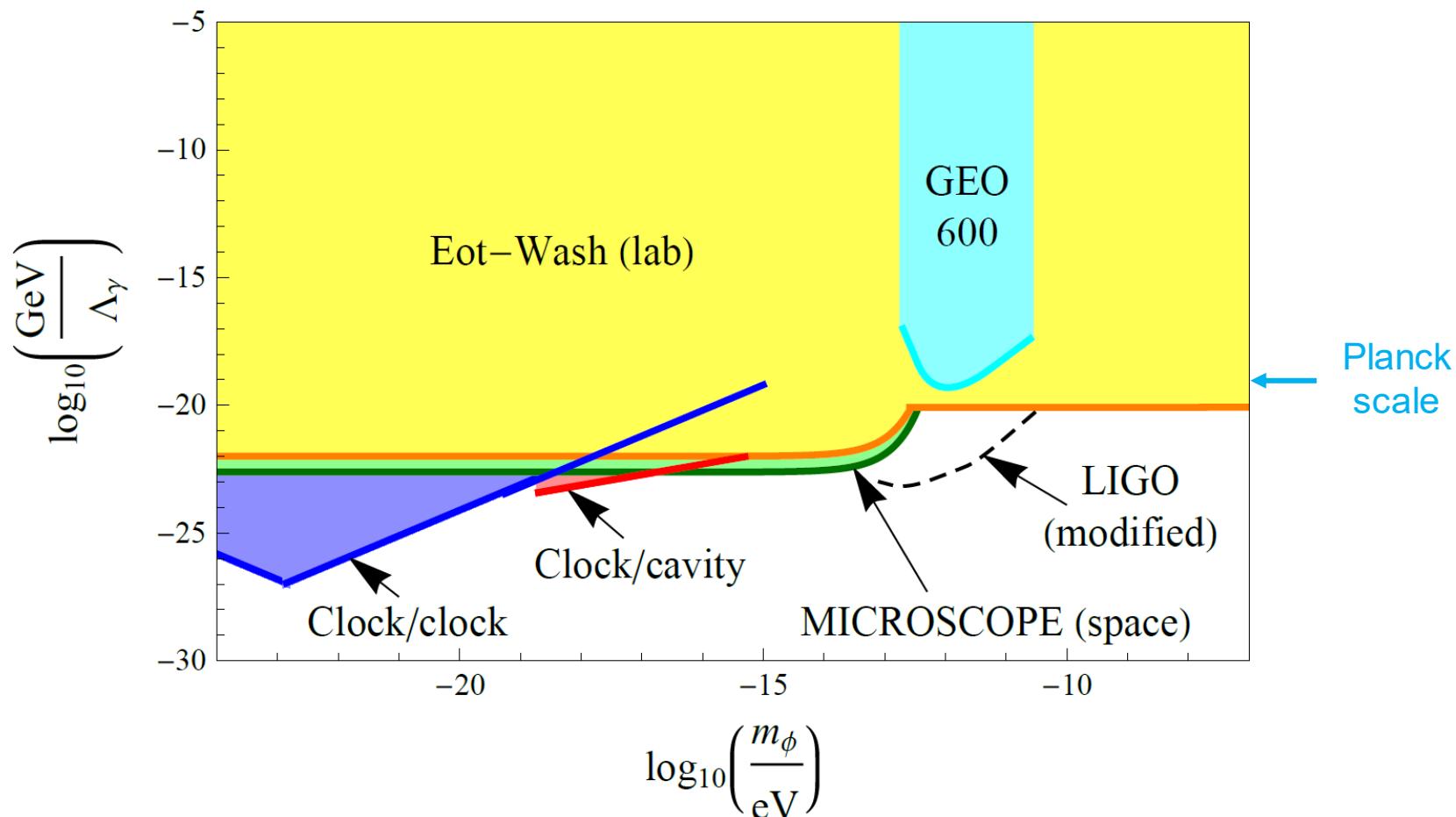


Constraints on Linear Interaction of Scalar Dark Matter with the Photon

Clock/clock: [Van Tilburg et al., *PRL* **115**, 011802 (2015)], [Hees et al., *PRL* **117**, 061301 (2016)];

Clock/cavity: [Kennedy et al., *PRL* **125**, 201302 (2020)]; **GEO600**: [Vermeulen et al., arXiv:2103.03783]

4 orders of magnitude improvement!

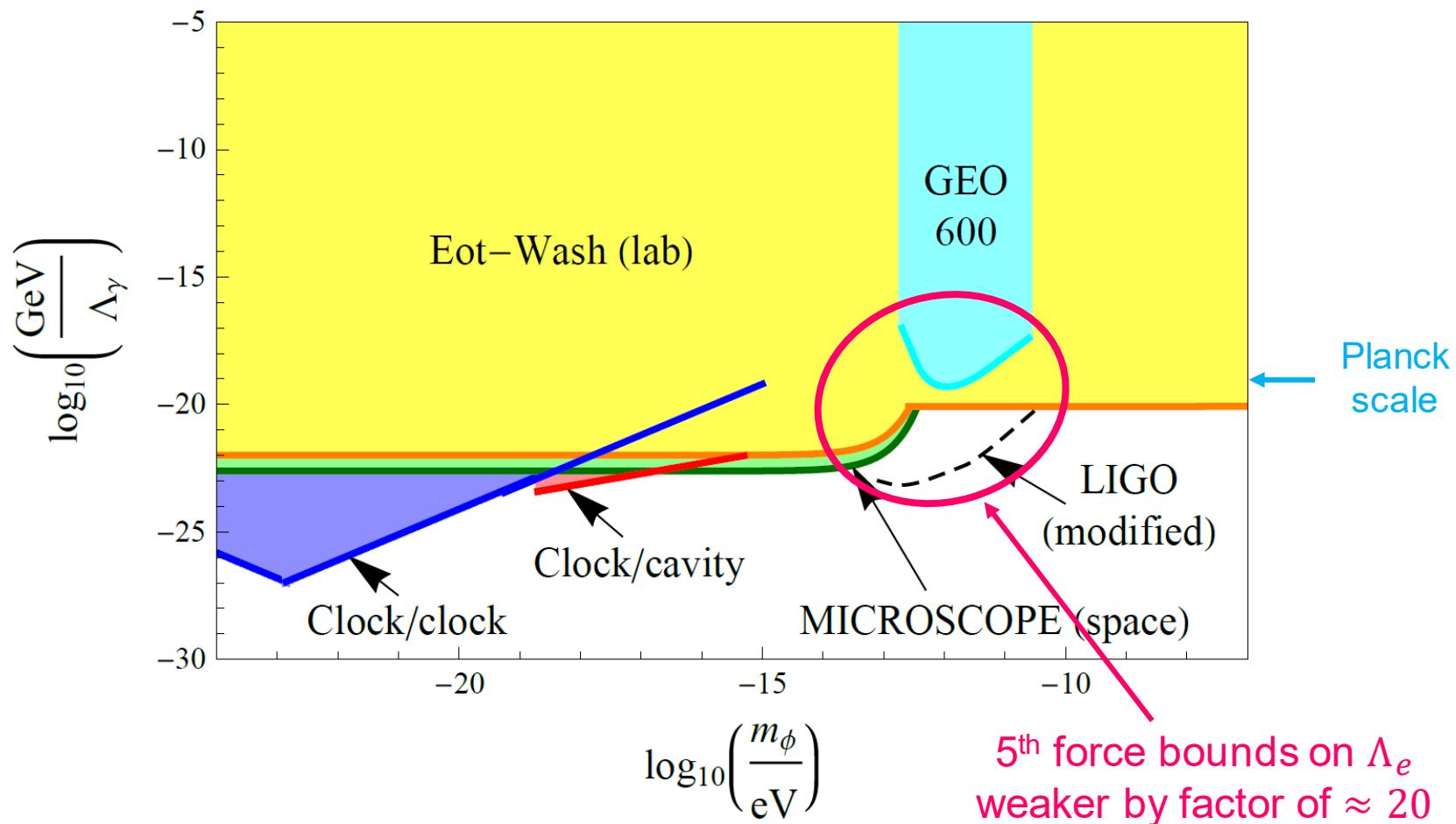


Constraints on Linear Interaction of Scalar Dark Matter with the Photon

Clock/clock: [Van Tilburg et al., *PRL* **115**, 011802 (2015)], [Hees et al., *PRL* **117**, 061301 (2016)];

Clock/cavity: [Kennedy et al., *PRL* **125**, 201302 (2020)]; **GEO600**: [Vermeulen et al., arXiv:2103.03783]

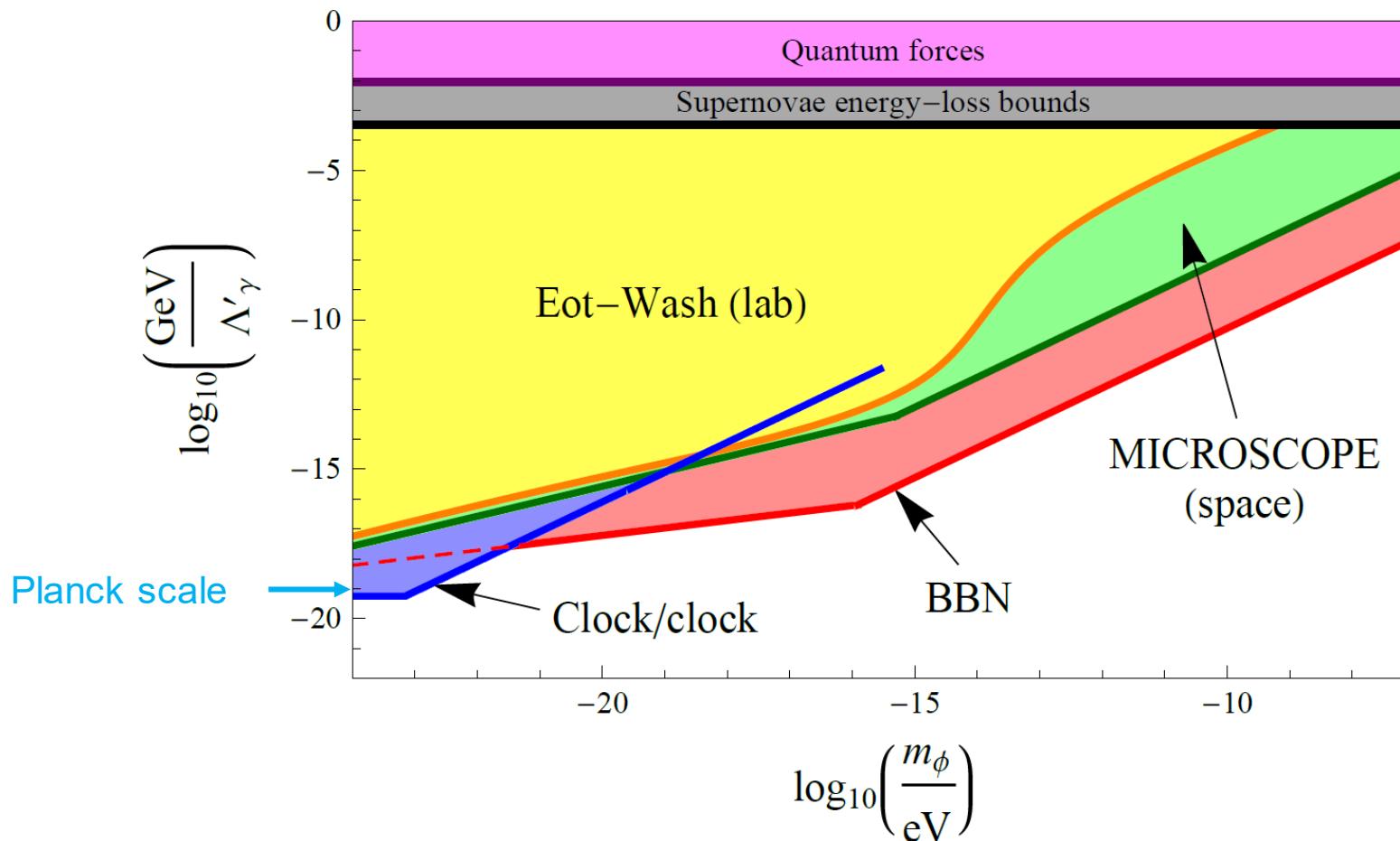
4 orders of magnitude improvement!



Constraints on Quadratic Interaction of Scalar Dark Matter with the Photon

Clock/clock + BBN constraints: [Stadnik, Flambaum, *PRL* **115**, 201301 (2015); *PRA* **94**, 022111 (2016)]; **MICROSCOPE + Eöt-Wash constraints:** [Hees et al., *PRD* **98**, 064051 (2018)]

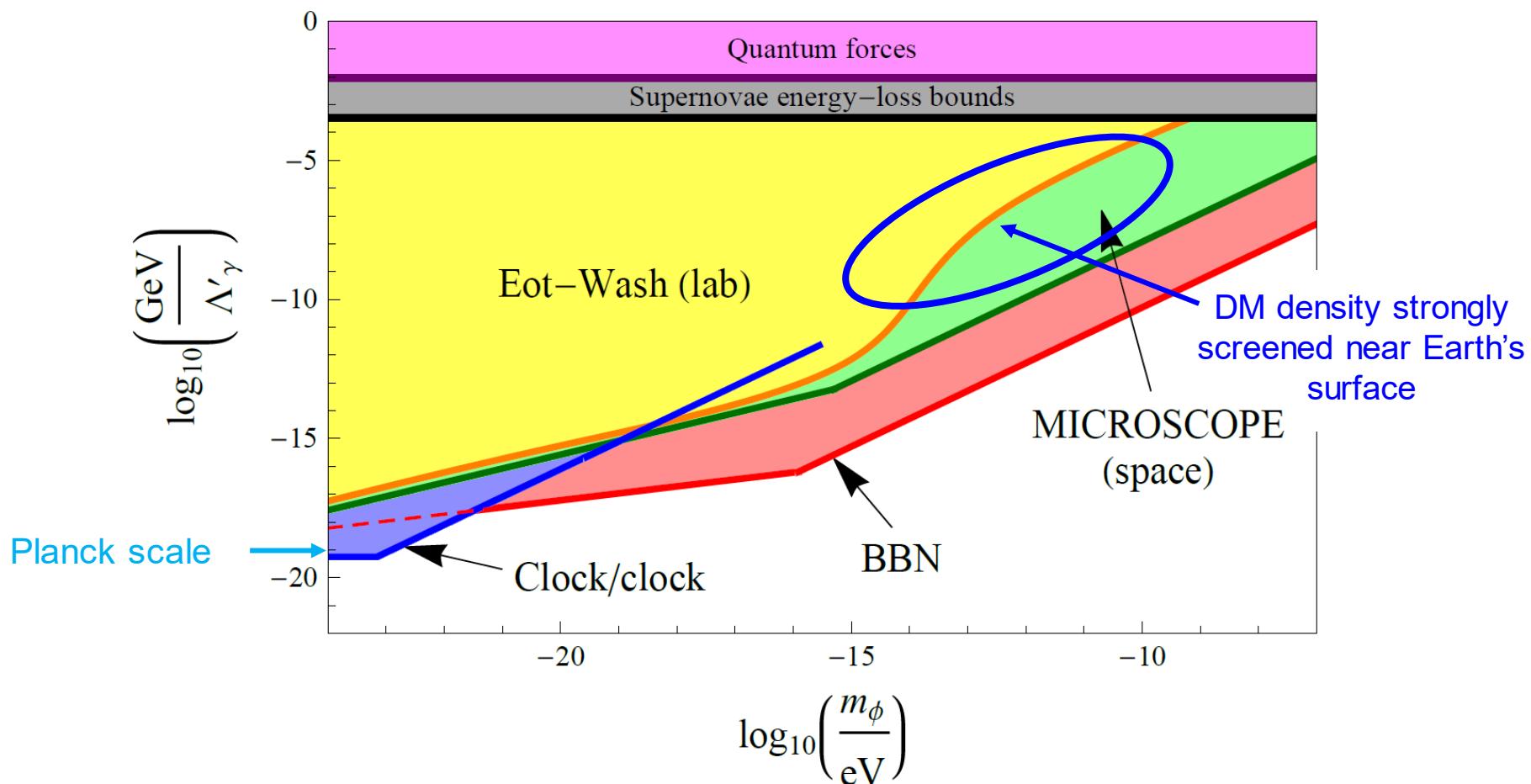
15 orders of magnitude improvement!



Constraints on Quadratic Interaction of Scalar Dark Matter with the Photon

Clock/clock + BBN constraints: [Stadnik, Flambaum, *PRL* **115**, 201301 (2015); *PRA* **94**, 022111 (2016)]; MICROSCOPE + Eöt-Wash constraints: [Hees et al., *PRD* **98**, 064051 (2018)]

15 orders of magnitude improvement!



Low-mass Spin-0 Dark Matter

Dark Matter



*More traditional axion detection methods
tend to focus on the **electromagnetic**
coupling*

*Here I focus on relatively new
detection methods based on
non-electromagnetic couplings*



Pseudoscalars
(Axions):

$$\varphi \xrightarrow{P} -\varphi$$

**Time-varying spin-
dependent effects**

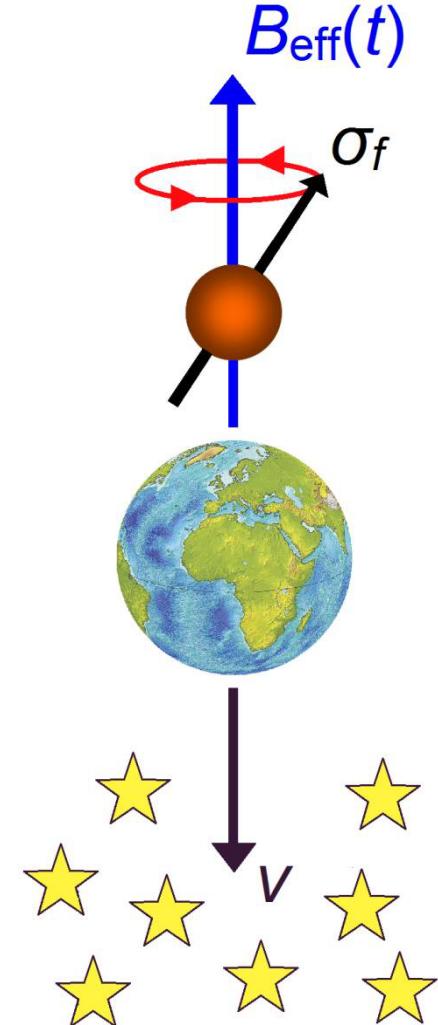
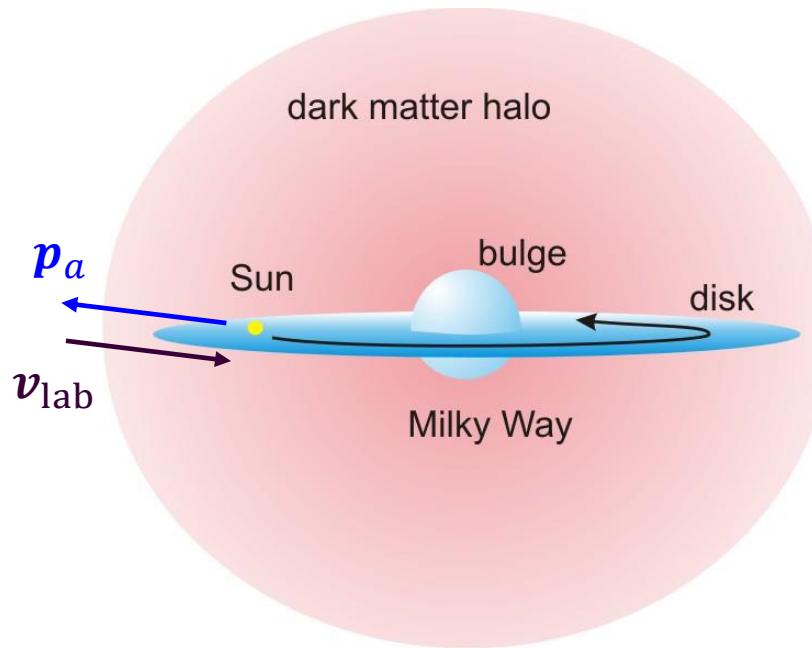
- Co-magnetometers
- Particle g -factors
- Spin-polarised torsion pendula
- Spin resonance (NMR, ESR)

“Axion Wind” Spin-Precession Effect

[Flambaum, talk at Patras Workshop, 2013], [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_f = -\frac{C_f}{2f_a} \partial_i [a_0 \cos(m_a t - \mathbf{p}_a \cdot \mathbf{x})] \bar{f} \gamma^i \gamma^5 f$$

$$\Rightarrow H_{\text{wind}}(t) = \boldsymbol{\sigma}_f \cdot \mathbf{B}_{\text{eff}}(t) \propto \boldsymbol{\sigma}_f \cdot \mathbf{p}_a \sin(m_a t)$$

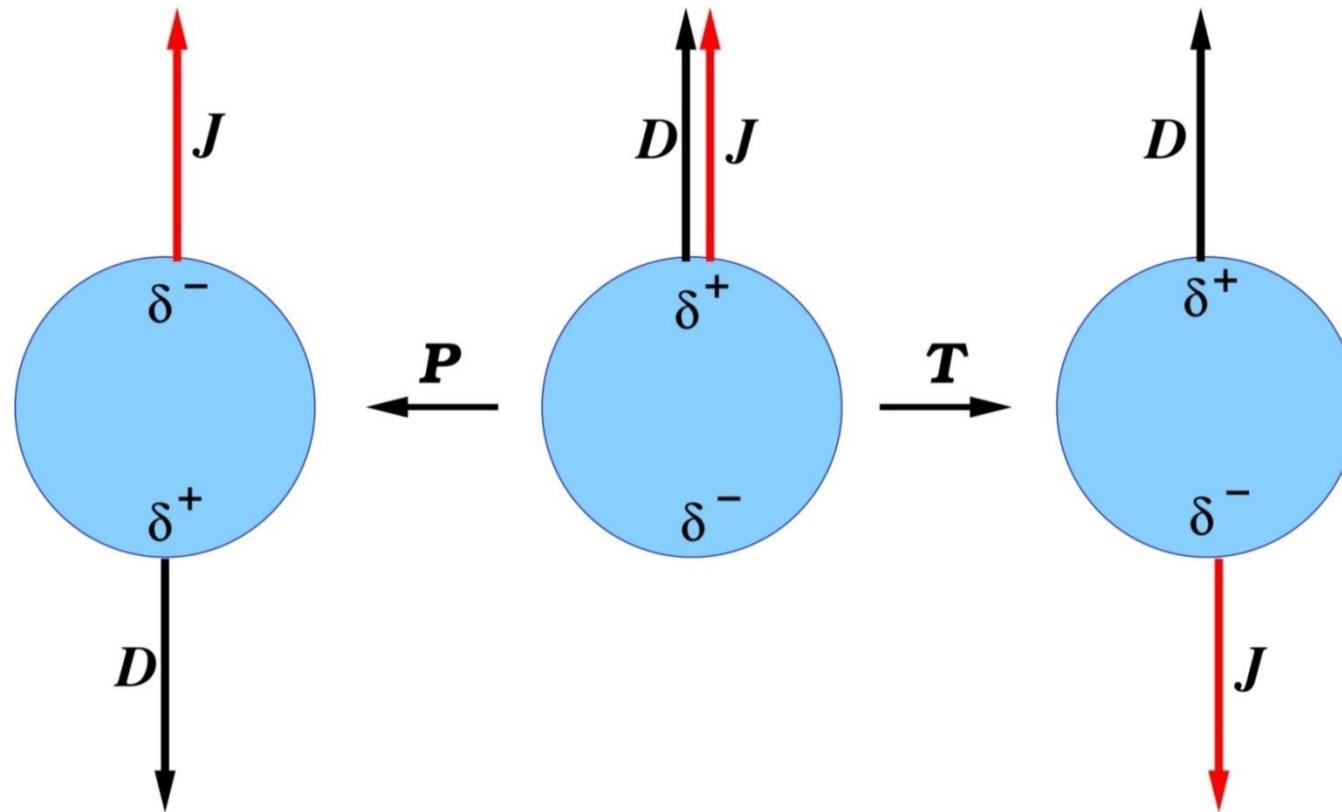


Oscillating Electric Dipole Moments

Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)]

Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

Electric Dipole Moment (EDM) = parity (P) and time-reversal-invariance (T) violating electric moment



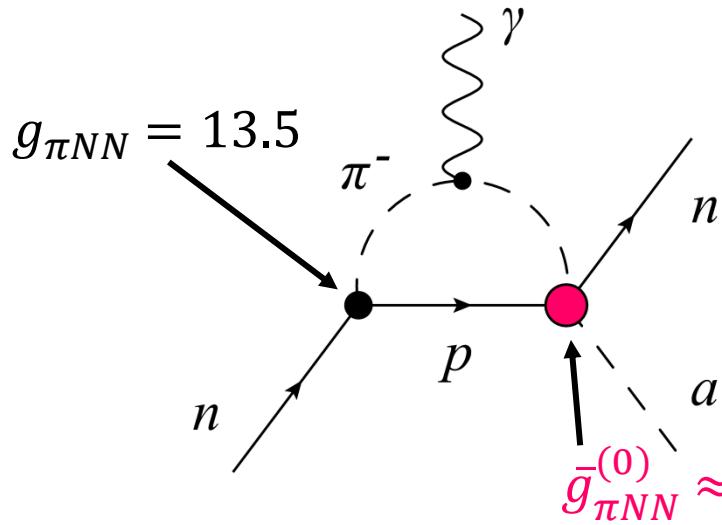
Oscillating Electric Dipole Moments

Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)]

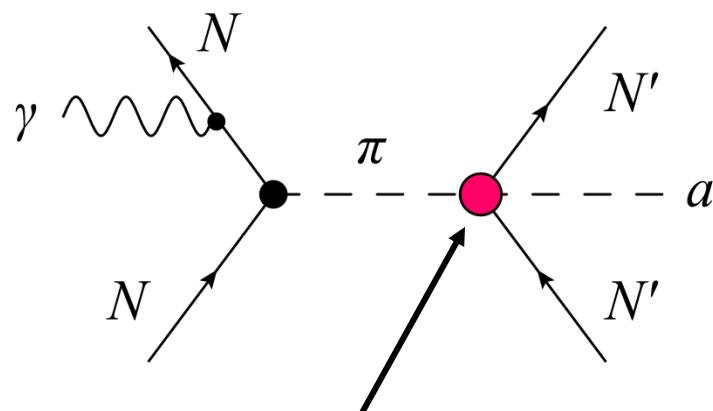
Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_G = \frac{C_G g^2}{32\pi^2 f_a} a_0 \cos(m_a t) G \tilde{G} \Rightarrow H_{\text{EDM}}(t) = \mathbf{d}(t) \cdot \mathbf{E},$$
$$\mathbf{d}(t) \propto J \cos(m_a t)$$

Nucleon EDMs



CP-violating intranuclear forces



In nuclei, **tree-level** *CP*-violating intranuclear forces dominate over **loop-induced** nucleon EDMs [loop factor = $1/(8\pi^2)$].

Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use spin-polarised sources: Atomic magnetometers,
cold/ultracold particles, torsion pendula

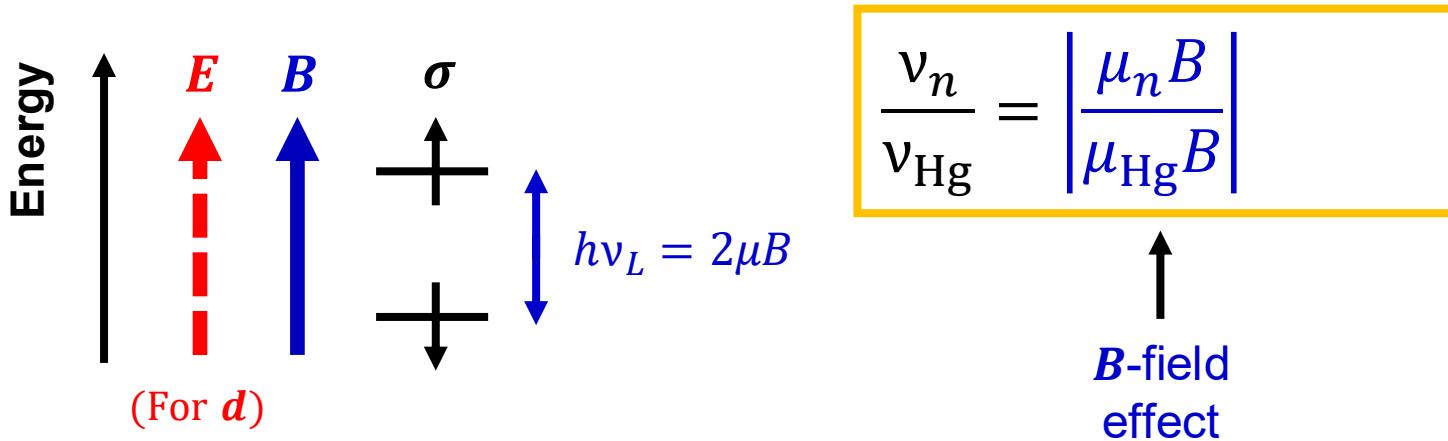
***Similar to previous searches for
Lorentz-invariance violation***

Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use spin-polarised sources: Atomic magnetometers,
cold/ultracold particles, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]

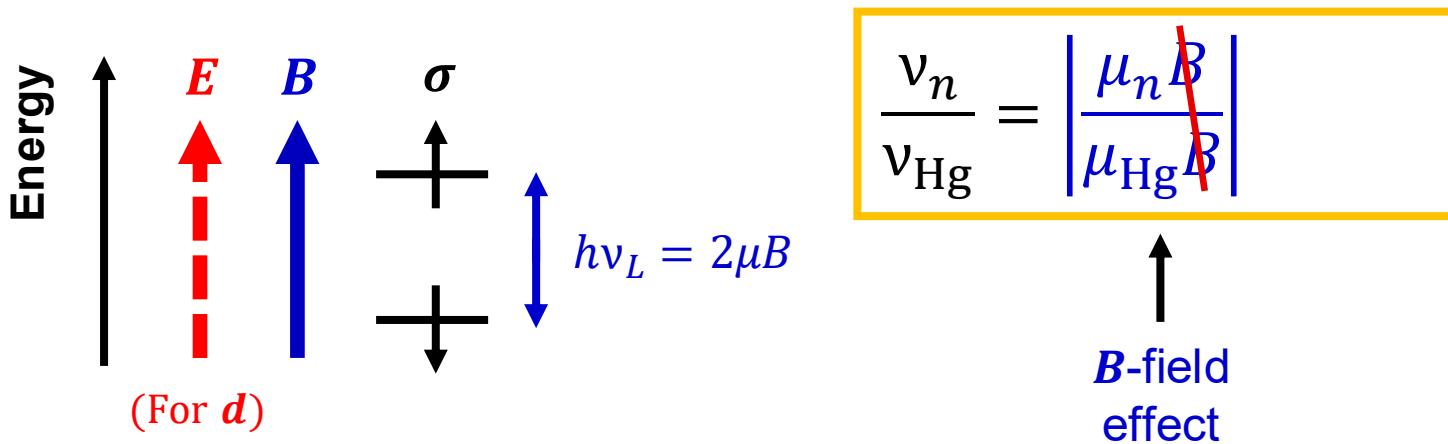


Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use spin-polarised sources: Atomic magnetometers,
cold/ultracold particles, torsion pendula

Experiment (*n*/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]

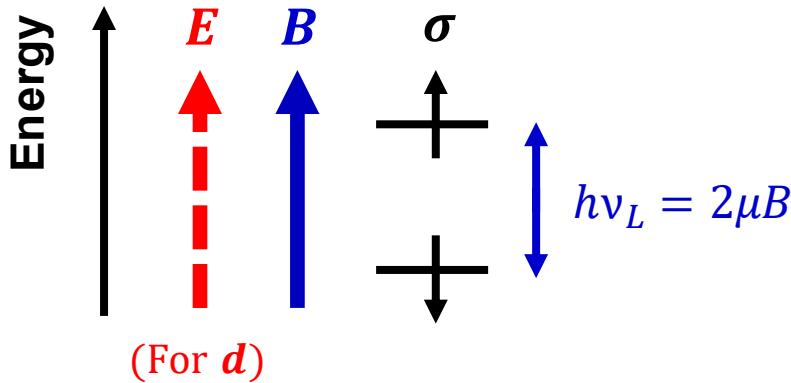


Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use spin-polarised sources: Atomic magnetometers,
cold/ultracold particles, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]



$$\frac{v_n}{v_{\text{Hg}}} = \left| \frac{\mu_n B}{\mu_{\text{Hg}} B} \right| + R(t)$$

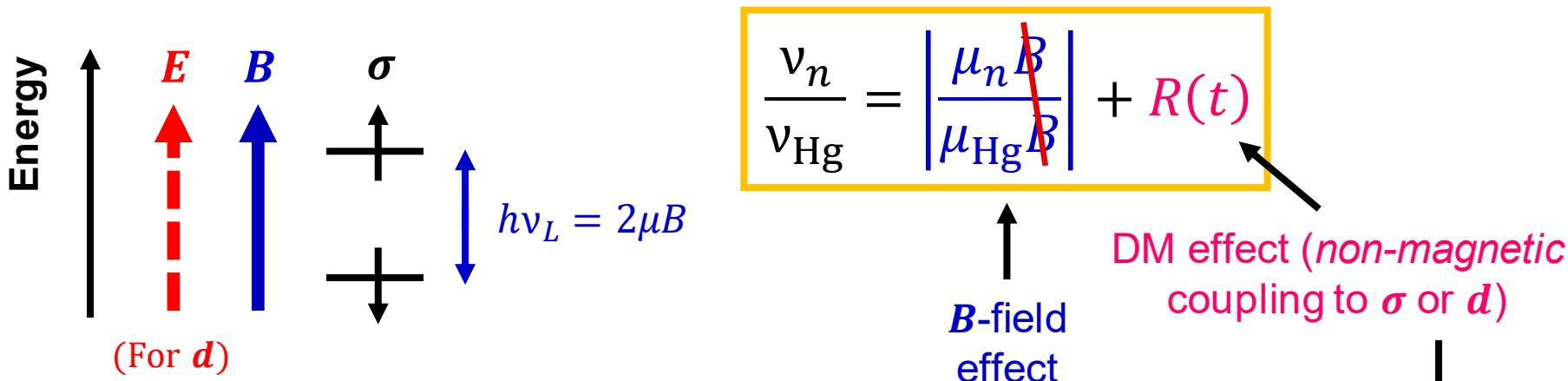
\uparrow
 B -field effect
 \uparrow
DM effect (*non-magnetic*
coupling to σ or d)

Searching for Spin-Dependent Effects

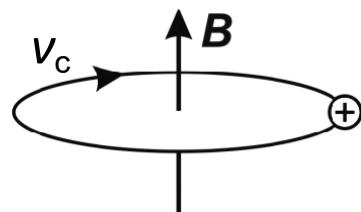
Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use spin-polarised sources: Atomic magnetometers,
cold/ultracold particles, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]



Proposal + Experiment (\bar{p}): [BASE collaboration, *Nature* **575**, 310 (2019)]



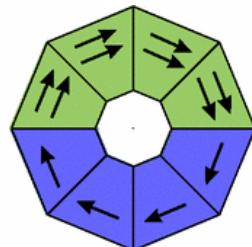
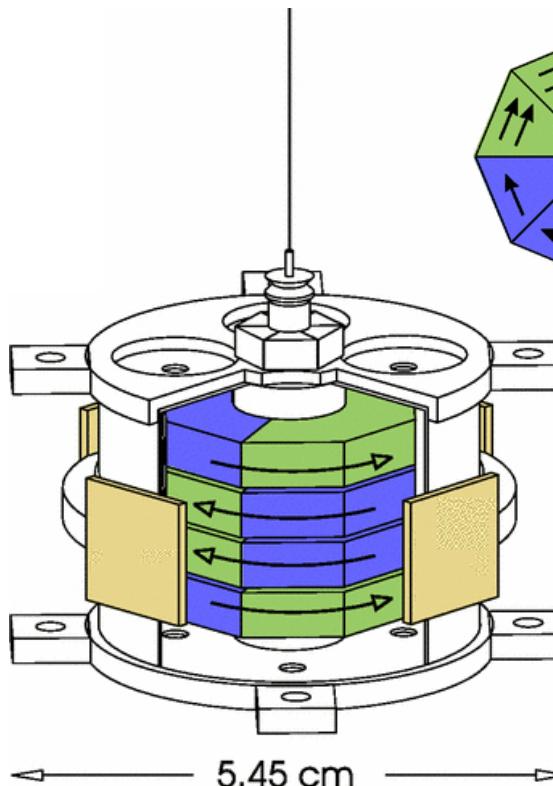
$$\left(\frac{v_L}{v_c} \right)_{\bar{p}} = \frac{|g_{\bar{p}}|}{2} + R(t)$$

Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use spin-polarised sources: *Atomic magnetometers*,
cold/ultracold particles, *torsion pendula*

Experiment (**Alnico/SmCo₅**): [Terrano *et al.*, *PRL* **122**, 231301 (2019)]



$$\begin{array}{c} \mu_{\text{Alnico}} \\ \longrightarrow \\ \longleftarrow \\ \mu_{\text{SmCo}_5} \end{array}$$

$$\mu_{\text{pendulum}} \approx 0$$

$$\begin{array}{c} (\sigma_e)_{\text{Alnico}} \\ \longrightarrow \\ \longleftarrow \\ (\sigma_e)_{\text{SmCo}_5} \end{array}$$

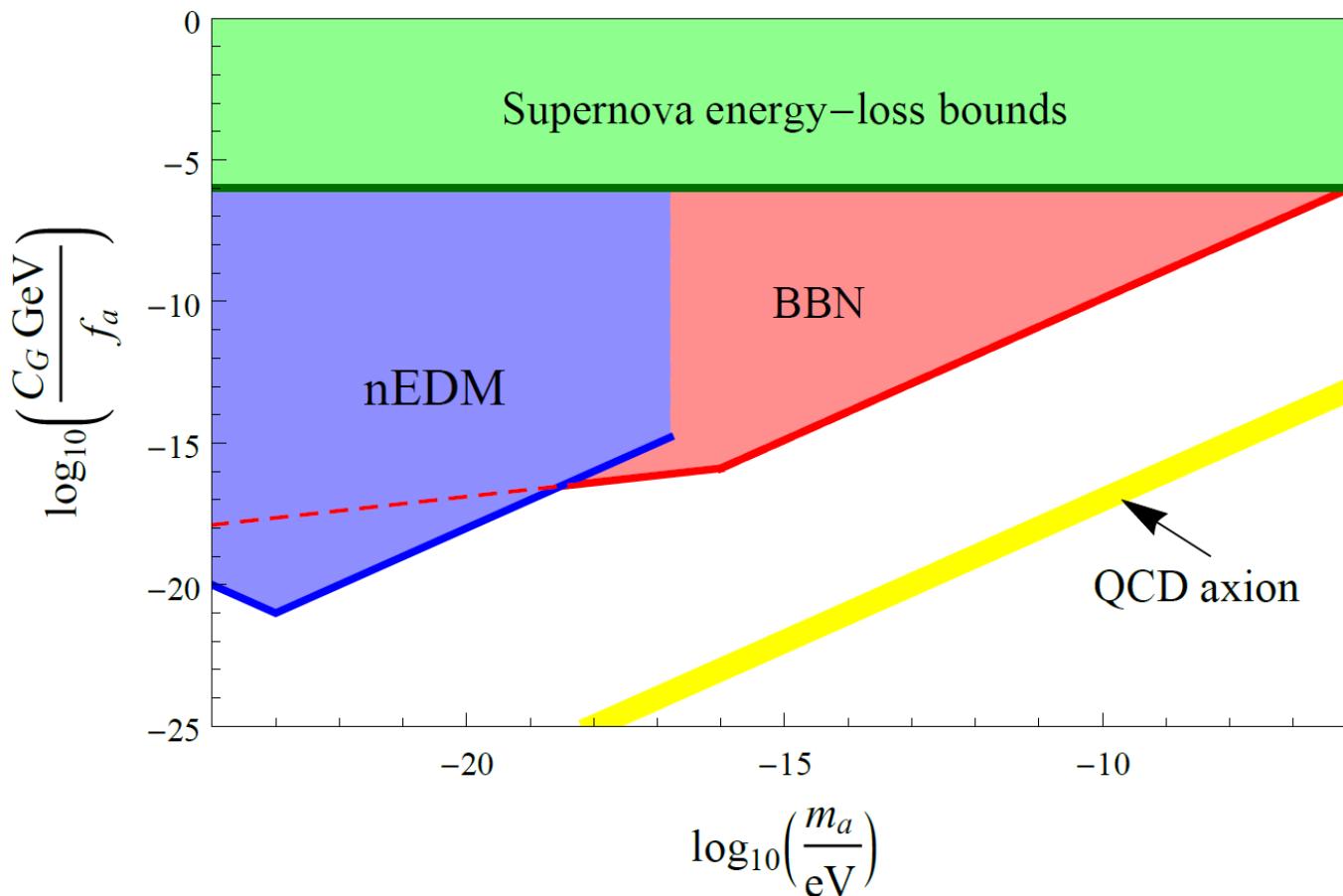
$$(\sigma_e)_{\text{pendulum}} \neq 0$$

$$\tau(t) \propto (\sigma_e)_{\text{pendulum}} \times B_{\text{eff}}(t)$$

Constraints on Interaction of Axion Dark Matter with Gluons

nEDM constraints: [nEDM collaboration, *PRX* 7, 041034 (2017)]

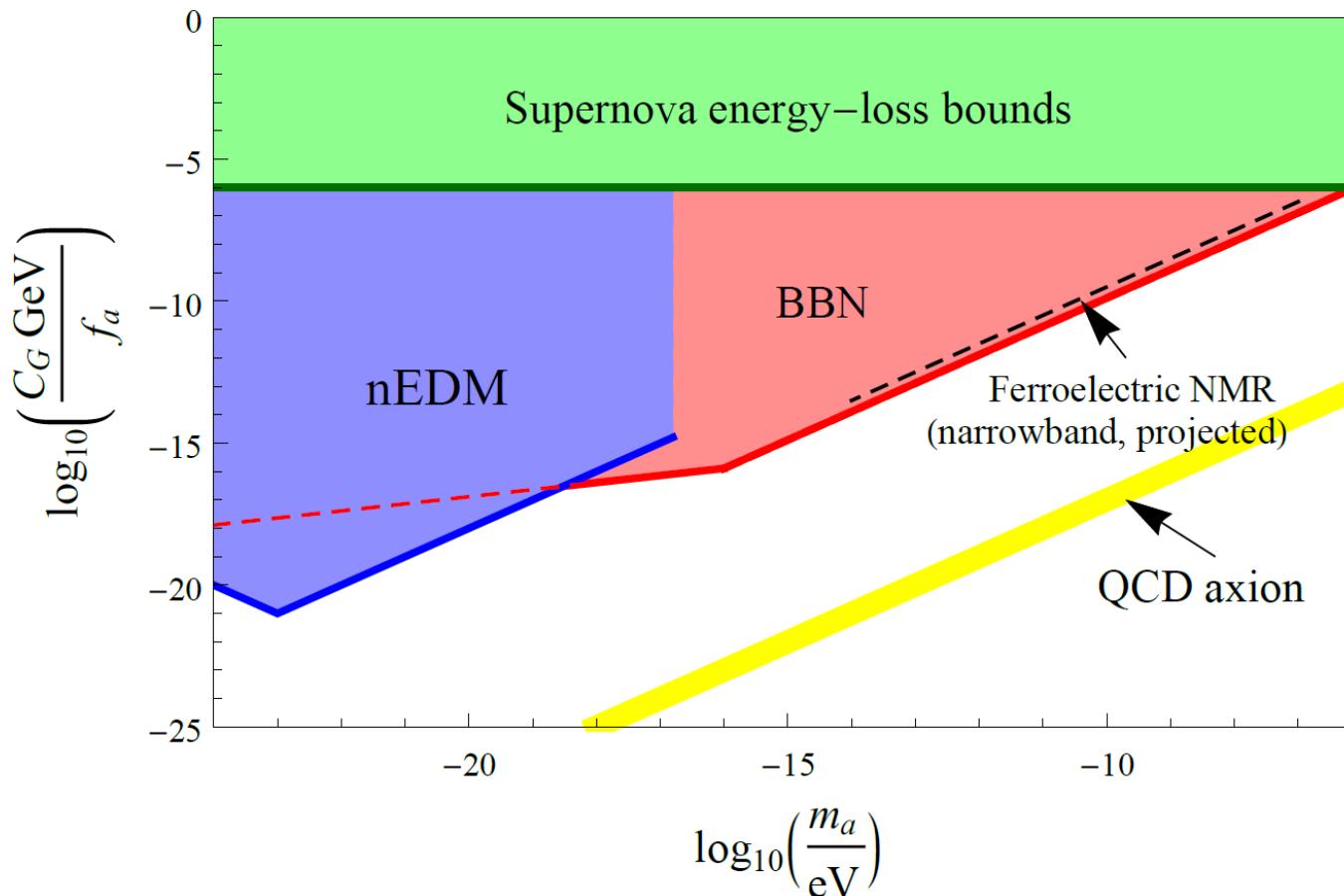
3 orders of magnitude improvement!



Constraints on Interaction of Axion Dark Matter with Gluons

nEDM constraints: [nEDM collaboration, *PRX* 7, 041034 (2017)]

3 orders of magnitude improvement!

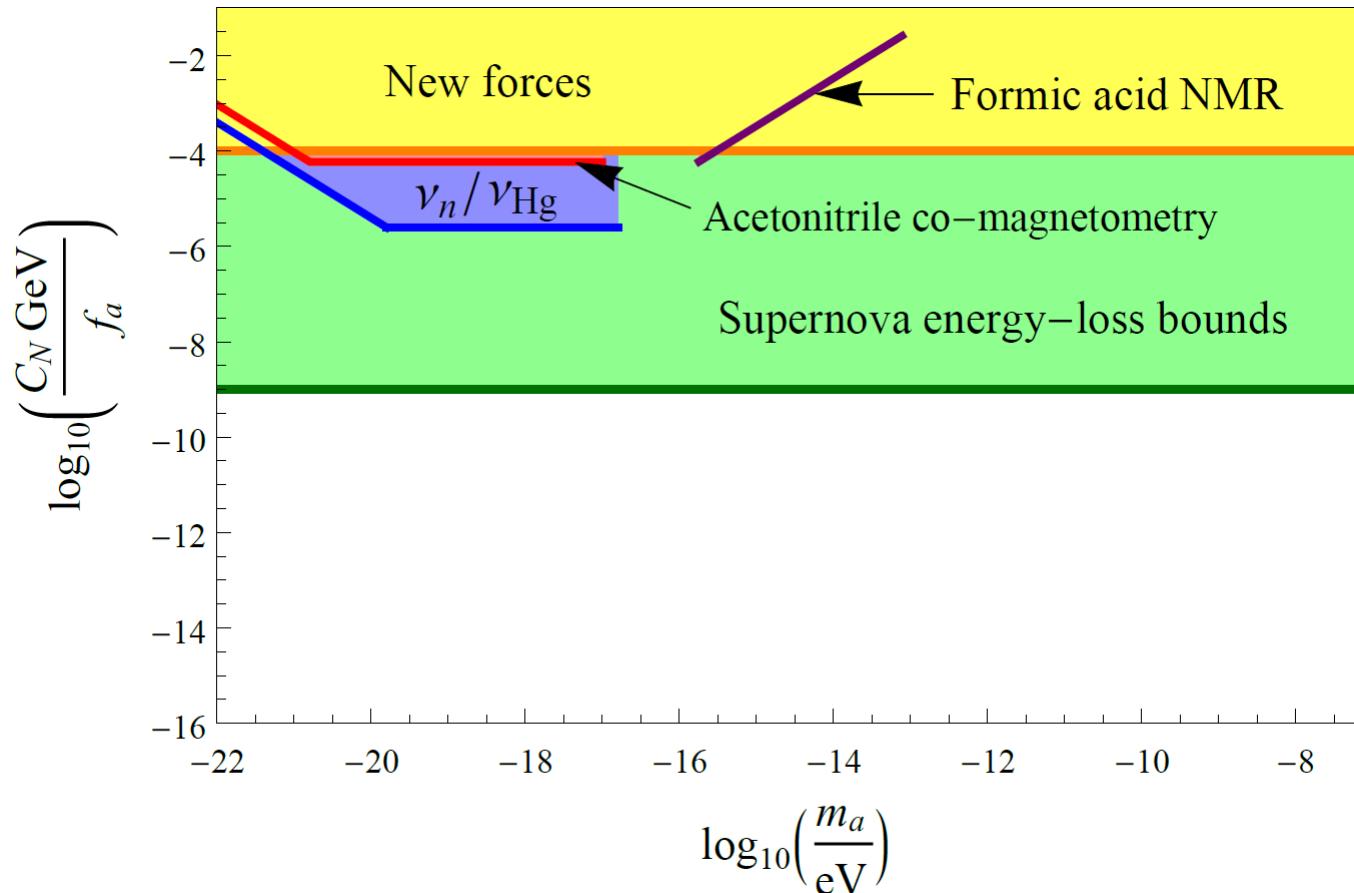


Constraints on Interaction of Axion Dark Matter with Nucleons

v_n/v_{Hg} constraints: [nEDM collaboration, *PRX* **7**, 041034 (2017)]

Acetonitrile constraints: [Wu *et al.*, *PRL* **122**, 191302 (2019)]

Formic acid NMR constraints: [Garcon *et al.*, *Sci. Adv.* **5**, eaax4539 (2019)]

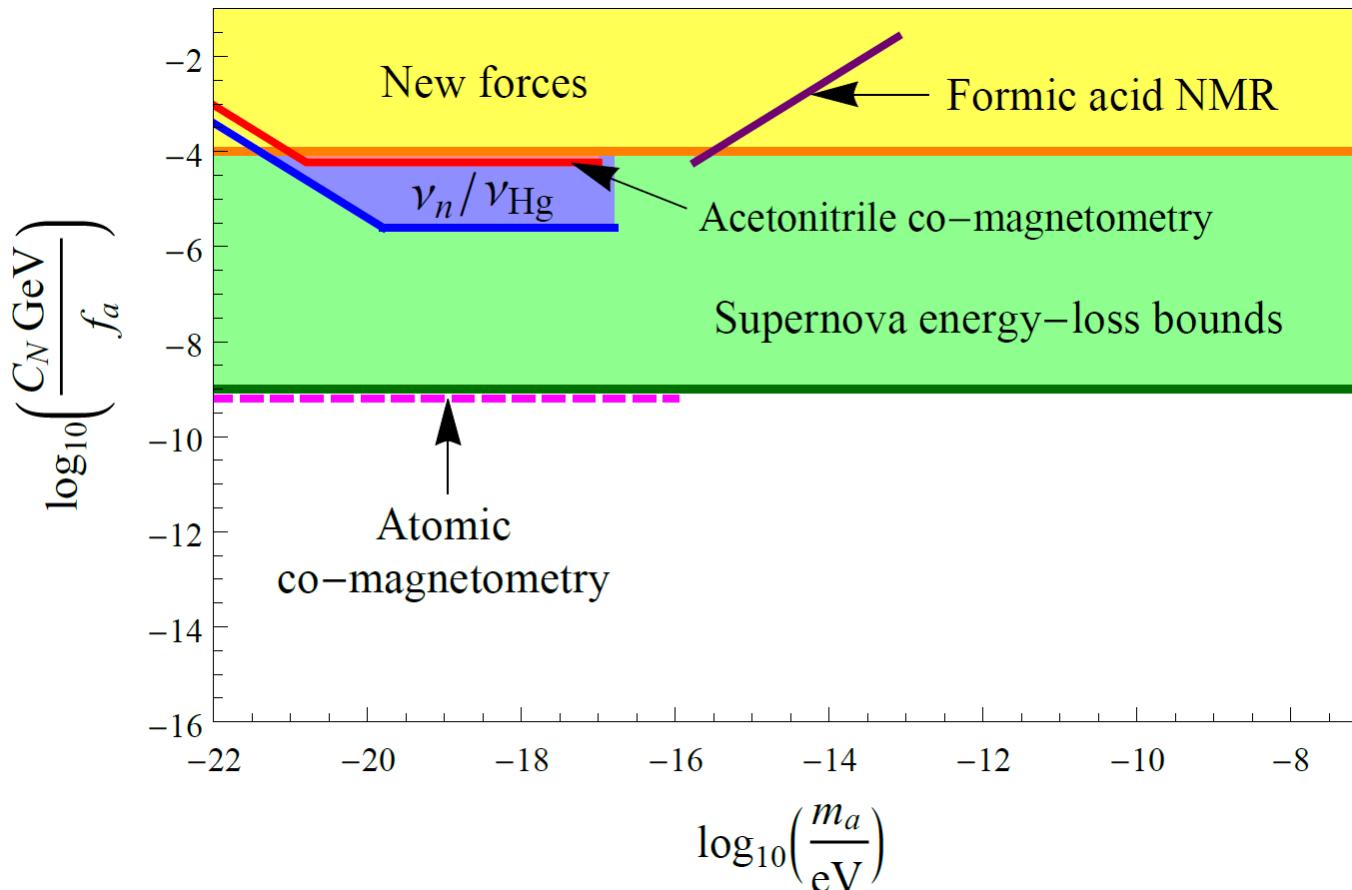


Constraints on Interaction of Axion Dark Matter with Nucleons

ν_n/ν_{Hg} constraints: [nEDM collaboration, *PRX* **7**, 041034 (2017)]

Acetonitrile constraints: [Wu *et al.*, *PRL* **122**, 191302 (2019)]

Formic acid NMR constraints: [Garcon *et al.*, *Sci. Adv.* **5**, eaax4539 (2019)]

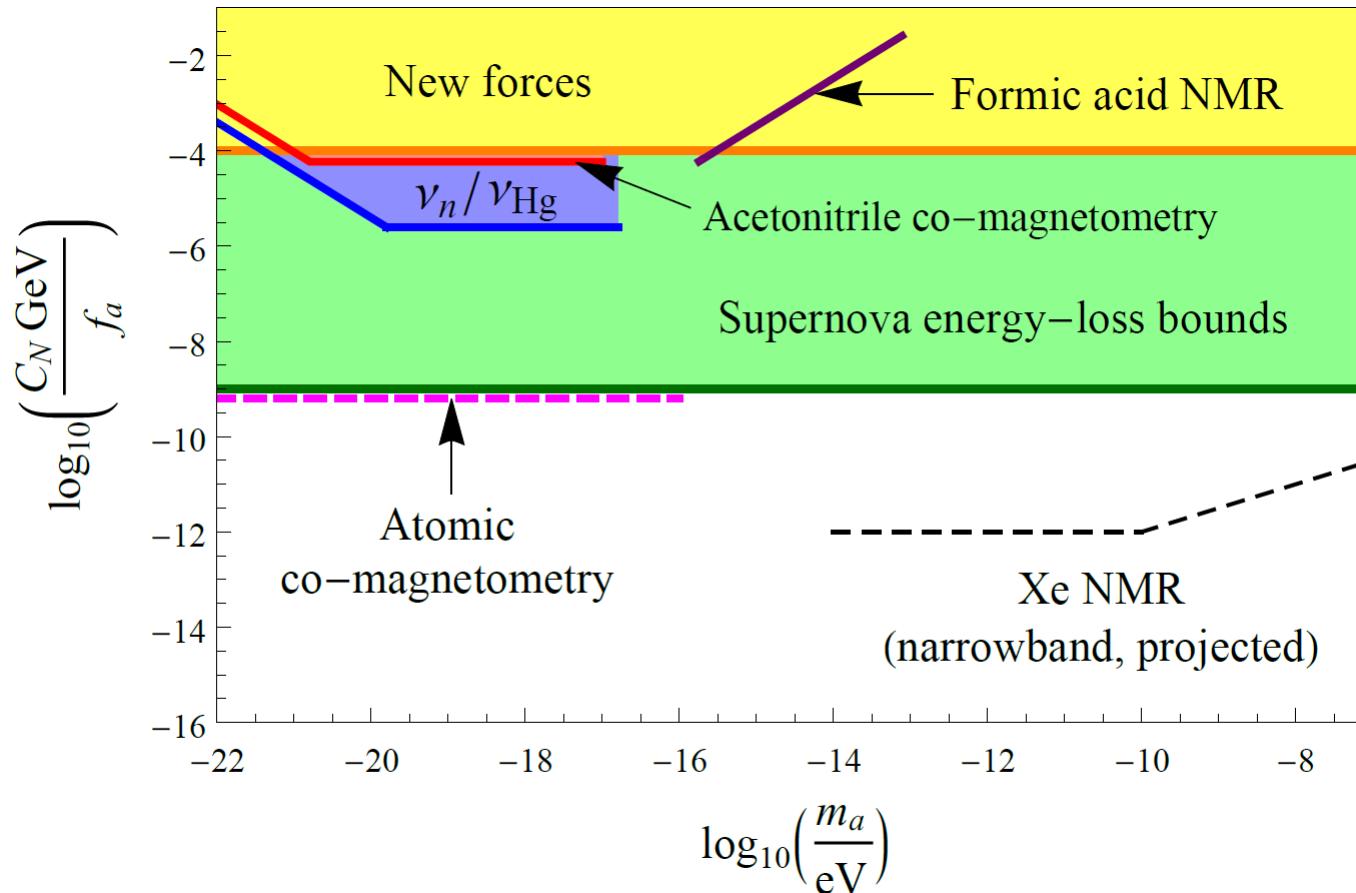


Constraints on Interaction of Axion Dark Matter with Nucleons

v_n/v_{Hg} constraints: [nEDM collaboration, *PRX* **7**, 041034 (2017)]

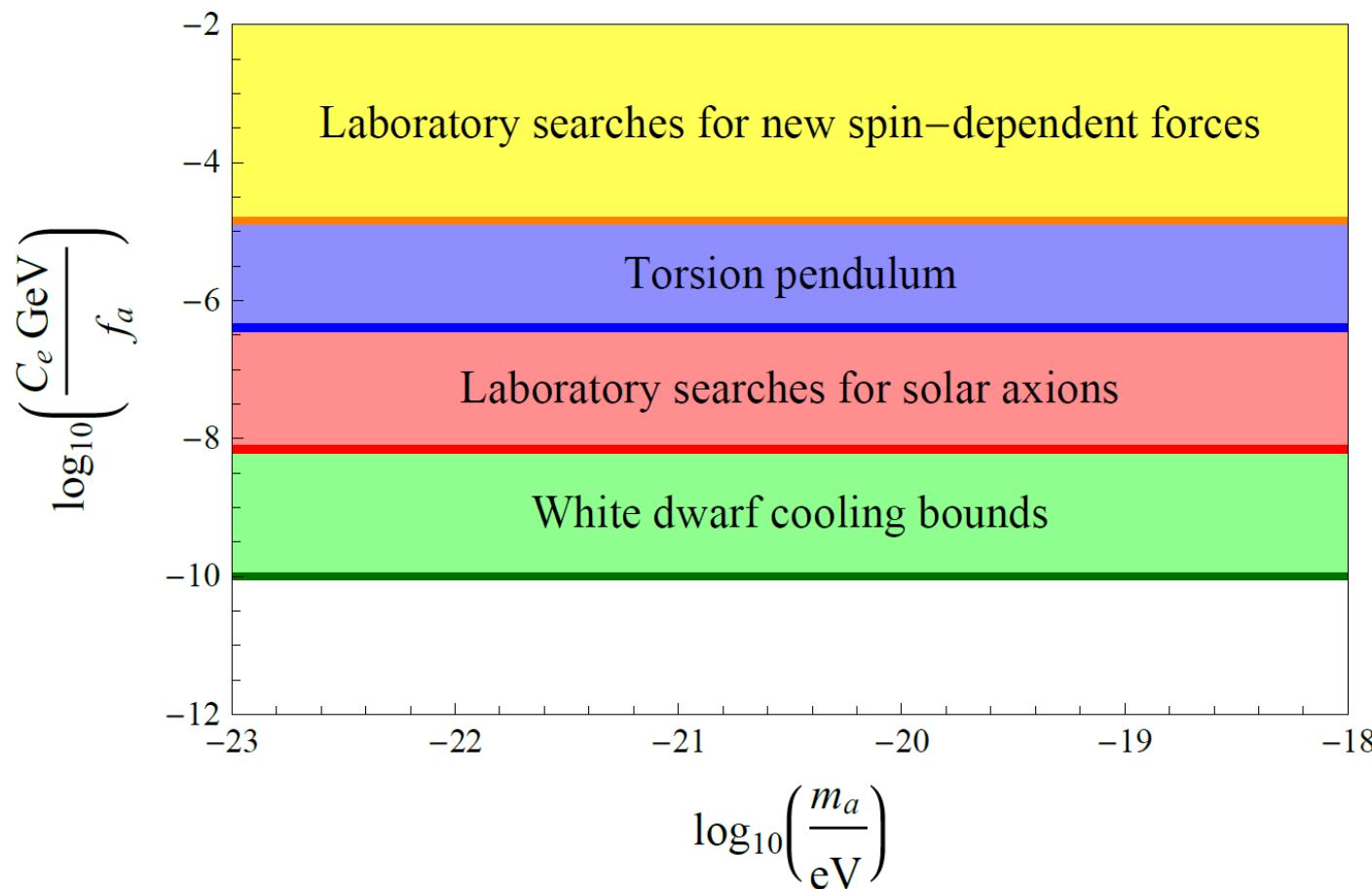
Acetonitrile constraints: [Wu *et al.*, *PRL* **122**, 191302 (2019)]

Formic acid NMR constraints: [Garcon *et al.*, *Sci. Adv.* **5**, eaax4539 (2019)]



Constraints on Interaction of Axion Dark Matter with the Electron

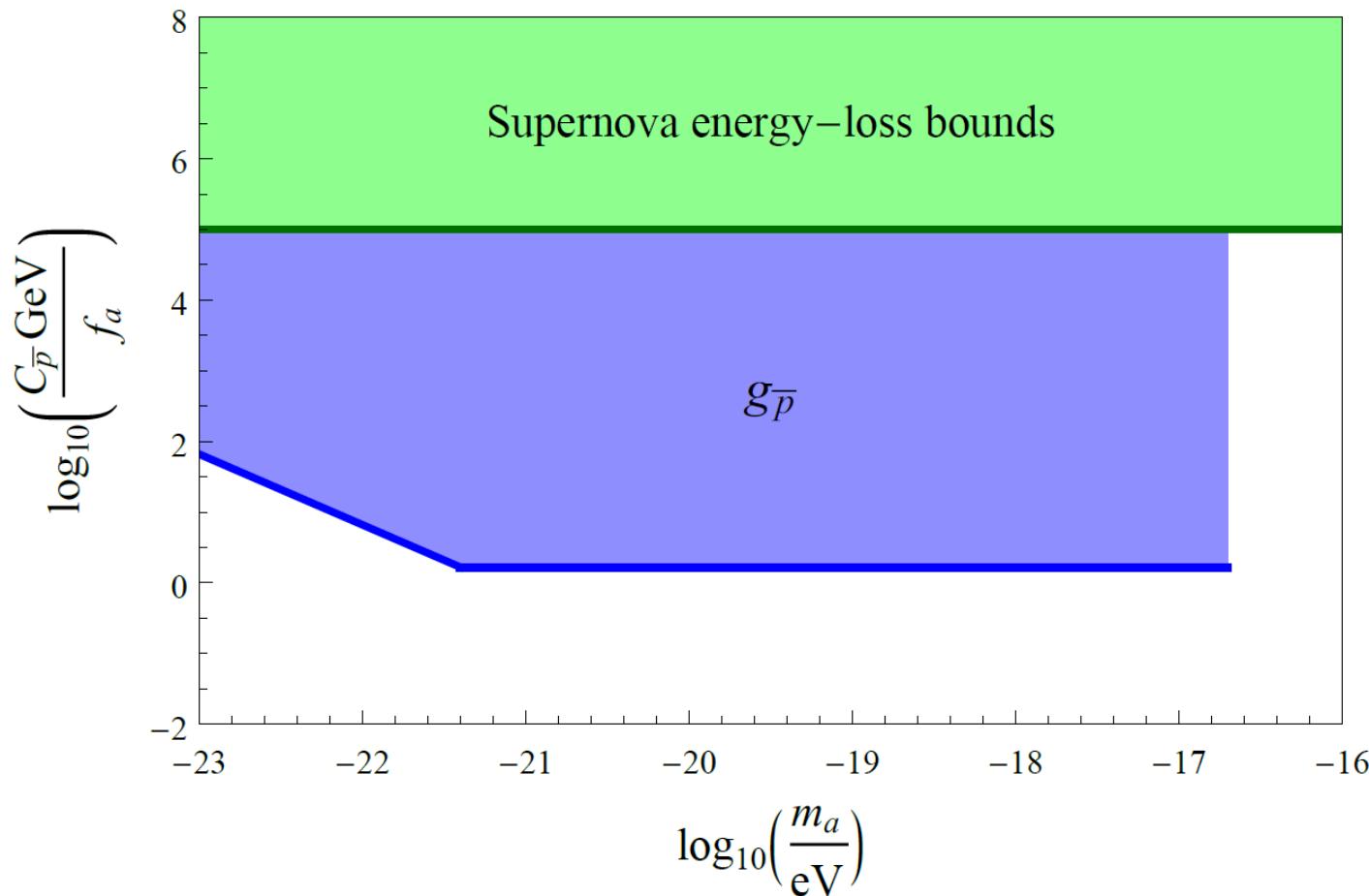
Torsion pendulum constraints: [Terrano *et al.*, *PRL* **122**, 231301 (2019)]



Constraints on Interaction of Axion Dark Matter with the Antiproton

Antiproton constraints: [BASE collaboration, *Nature* **575**, 310 (2019)]

5 orders of magnitude improvement!



Summary

- We have identified new signatures of ultra-low-mass dark matter that have allowed us and other groups to improve the sensitivity to underlying interaction strengths by up to **15 orders of magnitude**
- Novel approaches based on precision low-energy experiments (often “table-top scale”):
 - Spectroscopy (clocks) [Scalars]
 - Optical cavities and interferometry [Scalars]
 - Magnetometry and particle g -factors [Pseudoscalars]
 - Torsion pendula [Scalars and pseudoscalars]

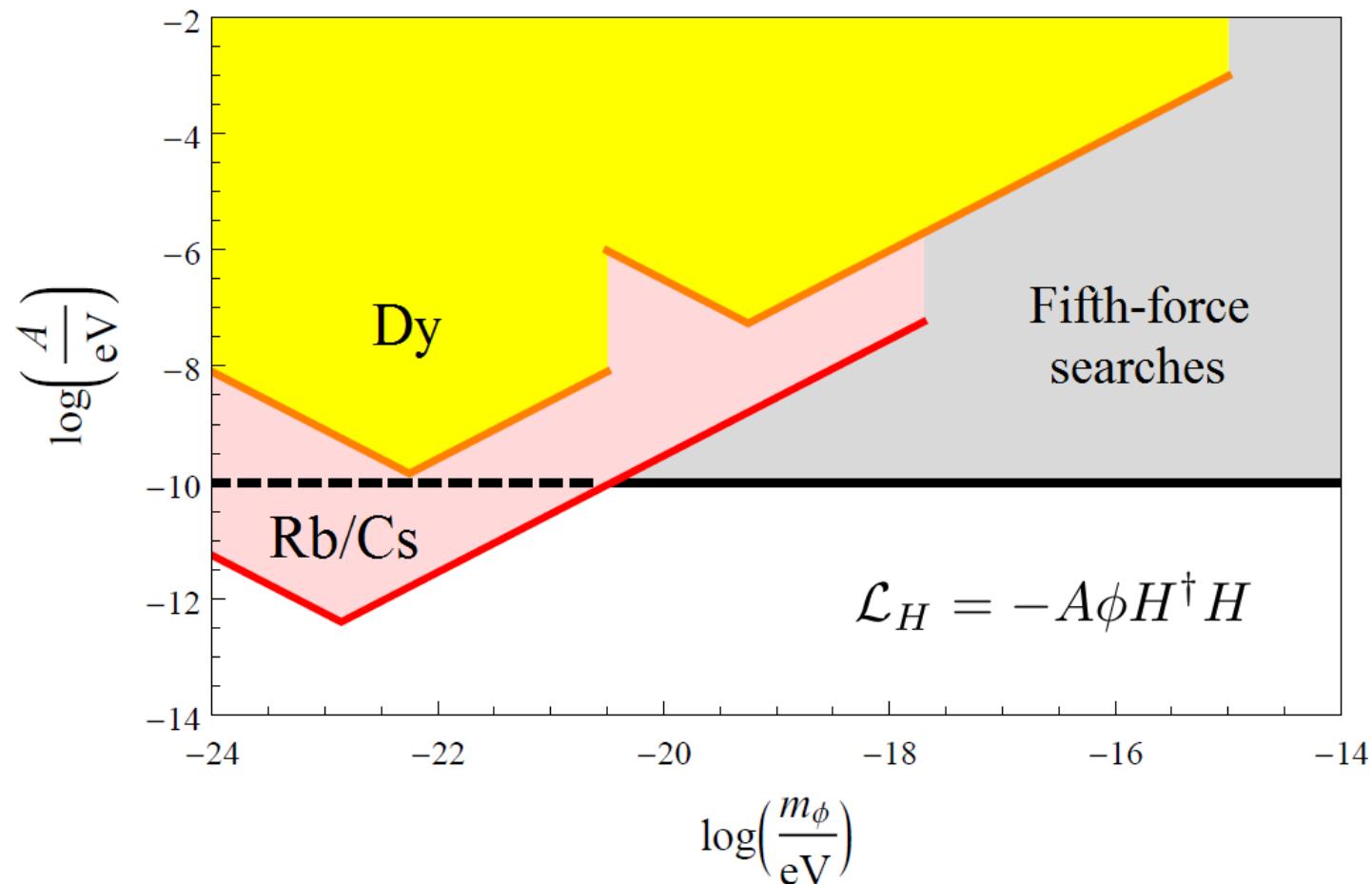
Back-Up Slides

Constraints on Linear Interaction of Scalar Dark Matter with the Higgs Boson

Rb/Cs constraints:

[Stadnik, Flambaum, *PRA* **94**, 022111 (2016)]

2 – 3 orders of magnitude improvement!

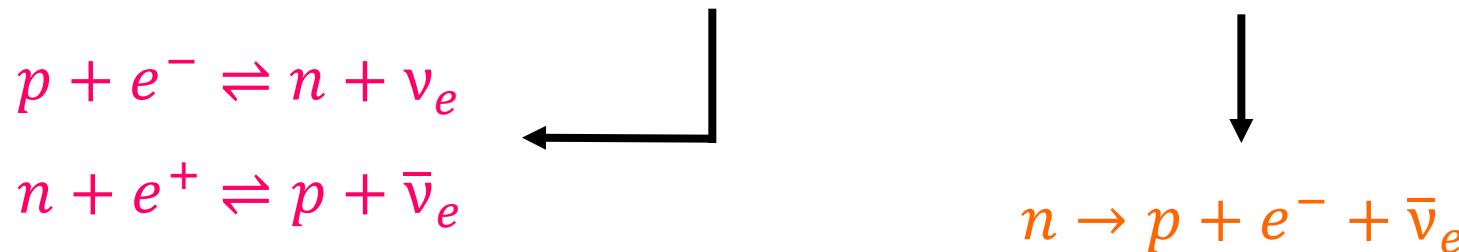


BBN Constraints on ‘Slow’ Drifts in Fundamental Constants due to Dark Matter

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

- Largest effects of DM in early Universe (highest ρ_{DM})
- Big Bang nucleosynthesis ($t_{\text{weak}} \approx 1 \text{ s} - t_{\text{BBN}} \approx 3 \text{ min}$)
- Primordial ${}^4\text{He}$ abundance sensitive to n/p ratio
(almost all neutrons bound in ${}^4\text{He}$ after BBN)

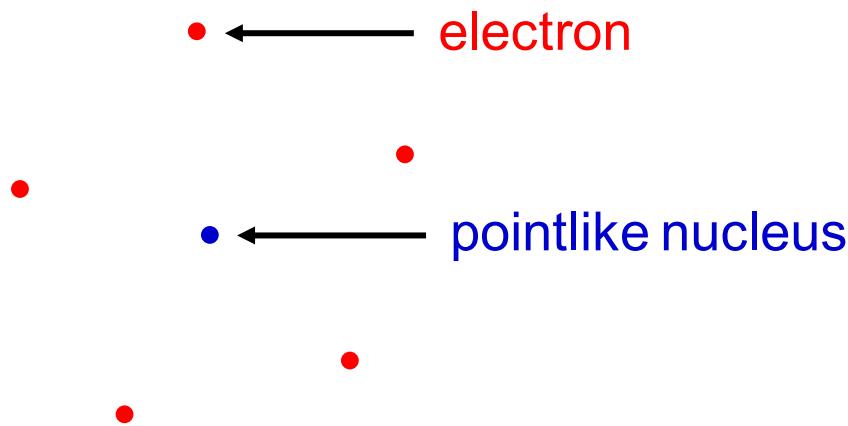
$$\frac{\Delta Y_p({}^4\text{He})}{Y_p({}^4\text{He})} \approx \frac{\Delta(n/p)_{\text{weak}}}{(n/p)_{\text{weak}}} - \Delta \left[\int_{t_{\text{weak}}}^{t_{\text{BBN}}} \Gamma_n(t) dt \right]$$



Schiff's Theorem

[Schiff, *Phys. Rev.* **132**, 2194 (1963)]

Schiff's Theorem: “In a neutral atom made up of point-like non-relativistic charged particles (interacting only electrostatically), the constituent EDMs are screened from an external electric field.”

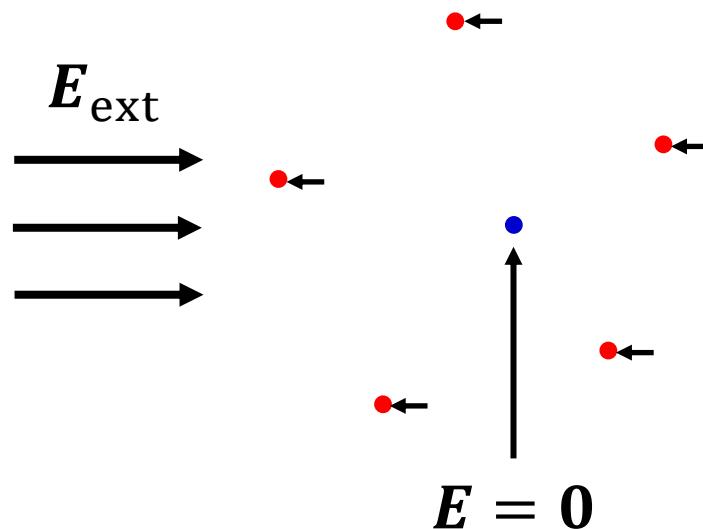


Classical explanation for nuclear EDM: A neutral atom does not accelerate in an external electric field!

Schiff's Theorem

[Schiff, *Phys. Rev.* **132**, 2194 (1963)]

Schiff's Theorem: “In a neutral atom made up of point-like non-relativistic charged particles (interacting only electrostatically), the constituent EDMs are screened from an external electric field.”



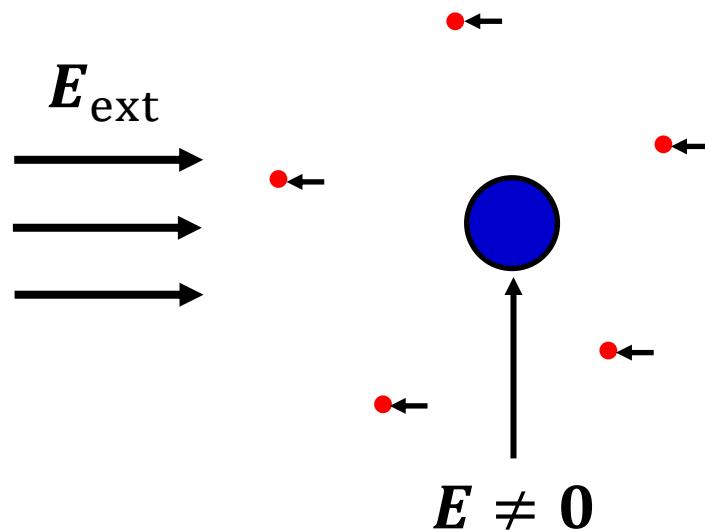
Classical explanation for nuclear EDM: A neutral atom does not accelerate in an external electric field!

Lifting of Schiff's Theorem

[Sandars, *PRL* **19**, 1396 (1967)],

[O. Sushkov, Flambaum, Khriplovich, *JETP* **60**, 873 (1984)]

In real (heavy) atoms: Incomplete screening of external electric field due to finite nuclear size, parametrised by *nuclear Schiff moment*.

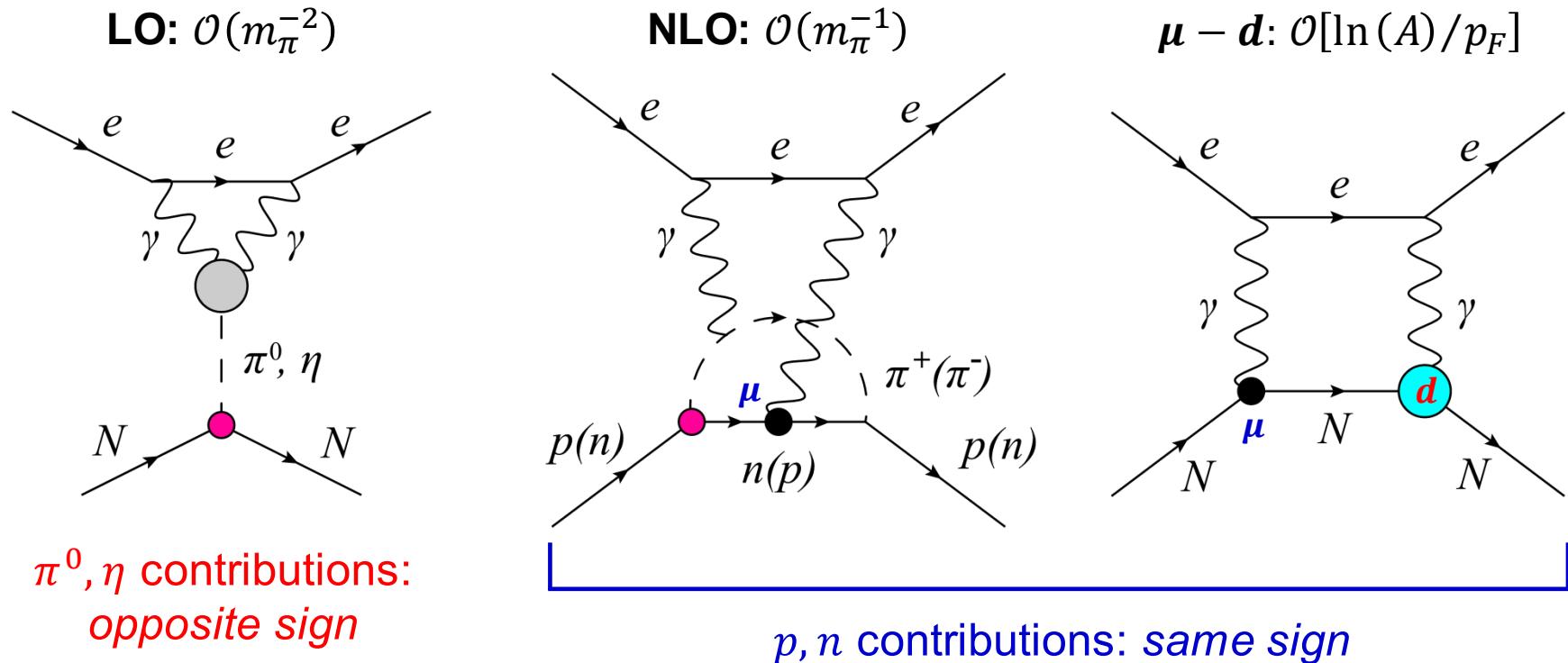


Hadronic CP Violation in Paramagnetic Molecules

[Flambaum, Pospelov, Ritz, Stadnik, *PRD* **102**, 035001 (2020)]

Hadronic CP-violating effects arise at 2-loop level, $\mathcal{O}(A)$ enhanced

Interaction of one of photons with nucleus is *magnetic* \Rightarrow no Schiff screening



Example – $\bar{\theta}_{\text{QCD}}$ term [$\bar{\theta} \leftrightarrow C_G a_0 \cos(m_a t)/f_a$]:

For $Z \sim 80$ & $A \sim 200$: $C_{\text{SP}}(\bar{\theta}) \approx [0.1_{\text{LO}} + 1.0_{\text{NLO}} + 1.7_{(\mu d)}] \times 10^{-2} \bar{\theta} \approx 0.03 \bar{\theta}$

Bounds on Hadronic CP Violation Parameters

ThO bounds: [Flambaum, Pospelov, Ritz, Stadnik, *PRD* **102**, 035001 (2020)]

| System | $ \bar{g}_{\pi NN}^{(1)} $ | $ \tilde{d}_u - \tilde{d}_d $ (cm) | $ d_p $ (e cm) | $ \bar{\theta} $ |
|--------|----------------------------|------------------------------------|-----------------------|-----------------------|
| ThO | 4×10^{-10} | 2×10^{-24} | 2×10^{-23} | 3×10^{-8} |
| n | 1.1×10^{-10} | 5×10^{-25} | — | 2.0×10^{-10} |
| Hg | 1×10^{-12} | 5×10^{-27} | 2.0×10^{-25} | 1.5×10^{-10} |
| Xe | 6.7×10^{-8} | 3×10^{-22} | 3.2×10^{-22} | 3.2×10^{-6} |

* These limits can formally be null within nuclear uncertainties

Current bounds from molecules are $\sim 10 - 100$ times weaker than from Hg & n , but are $\sim 10 - 100$ times stronger than bounds from Xe